# STATISTICS OF MULTIPLIERS FOR HYPERBOLIC RATIONAL MAPS 

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#### Abstract

In this article, we consider a counting problem for orbits of hyperbolic rational maps on the Riemann sphere, where constraints are placed on the multipliers of orbits. Using arguments from work of Dolgopyat, we consider varying and potentially shrinking intervals, and obtain a result which resembles a local central limit theorem for the logarithm of the absolute value of the multiplier and an equidistribution theorem for the holonomies.


## 1. Introduction

A major theme in the theory of dynamical systems is the study of the distribution of periodic orbits. This is particularly well-developed for hyperbolic systems, where one finds precise asymptotics and equidistribution results. Here, equidistribution can refer to spatial results, where averages of orbital measures converge to a prescribed limiting measure, or to equidistribution with respect to some symmetry of the system. This paper can be seen in the context of the second setting. A major advance in this theory came from the work of Dolgopyat [2], which introduced an approach to obtaining more precise results.

Let us now be more precise about our setting. Let $f: J \rightarrow J$ be a hyperbolic rational map restricted to its Julia set and $0<\delta<2$ to be the Hausdorff dimension of $J$. (See the next section for formal definitions.) A periodic orbit $\tau=\left\{z, f(z), \ldots, f^{n-1}(z)\right\}$ (with $f^{n}(z)=z$ ) is called primitive if $f^{m}(z) \neq z$ for all $1 \leq m<n$. We denote the set of primitive periodic orbits by $\mathcal{P}$. For each $\tau=\left\{z, f(z), \ldots, f^{n-1}(z)\right\} \in \mathcal{P}$, we define its multiplier

$$
\lambda(\tau):=\left(f^{n}\right)^{\prime}(z)
$$

and its holonomy

$$
\hat{\lambda}(\tau):=\frac{\lambda(\tau)}{|\lambda(\tau)|} \in \mathbb{S}^{1},
$$

where $\mathbb{S}^{1}$ denotes the unit circle in $\mathbb{C}$. A beautiful recent result of Oh and Winter [7] states that, apart from a small set of completely classified exceptional cases, there exists $\varepsilon>0$ such that

$$
\begin{equation*}
\#\{\tau \in \mathcal{P}:|\lambda(\tau)|<t\}=\operatorname{Li}\left(t^{\delta}\right)+O\left(t^{\delta-\varepsilon}\right) \tag{1.1}
\end{equation*}
$$

and, for any $\psi \in C^{4}\left(\mathbb{S}^{1}\right)$,

$$
\sum_{\tau \in \mathcal{P}:|\lambda(\tau)|<t} \psi(\hat{\lambda}(\tau))=\left(\int_{0}^{1} \psi\left(e^{2 \pi i \theta}\right) d \theta\right) \operatorname{Li}\left(t^{\delta}\right)+O\left(t^{\delta-\varepsilon}\right)
$$

as $t \rightarrow \infty$. Here, Li denotes the logarithmic integral $\operatorname{Li}(x)=\int_{2}^{x}(\log u)^{-1} d u \sim x / \log x$, as $x \rightarrow \infty$ and we write $f(x)=O(g(x))$ as $x \rightarrow \infty$ whenever there exists $C>0$ and $x_{0} \in \mathbb{R}$ such that for all $x \geq x_{0}$ we have that $|f(x)| \leq C g(x)$. We also write $f(x) \sim g(x)$ as $x \rightarrow \infty$ whenever $\lim _{x \rightarrow \infty} f(x) / g(x)=1$.

In this paper, we take a slightly different viewpoint. Instead of counting $\tau=\left\{z, f(z), \ldots, f^{n-1}(z)\right\}$ according to the modulus of its multiplier $|\lambda(\tau)|$, we count by the period $|\tau|=n$ but impose constraints on deviations of $\log |\lambda(\tau)|$ from the period and on the holonomy. More precisely, for $\alpha \in \mathbb{R}$, an interval $I \subset \mathbb{R}$ and an $\operatorname{arc} S \subset \mathbb{S}^{1}$, and writing $\mathcal{P}_{n}=\{\tau \in \mathcal{P}:|\tau|=n\}$, we aim to study the behaviour of

$$
\pi(n, \alpha, I, S):=\#\left\{\tau \in \mathcal{P}_{n}: \log |\lambda(\tau)|-n \alpha \in I \text { and } \hat{\lambda}(\tau) \in S\right\}
$$

as $n \rightarrow \infty$. We need to impose a restriction on $\alpha$ and, to do this, define the closed interval

$$
\mathcal{I}_{f}:=\left\{\int \log \left|f^{\prime}\right| d \mu: \mu \in \mathcal{M}_{f}\right\}
$$

where $\mathcal{M}_{f}$ is the set of $f$-invariant probability measures on $J$. We also asume that the Julia set of $f$ is not contained in a circle in $\widehat{\mathbb{C}}$ since otherwise all holonomies are real. We write $\ell$ for Lebesgue measure on $\mathbb{R}$ and $\nu$ for the normalised Haar measure on $\mathbb{S}^{1}$.

Theorem 1.1. Let $f: J \rightarrow J$ be a hyperbolic rational map of degree $d \geq 2$ restricted to its Julia set such that $J$ is not contained in a circle in $\widehat{\mathbb{C}}$. Then, for $\alpha \in \operatorname{int}\left(\mathcal{I}_{f}\right)$, there exists $\sigma_{\alpha}>0$ and $\xi_{\alpha} \in \mathbb{R}$ such that

$$
\pi(n, \alpha, I, S) \sim \frac{\nu(S)}{\sigma_{\alpha} \sqrt{2 \pi}} \int_{I} e^{-\xi_{\alpha} x} d x \frac{e^{H(\alpha) n}}{n^{3 / 2}}, \quad \text { as } n \rightarrow \infty
$$

where

$$
H(\alpha)=\sup \left\{h_{f}(\mu): \mu \in \mathcal{M}_{f} \text { and } \int \log \left|f^{\prime}\right| d \mu=\alpha\right\} .
$$

In particular, if $\alpha=\int \log \left|f^{\prime}\right| d \mu_{\max }$, where $\mu_{\max }$ is the measure of maximal entropy then

$$
\pi(n, \alpha, I, S) \sim \frac{\nu(S) \ell(I)}{\sigma_{\alpha} \sqrt{2 \pi}} \frac{d^{n}}{n^{3 / 2}} \quad \text { as } n \rightarrow \infty .
$$

We can also allow $I$ and $S$ to shrink at suitably slow rates as $n$ increases. The corresponding result will appear below as Theorem 2.1.

## 2. Hyperbolic rational maps

Let $f: \widehat{\mathbb{C}} \rightarrow \widehat{\mathbb{C}}$ be a rational map of degree $d \geq 2$. Recalling the definitions in the introduction, a periodic orbit can be classified as repelling, attracting or indifferent depending on whether its multiplier has modulus greater than, less than, or equal to one, respectively. Then, the Julia set of $f$ is defined as the closure of the union of repelling periodic orbits and denoted by $J=J(f)$. It is a compact $f^{ \pm 1}$-invariant subset of $\widehat{\mathbb{C}}$ and the reader is referred to Milnor's classical text [5] for an excellent and systematic introduction to the dynamics of functions of one complex variable. In particular, we note that such a map has topological entropy $h(f)=\log d$ and $\#\left\{z \in \mathbb{C}: f^{n}(z)=z\right\}=d^{n}$.

We say that a rational map $f: \widehat{\mathbb{C}} \rightarrow \widehat{\mathbb{C}}$ is hyperbolic if $f$ is eventually expanding on $J$, that is there exist constants $c>0$ and $\gamma>1$ such that

$$
\begin{equation*}
\left|\left(f^{n}\right)^{\prime}(z)\right| \geq c \gamma^{n} \tag{2.1}
\end{equation*}
$$

for all $z \in J$ and all $n \geq 1$.
For such a map, it is known that at most $2 d-2$ primitive periodic orbits are not repelling. Therefore, to study asymptotic counting problems for periodic orbits of $f$ we can focus, without any loss of generality, to the study of the repelling periodic orbits. We write $\delta$ for the Hausdorff dimension of $J$; this satisfies $0<\delta<2$ [18]. We will impose an additional hypothesis on $f$ : we suppose that $J$ is not contained in any circle in $\widehat{\mathbb{C}}$. In particular, this implies that $f$ is not conjugate by a Möbius transformation to a monomial $z \mapsto z^{ \pm d}$ for any $d \in \mathbb{N}$.

We will now give a more precise version of our results. As in the introduction, $\mathcal{M}_{f}$ is the set of $f$-invariant probability measures on $J$, which is convex and compact with respect to the weak* topology. Hence, the image of $\mathcal{M}_{f}$ onto the reals under the continuous projection

$$
\mu \mapsto \int_{J} \log \left|f^{\prime}\right| d \mu
$$

is an interval, which we denote by $\mathcal{I}_{f}$. Since we are assuming that $f$ is not Möbius conjugate to a monomial, $\mathcal{I}_{f}$ has non-empty interior. (If $\mathcal{I}_{f}$ is a single point then $\log \left|f^{\prime}\right|$ is cohomologous to a constant, which is incompatible with (1.1).) We define

$$
H(\alpha):=\sup \left\{h_{f}(\mu): \mu \in \mathcal{M}_{f} \text { with } \int \log \left|f^{\prime}\right| d \mu=\alpha\right\}
$$

where $h_{f}(\mu)$ denotes the measure-theoretic entropy. There is a unique $\mu_{\alpha} \in \mathcal{M}_{f}$ that realises this supremum above and a unique $\xi_{\alpha} \in \mathbb{R}$ such that

$$
h_{f}\left(\mu_{\alpha}\right)+\xi_{\alpha} \int \log \left|f^{\prime}\right| d \mu_{\alpha}=\sup \left\{h_{f}(\mu)+\xi_{\alpha} \int \log \left|f^{\prime}\right| d \mu: \mu \in \mathcal{M}_{f}\right\}
$$

We also define the variance of $\log \left|f^{\prime}\right|-\alpha$ by

$$
\sigma_{\alpha}^{2}:=\lim _{n \rightarrow \infty} \frac{1}{n} \int\left(\log \left|\left(f^{n}\right)^{\prime}\right|-n \alpha\right)^{2} d \mu_{\alpha}
$$

Our hypothesis on $f$ implies that $\sigma_{\alpha}^{2}>0$. These statements will be proved in the next section.
We want to consider the quantity $\pi(n, \alpha, I, S)$ defined in the introduction. However, we also wish to consider a situation where $I$ and $S$ shrink as $n \rightarrow \infty$. To do this, let $K \subset \mathbb{R}$ be a compact set, let $\left(I_{n}\right)_{n=1}^{\infty}$ be a sequence of intervals contained in $K$ and let $\left(S_{n}\right)_{n=1}^{\infty}$ be a sequence of arcs in $\mathbb{S}^{1}$. We are mainly interested in the two special cases where the sequences $\left(I_{n}\right)_{n=1}^{\infty}$ and $\left(S_{n}\right)_{n=1}^{\infty}$ are constant, corresponding to the case of a fixed interval and a fixed arc as in the introduction, and where the sequences $\left(\ell\left(I_{n}\right)\right)_{n=1}^{\infty}$ and $\left(\nu\left(S_{n}\right)\right)_{n=1}^{\infty}$ tend to zero, hence realising shrinking intervals. Similar asymptotic counting problems were considered in [12], [13] and [14].

We say that a sequence $\left(s_{n}\right)_{n=1}^{\infty}$ has sub-exponential growth if $\lim \sup _{n \rightarrow \infty}\left|\log s_{n}\right| / n=0$. We have the following theorem.

Theorem 2.1. Let $f: \widehat{\mathbb{C}} \rightarrow \widehat{\mathbb{C}}$ be a hyperbolic rational map of degree at least 2 such that its Julia set is not contained in a circle in $\widehat{\mathbb{C}}$. Let $K \subset \mathbb{R}$ be a compact set, let $\left(I_{n}\right)_{n=1}^{\infty}$ be a sequence of intervals in $K$ and let $\left(S_{n}\right)_{n=1}^{\infty}$ be a sequence of arcs in $\mathbb{S}^{1}$. Furthermore, suppose that $\left(\ell\left(I_{n}\right)^{-1}\right)_{n=1}^{\infty}$ and $\left(\nu\left(S_{n}\right)^{-1}\right)_{n=1}^{\infty}$ have sub-exponential growth. Then, for each $\alpha \in \operatorname{int}\left(\mathcal{I}_{f}\right)$, we have that

$$
\begin{equation*}
\pi\left(n, \alpha, I_{n}, S_{n}\right) \sim \frac{\nu\left(S_{n}\right)}{\sigma_{\alpha} \sqrt{2 \pi}} \int_{I_{n}} e^{-\xi_{\alpha} x} d x \frac{e^{H(\alpha) n}}{n^{3 / 2}}, \quad \text { as } n \rightarrow \infty \tag{2.2}
\end{equation*}
$$

In particular, if in addition we have that $\lim _{n \rightarrow \infty} \ell\left(I_{n}\right)=0$ and $p_{n} \in I_{n}$ then

$$
\begin{equation*}
\pi\left(n, \alpha, I_{n}, S_{n}\right) \sim \frac{\nu\left(S_{n}\right) \ell\left(I_{n}\right) e^{-\xi_{\alpha} p_{n}}}{\sigma_{\alpha} \sqrt{2 \pi}} \frac{e^{H(\alpha) n}}{n^{3 / 2}}, \quad \text { as } n \rightarrow \infty \tag{2.3}
\end{equation*}
$$

Corollary 2.2. If $\alpha=\int \log \left|f^{\prime}\right| d \mu_{\max }$, where $\mu_{\max }$ is the measure of maximal entropy then

$$
\begin{equation*}
\pi\left(n, \alpha, I_{n}, S_{n}\right) \sim \frac{\nu\left(S_{n}\right) \ell\left(I_{n}\right)}{\sigma_{\alpha} \sqrt{2 \pi}} \frac{d^{n}}{n^{3 / 2}}, \quad \text { as } n \rightarrow \infty \tag{2.4}
\end{equation*}
$$

## 3. ThERMODYNAMIC FORMALISM FOR HYPERBOLIC RATIONAL MAPS

The main purpose of this section is to describe how one can study the dynamics of a hyperbolic rational map using transfer operators and to obtain some decay estimates for them. We begin by recalling the essential features of this approach but for more details the reader is referred to [15]. We fix a hyperbolic rational map $f$ of degree $d \geq 2$. Further, we assume that the Julia set of $f$ is not contained inside a circle in $\widehat{\mathbb{C}}$.
3.1. Markov Partitions. For any small $\varepsilon>0$, we can find a Markov partition for $J$ : compact subsets $P_{1}, \ldots, P_{N}$ of $J$ each of diameter at most $\varepsilon$, such that
(1) $J=\bigcup_{i=1}^{N} P_{i}$,
(2) $\overline{\operatorname{int}\left(P_{i}\right)}=P_{i}, i=1, \ldots, N$,
(3) $\operatorname{int}\left(P_{i}\right) \cap \operatorname{int}\left(P_{j}\right)=\emptyset$, whenever $i \neq j$,
(4) for each $i=1, \ldots, N, f\left(P_{i}\right)=\bigcup_{j \in \mathcal{N}_{i}} P_{j}$, where $\mathcal{N}_{i}=\left\{j \in\{1, \ldots, N\}: f\left(P_{i}\right) \cap \operatorname{int}\left(P_{j}\right) \neq \emptyset\right\}$ (where closure and interior is taken relative to $J$ ).

Given a Markov partition $P_{1}, \ldots, P_{N}$, we can find open neighbourhoods $U_{j} \supset P_{j}$ such that:
(1) $f$ is injective on the closure of each $U_{j}$ and on the union $U_{i} \cup U_{j}$, whenever $U_{i} \cap U_{j} \neq \emptyset$,
(2) each $P_{i}$ is not contained in $\bigcup_{j \neq i} \overline{U_{j}}$,
(3) for each pair $i, j$ with $f\left(P_{i}\right) \supset P_{j}$ there is a local inverse $g_{i j}: U_{j} \rightarrow U_{i}$ for $f$.

We write $U=\coprod_{i=1}^{N} U_{i}$ for the disjoint union of the neighbourhoods $U_{i}$.
The structure of the partition allows us to define an $N \times N$ matrix $M$ with zero-one entries, where

$$
M_{i j}=\left\{\begin{array}{l}
1 \text { if } f\left(P_{i}\right) \supset P_{j} \\
0 \text { otherwise }
\end{array}\right.
$$

3.2. Ruelle Transfer Operators and the Pressure Function. By the hyperbolicity assumption the Julia set of $f$, and hence $U$, does not contain any critical points, that is points where the derivative of $f$ vanishes. We can therefore define the following real analytic functions related to $f$, which will help us in the study of multipliers and holonomies of periodic orbits.

Definition 3.1. We define the distortion function

$$
r(z)=\log \left|f^{\prime}(z)\right|
$$

and the rotation function

$$
\theta(z)=\arg \left(f^{\prime}(z)\right) \in \mathbb{R} / 2 \pi \mathbb{Z}
$$

which are both defined on $U$.
For a function $w: U \rightarrow \mathbb{R}($ or $\mathbb{C})$ and $n \geq 1$, we write

$$
w^{n}(z)=\sum_{j=0}^{n-1} w\left(f^{j}(z)\right)
$$

(The context should make clear that this is not an iterate.) Hence, when $\tau=\left\{z, f(z), \ldots, f^{n-1}(z)\right\} \in$ $\mathcal{P}_{n}$, we have $\lambda(\tau)=\left(f^{n}\right)^{\prime}(z)=e^{r^{n}(z)+i \theta^{n}(z)}$.

We proceed to define the Ruelle transfer operators as well as recalling some concepts from thermodynamic formalism. Write $C^{1}(U)$ for functions in $C^{1}(U, \mathbb{C})$ with bounded derivatives. Then, for $F \in C^{1}(U)$, we define the transfer operator $\mathcal{L}_{F}: C^{1}(U) \rightarrow C^{1}(U)$ by

$$
\left(\mathcal{L}_{F} w\right)(x):=\sum_{i: M_{i j}=1} e^{F\left(g_{i j} x\right)} w\left(g_{i j} x\right) \text { when } x \in U_{j}
$$

Furthermore, we define a family of modified $C^{1}$ norms on $C^{1}(U)$ by

$$
\|w\|_{(t)}:=\left\{\begin{array}{lc}
\|w\|_{\infty}+\frac{\left\|w^{\prime}\right\|_{\infty}}{t} & \text { if } t \geq 1 \\
\|w\|_{\infty}+\left\|w^{\prime}\right\|_{\infty} & \text { if } 0<t<1
\end{array}\right.
$$

The reason for this is that at the end of this section we will encounter a family of transfer operators which is not uniformly bounded using the usual $C^{1}$ norm. However, these modified norms $\|\cdot\|_{(t)}$ will help us find sufficiently good bounds for large values of $t$.

Definition 3.2. Given a continuous function $g: J \rightarrow \mathbb{R}$ we define the topological pressure of $g$ by

$$
P(g):=\sup \left\{h_{f}(\mu)+\int g d \mu: \mu \in \mathcal{M}_{f}\right\}
$$

Moreover, we call $\mu$ an equilibrium state of $g$ if $P(g)=h_{f}(\mu)+\int g d \mu$.
If $g$ is a Hölder continuous function then it has a unique equilibrium state, which is fully supported and ergodic; we denote this by $m_{g}$. Given two functions $g, h$ we have the inequality

$$
\begin{equation*}
|P(g)-P(h)| \leq\|g-h\|_{\infty} \tag{3.1}
\end{equation*}
$$

Two continuous functions $g$ and $h$ are called cohomologous if there exists a continuous function $u$ : $J \rightarrow \mathbb{R}$ such that $g-h=u \circ f-u$. If $g$ and $h$ are Hölder continuous then $m_{g}=m_{h}$ if and only if $g-h$ is cohomologous to a constant. If $g$ and $h$ are Hölder continuous then the function $\mathbb{R} \rightarrow \mathbb{R}: t \mapsto P(t g+h)$ is real analytic and

$$
\begin{align*}
\left.\frac{d P(t g+h)}{d t}\right|_{t=0} & =\int g d m_{h}  \tag{3.2}\\
\left.\frac{d^{2} P(t g+h)}{d t^{2}}\right|_{t=0} & =\lim _{n \rightarrow \infty} \frac{1}{n} \int\left(g^{n}(x)-n \int g d m_{h}\right)^{2} d m_{h} \tag{3.3}
\end{align*}
$$

see $[17,8]$. Furthermore, if $g$ is not cohomologous to a constant then $t \mapsto P(t g+h)$ is strictly convex and

$$
\begin{equation*}
\left.\frac{d^{2} P(t g+h)}{d t^{2}}\right|_{t=0}>0 \tag{3.4}
\end{equation*}
$$

We will now prove some of the statements made in the previous section. We have the following result.

Lemma 3.3. For each $\alpha \in \operatorname{int}\left(\mathcal{I}_{f}\right)$, there is a unique $\xi=\xi(\alpha) \in \mathbb{R}$ such that $H(\alpha)=h_{f}\left(m_{\xi r}\right)$ and

$$
\int r d m_{\xi r}=\alpha
$$

Proof. Since the Julia set of $f$ is not contained in a circle in $\widehat{\mathbb{C}}$, then $f$ is not conjugate to a monomial and so the distortion function $r$ is not cohomologous to a constant. Therefore, the function $\mathfrak{p}: \mathbb{R} \rightarrow \mathbb{R}$ defined by $\mathfrak{p}(t)=P(t r)$ is strictly convex.

Now consider the set

$$
\mathcal{D}:=\left\{\mathfrak{p}^{\prime}(\xi): \xi \in \mathbb{R}\right\}=\left\{\int r d m_{\xi r}: \xi \in \mathbb{R}\right\} \subset \mathcal{I}_{f}
$$

Since $\mathfrak{p}$ is strictly convex, $\mathcal{D}$ is an open interval. By the definition of pressure, for all $\mu \in \mathcal{M}_{f}$,

$$
\mathfrak{p}(t) \geq h_{f}(\mu)+t \int r d \mu
$$

In particular, the graph of the convex function $\mathfrak{p}$ lies above a line with slope $\int r d \mu$ (possibly touching it tangentially) and so $\int r d \mu \in \overline{\mathcal{D}}$. Thus, since $\mu$ is arbitrary, $\operatorname{int}\left(\mathcal{I}_{f}\right) \subset \overline{\mathcal{D}}$, and so we have $\mathcal{D}=\operatorname{int}\left(\mathcal{I}_{f}\right)$. Thus, for $\alpha \in \operatorname{int}\left(\mathcal{I}_{f}\right)$, there is a unique $\xi=\xi(\alpha) \in \mathbb{R}$ with

$$
\alpha=\mathfrak{p}^{\prime}(\xi)=\int r d m_{\xi r}
$$

Since the map $\mu \mapsto h_{f}(\mu)$ is upper semi-continuous [4], the supremum in

$$
H(\alpha)=\sup \left\{h_{f}(\mu): \int r d \mu=\alpha\right\}
$$

is attained. Since $m_{\xi r}$ is the equilibrium state for $\xi r$, we have, for any $\mu \in \mathcal{M}_{f}$ with $\mu \neq m_{\xi r}$,

$$
h_{f}\left(m_{\xi r}\right)+\xi \int r d m_{\xi r}>h_{f}(\mu)+\xi \int r d \mu .
$$

In particular, if $\int r d \mu=\alpha$ then $h_{f}\left(m_{\xi r}\right)>h_{f}(\mu)$. Therefore, $m_{\xi r}$ is the unique measure with the desired properties.

Remark 3.4. Above, we used that $r$ is not cohomologous to a constant. In fact, Oh and Winter proved a stronger statement ([7, Corollary 6.2]) that if $J$ is not contained in a circle then $r$ satisfies the non-lattice property, i.e. that it is not cohomologous to any function of the form $a+b u$, with $a, b \in \mathbb{R}$ and $u: J \rightarrow \mathbb{Z}$.

Setting $\mu_{\alpha}=m_{\xi_{\alpha} r}$, we have the measure whose existence is claimed in section 2. Furthermore,

$$
\sigma_{\alpha}^{2}=\lim _{n \rightarrow \infty} \frac{1}{n} \int\left(r^{n}-n \alpha\right)^{2} d \mu_{\alpha}=\mathfrak{p}^{\prime \prime}(\xi)>0,
$$

where we have used that $m_{\xi r}=m_{\xi(r-\alpha)}$.
For the rest of the paper, we will fix $\alpha \in \operatorname{int}\left(\mathcal{I}_{f}\right)$ and set $\xi=\xi_{\alpha}$ as in Lemma 3.3. We will also write $R:=r-\alpha$ and $R^{n}(x):=r^{n}(x)-n \alpha$ and note that, by Lemma 3.3, we have that $H(\alpha)=P(\xi R)$.

We will need to consider the function $s \mapsto e^{P(s R)}, s \in \mathbb{R}$. This function is real analytic and has an analytic extension in a neighborhood of the real line. The following lemma will prove useful in our analysis.

Lemma 3.5. [8, Proposition 4.7] The function $t \mapsto e^{P((\xi+i t) R)}$ is analytic and for some $\varepsilon>0$ we can write for each $t \in[-\varepsilon, \varepsilon]$

$$
e^{P((\xi+i t) R)}=e^{P(\xi R)}\left(1-\frac{\sigma_{\alpha}^{2} t^{2}}{2}+O\left(|t|^{3}\right)\right)
$$

where the implied constant is uniform on $[-\varepsilon, \varepsilon]$.
3.3. Decay Estimates. The approach in this subsection is motivated by Dolgopyat's seminal work on exponential mixing of Anosov flows in [2]. This was later used by Pollicott and Sharp to obtain an analogue of the prime number theorem with an exponential error term for closed geodesics on a compact negatively curved surface [11]. Naud adapted Dolgopyat's analysis to prove a similar result for convex co-compact surfaces [6] as well as Oh and Winter whose work was in the current setting of expanding rational maps [7]. We use a similar approach to obtain bounds on the spectral radii of a family of transfer operators in order to extract our asymptotic result in the final section.

Recall $f: \widehat{\mathbb{C}} \rightarrow \widehat{\mathbb{C}}$ is a hyperbolic rational map of degree at least 2 and $\alpha, \xi$ are fixed constants as in Lemma 3.3. We now consider the family of Ruelle transfer operators $\mathcal{L}_{(\xi+i b, k)}$, for $b \in \mathbb{R}$ and $k \in \mathbb{Z}$, where $\mathcal{L}_{(s, k)}:=\mathcal{L}_{s R+i k \theta}$.

We recall the following theorem.
Theorem 3.6 (Ruelle-Perron-Frobenius Theorem, [15]). Let $u \in C^{1}(U)$ be real valued. Then

- the operator $\mathcal{L}_{u}$ has a simple maximal positive eigenvalue $\lambda=e^{P(u)}$ with an associated strictly positive eigenfunction $\psi \in C^{1}(J)$,
- the rest of the spectrum is contained in a disk of radius strictly smaller than $e^{P(u)}$ and
- there is a unique probability measure $\mu$ on $J$ such that $\mathcal{L}_{u}^{*} \mu=e^{P(u)} \mu$ and $\int \psi d \mu=1$. If $v \in C^{1}(U)$ is real valued then the spectral radius of $\mathcal{L}_{u+i v}$ is bounded above by $e^{P(u)}$.

By Theorem 3.6, the spectral radius of $\mathcal{L}_{(\xi+i b, k)}$ is bounded above by $e^{P(\xi R)}$. The aim of this subsection is to show that in fact, when the Julia set of $f$ is not contained in a circle in $\widehat{\mathbb{C}}$ we can
bound the spectral radius of $\mathcal{L}_{(\xi+i b, k)}$ away from $e^{P(\xi R)}$ when $(b, k) \neq(0,0)$. To achieve this we fix arbitrary $b \in \mathbb{R}$ and $k \in \mathbb{Z}$ and consider the transfer operator $\mathcal{L}_{(\xi+i b, k)}$.

As a first step we want to consider a normalised transfer operator and we thus add a coboundary and a constant to $\xi R$. Write $v=b R+k \theta$ and

$$
u=\xi R+\log \psi-\log \psi \circ f-P(\xi R)
$$

with $\psi$ the positive eigenfunction of $\mathcal{L}_{(\xi, 0)}$ with corresponding eigenvalue $e^{P(\xi R)}$ guaranteed by Theorem 3.6. We then get that

$$
\begin{equation*}
\mathcal{L}_{u} \mathbb{1}=\mathbb{1} \text { and } \mathcal{L}_{(\xi+i b, k)}=e^{P(\xi R)} M \mathcal{L}_{u+i v} M^{-1} \tag{3.5}
\end{equation*}
$$

where $M$ is the multiplication operator by $\psi$. Thus, to show that the spectral radius of $\mathcal{L}_{(\xi+i b, k)}$ is less than $e^{P(\xi R)}$ it suffices to show that the spectral radius of $\mathcal{L}_{u+i v}$ is less than 1 . Write $F_{b, k}=u+i v$. Below we show that the spectral radius of our operator, denoted by $\operatorname{spr}\left(\mathcal{L}_{F_{b, k}}\right)$, is strictly less than 1 .

Let $\mu$ be the unique probability measure on $J$ satisfying $\mathcal{L}_{u}^{*} \mu=\mu$, as guaranteed by Theorem 3.6. We regard $\mu$ as a measure on $U$ by taking $\mu=\sum \mu_{j}$ where $\mu_{j}$ is the restriction of $\mu$ to the copy of $P_{j}$ sitting inside $U_{j}$. Since the boundary points of $\mu_{j}$, that is points in $U_{j} \backslash P_{j}$, have zero mass, $\mu$ is a probability measure on $U$.

Definition 3.7. We say that a probability measure $m$ on $J$ has the doubling property if there exists a positive constant $C$ such that for all $x \in J$ and all $\varepsilon>0$ we have that

$$
m(B(x, 2 \varepsilon)) \leq C \cdot m(B(x, \varepsilon))
$$

We know that in fact $\mu$ is a doubling measure (see [9, Theorem 8]). Moreover, as in [7, Proposition 4.5], it follows that the restrictions $\mu_{j}$ satisfy the doubling property as probability measures on $P_{j}$. We therefore have all the properties required to get the following theorem.
Theorem 3.8 (Theorem 2.7, [7]). Suppose that the Julia set of $f$ is not contained in a circle in $\widehat{\mathbb{C}}$. Then there exist $C>0$ and $\rho \in(0,1)$ such that for any $w \in C^{1}(U)$ with $\|w\|_{(|b|+|k|)} \leq 1$ and any $n \in \mathbb{N}$

$$
\left\|\mathcal{L}_{F_{b, k}}^{n} w\right\|_{L^{2}(\mu)} \leq C \rho^{n}
$$

whenever $|b|+|k| \geq 1$.
Using a standard argument (see $[2,6]$ ) we can convert the bounds on the $\|\cdot\|_{L^{2}(\nu)}$ norm to bounds for the modified $\|\cdot\|_{(t)}$ norm. Then noting that $\|\cdot\|_{C^{1}} \leq(|b|+|k|)\|\cdot\|_{(|b|+|k|)}$ for $|b|+|k| \geq 1$ we get the following corollary.

Corollary 3.9. Suppose that the Julia set of $f$ is not contained in a circle in $\widehat{\mathbb{C}}$. Then, for any $\varepsilon>0$, there exist $C_{\varepsilon}>0$ and $\rho_{\varepsilon} \in(0,1)$ such that for all $b \in \mathbb{R}$ and all $k \in \mathbb{Z}$ with $|b|+|k|>1$ we have that

$$
\left\|\mathcal{L}_{F_{b, k}}^{n}\right\|_{C^{1}} \leq C_{\varepsilon}(|b|+|k|)^{1+\varepsilon} \rho_{\varepsilon}^{n}
$$

for all $n \in \mathbb{N}$. In particular, $\operatorname{spr}\left(\mathcal{L}_{F_{b, k}}\right)<\rho_{\varepsilon}<1$.
Now that we have established the required bounds on the $C^{1}$ norm of our transfer operators we proceed to bound the sums

$$
Z_{n}(s, k):=\sum_{f^{n} x=x} e^{s R^{n}(x)+i k \theta^{n}(x)}
$$

for $s=\xi+i b$. This next result follows essentially from Ruelle's work in [16], except that we require explicit dependence on $b$ and $k$. A proof can be found in the appendix of [6] without the dependence on $k \in \mathbb{Z}$, which appeared as Proposition 6.1 in [7]. In the statement below, $\chi_{j}$ is the characteristic function of $U_{j}$ for each $1 \leq j \leq m$. (Note that, since $U$ is the disjoint union of the sets $U_{j}$, for each such $j$ we have that $\chi_{j} \in \bar{C}^{1}(U)$.)

Proposition 3.10. Fix an arbitrary $b_{0}>0$. There exists $x_{j} \in P_{j}$, for $1 \leq j \leq m$, such that for any $\eta>0$, there exists $C_{\eta}>0$ such that for all $n \geq 2$ and any $k \in \mathbb{Z}$

$$
\left|Z_{n}(\xi+i b, k)-\sum_{j=1}^{N} \mathcal{L}_{(\xi+i b, k)}^{n}\left(\chi_{j}\right)\left(x_{j}\right)\right| \leq C_{\eta}(|b|+|k|) \sum_{p=2}^{n}\left\|\mathcal{L}_{(\xi+i b, k)}^{n-p}\right\|_{C^{1}}\left(\gamma^{-1} e^{\eta+P(\xi R)}\right)^{p}
$$

for all $|b|+|k|>b_{0}$.
We are now ready to prove the decay estimates that will give us the proof of Theorem 2.1 in the next section. Fixing $\varepsilon>0$ then by Corollary 3.9 and Proposition 3.10, we get that for all $|b|+|k|>1$,

$$
\begin{aligned}
\left|Z_{n}(\xi+i b, k)\right| & \leq\left|Z_{n}(\xi+i b, k)-\sum_{j=1}^{N} \mathcal{L}_{(\xi+i b, k)}^{n}\left(\chi_{j}\right)\left(x_{j}\right)\right|+N C_{\varepsilon}(|b|+|k|)^{1+\varepsilon}\left(\rho_{\varepsilon} e^{P(\xi)}\right)^{n} \\
& \leq C_{\eta} C_{\varepsilon}(|b|+|k|)^{2+\varepsilon}\left(\rho_{\varepsilon} e^{P(\xi)}\right)^{n} \sum_{p=2}^{n}\left(\frac{e^{\eta}}{\gamma \rho_{\varepsilon}}\right)^{p}+N C_{\varepsilon}(|b|+|k|)^{1+\varepsilon}\left(\rho_{\varepsilon} e^{P(\xi)}\right)^{n}
\end{aligned}
$$

We note that it is possible to choose $1>\rho_{\varepsilon}>1 / \gamma$ (recall that $\gamma$ is the expansion rate given in (2.1)). Provided $\eta$ is small enough such that $e^{\eta} / \gamma \rho_{\varepsilon}<1$ we get that for some $C>0$

$$
\begin{equation*}
\left|Z_{n}(\xi+i b, k)\right| \leq C(|b|+|k|)^{2+\varepsilon}\left(\rho_{\varepsilon} e^{P(\xi R)}\right)^{n} \tag{3.6}
\end{equation*}
$$

Finally, we will also need a more elementary result to bound the sums $Z_{n}(\xi+i b, 0)$ for small $b \in \mathbb{R}$. These estimates can be derived as in the symbolic case in [8].

Lemma 3.11. Let $K \subset \mathbb{R}$ be a compact set. There exists $\varepsilon>0$ such that for each $n \in \mathbb{N}$ and some $\beta \in(0,1)$ we have that
(1) for $b \in K \backslash(-\varepsilon, \varepsilon)$ we can bound $Z_{n}(\xi+i b, 0)=O\left(\beta^{n} e^{H(\alpha) n}\right)$ and
(2) for $b \in(-\varepsilon, \varepsilon)$ we have

$$
Z_{n}(\xi+i b, 0)=e^{n P((\xi+i b) R)}+O\left(\beta^{n} e^{H(\alpha) n}\right)
$$

Proof. For part (1), we use the fact that, since $R$ is non-lattice (see Remark 3.4) we have that $\operatorname{spr}\left(\mathcal{L}_{(\xi+i b, 0)}\right)<e^{P(\xi R)}$ for $b \neq 0$, with a uniform bound on $K \backslash(-\varepsilon, \varepsilon)$, and Proposition 3.10. Part (2) follows from the spectral gap in the Ruelle-Perron-Frobenius theorem, which is uniform over an interval $(-\varepsilon, \varepsilon)$.

## 4. Proof of Theorem 2.1

Throughout this section we fix a hyperbolic rational map $f: \widehat{\mathbb{C}} \rightarrow \widehat{\mathbb{C}}$ of degree at least 2 . We suppose that its Julia set is not contained inside a circle in $\widehat{\mathbb{C}}$ and we fix $\alpha$ a constant in the interior of $\mathcal{I}_{f}$. We set $\xi=\xi(\alpha)$ to be the unique real number given by Lemma 3.3. Let $K \subset \mathbb{R}$ be a compact set, let $\left(I_{n}\right)_{n=1}^{\infty}$ be a sequence of intervals in $K$ and let $\left(S_{n}\right)_{n=1}^{\infty}$ be a sequence of arcs in $\mathbb{S}^{1}$. For convenience we parametirise $\mathbb{R} / \mathbb{Z}$ as $\left[-\frac{1}{2}, \frac{1}{2}\right]$ and assume that the sequence of $\operatorname{arcs}\left(S_{n}\right)_{n=1}^{\infty}$ is contained inside a fixed reference $\operatorname{arc} S=\left[-\frac{\kappa}{2}, \frac{\kappa}{2}\right]$ of length $\kappa<1$.

For each $n \in \mathbb{N}$ we denote by $p_{n}$ the midpoint of the interval $I_{n}$ and by $\vartheta_{n}$ the midpoint of the $\operatorname{arc} S_{n}$. Denote also their lengths by $\ell_{n}=\ell\left(I_{n}\right)$ and $\kappa_{n}=\nu\left(S_{n}\right)$. Furthermore, suppose that $\left(\ell_{n}^{-1}\right)_{n=1}^{\infty}$
and $\left(\kappa_{n}^{-1}\right)_{n=1}^{\infty}$ have sub-exponential growth. Then we can write

$$
\begin{aligned}
\pi\left(n, \alpha, I_{n}, S_{n}\right) & =\sum_{\tau \in \mathcal{P}_{n}} \mathbb{1}_{I_{n}}(\log |\lambda(\tau)|-n \alpha) \mathbb{1}_{S_{n}}(\hat{\lambda}(\tau)) \\
& =\sum_{\tau \in \mathcal{P}_{n}} \mathbb{1}_{\left[-\frac{1}{2}, \frac{1}{2}\right]}\left(\ell_{n}^{-1}\left(\log |\lambda(\tau)|-n \alpha-p_{n}\right)\right) \mathbb{1}_{\left[-\frac{\kappa}{2}, \frac{\kappa}{2}\right]}\left(\frac{\kappa}{\kappa_{n}}\left(\hat{\lambda}(\tau)-\vartheta_{n}\right)\right)
\end{aligned}
$$

4.1. Some auxiliary estimates. We fix $\phi \in C^{4}\left(\mathbb{R}, \mathbb{R}_{\geq 0}\right)$ compactly supported and $\psi \in C^{4}\left(\mathbb{S}^{1}, \mathbb{R}_{\geq 0}\right)$ and consider the auxiliary counting number:

$$
\pi_{\phi, \psi}(n):=\sum_{\tau \in \mathcal{P}_{n}} \phi\left(\ell_{n}^{-1}\left(\log |\lambda(\tau)|-n \alpha-p_{n}\right)\right) \psi\left(\frac{\kappa}{\kappa_{n}}\left(\hat{\lambda}(\tau)-\vartheta_{n}\right)\right)
$$

We study the asymptotic behaviour of $\pi_{\phi, \psi}$ to infer our result using an approximation argument in the next subsection.

We begin by changing the summation over $\mathcal{P}_{n}$, that is primitive periodic orbits of length $n$, to a sum over the set of fixed points of the iterated map $f^{n}$. Clearly, a primitive periodic orbit corresponds to $n$ distinct points in this set. However this set also contains points belonging in primitive periodic orbits of shorter lengths. In the following lemma we bound the error from these shorter primitive orbits.

Lemma 4.1. For all $\eta>0$ we have that

$$
\pi_{\phi, \psi}(n)=\frac{1}{n} \sum_{f^{n} x=x} \phi\left(\ell_{n}^{-1}\left(R^{n}(x)-p_{n}\right)\right) \psi\left(\frac{\kappa}{\kappa_{n}}\left(\theta^{n}(x)-\vartheta_{n}\right)\right)+O\left(e^{(H(\alpha)+\eta) n / 2}\right)
$$

Proof. Call a fixed point $x$ of the iterated map $f^{n}$ non-primitive when there exists $d$, a proper divisor of $n$, such that $f^{d} x=x$. We can then get the following bound,

$$
\begin{aligned}
& \frac{1}{n} \sum_{f^{n} x=x} \phi\left(\ell_{n}^{-1}\left(R^{n}(x)-p_{n}\right)\right) \psi\left(\frac{\kappa}{\kappa_{n}}\left(\theta^{n}(x)-\vartheta_{n}\right)\right)-\pi_{\phi, \psi}(n) \\
= & \frac{1}{n} \sum_{\substack{f^{n} x=x \\
\text { non-primitive }}} \phi\left(\ell_{n}^{-1}\left(R^{n}(x)-p_{n}\right)\right) \psi\left(\frac{\kappa}{\kappa_{n}}\left(\theta^{n}(x)-\vartheta_{n}\right)\right) \\
= & O\left(\frac{\|\psi\|_{\infty}}{n} \sum_{\substack{d \mid n \\
d \leq n / 2}} \sum_{f^{d} x=x} \phi\left(\ell_{n}^{-1}\left(R^{n}(x)-p_{n}\right)\right)\right)=O\left(\frac{1}{n} \sum_{d \leq n / 2} \sum_{f^{d} x=x} \frac{\phi\left(\ell_{n}^{-1}\left(R^{n}(x)-p_{n}\right)\right.}{e^{\xi R^{d}(x)}} e^{\xi R^{d}(x)}\right)
\end{aligned}
$$

We are only interested in periodic points which satisfy $\ell_{n}^{-1}\left(R^{n}(x)-p_{n}\right) \in \operatorname{supp} \phi$ that is $R^{n}(x) \in$ $p_{n}+\ell_{n} \operatorname{supp} \phi$. Recalling that the intervals $I_{n}$ were chosen inside a compact set $K$ we conclude that for such a periodic point the absolute value of $R^{n}(x)$ is bounded. Therefore for a non-primitive periodic point $x$, satisfying $f^{d} x=x$ for $d$ as above, we get that $R^{d}(x)=\frac{d}{n} R^{n}(x)$ and thus $e^{\xi R^{d}(x)}$ is bounded from below. From this we conclude using Lemma 3.11 that for any $\eta>0$,

$$
\begin{aligned}
& \frac{1}{n} \sum_{d \leq n / 2} \sum_{f^{d} x=x} \frac{\phi\left(\ell_{n}^{-1}\left(R^{n}(x)-p_{n}\right)\right.}{e^{\xi R^{d}(x)}} e^{\xi R^{d}(x)}=O\left(\frac{\|\phi\|_{\infty}}{n} \sum_{d \leq n / 2} \sum_{f^{d} x=x} e^{\xi R^{d}(x)}\right)=O\left(\frac{\|\phi\|_{\infty}}{n} \sum_{d \leq n / 2} Z_{d}(\xi, 0)\right) \\
= & O\left(\frac{1}{n} \sum_{d \leq n / 2} e^{(P(\xi R)+\eta) d}\right)=O\left(e^{(H(\alpha)+\eta) n / 2}\right)
\end{aligned}
$$

Consequently, we define

$$
\begin{equation*}
\tilde{\pi}_{\phi, \psi}(n)=\frac{1}{n} \sum_{f^{n} x=x} \phi\left(\ell_{n}^{-1}\left(R^{n}(x)-p_{n}\right)\right) \psi\left(\frac{\kappa}{\kappa_{n}}\left(\theta^{n}(x)-\vartheta_{n}\right)\right) . \tag{4.1}
\end{equation*}
$$

Setting

$$
\phi_{n}(x):=\phi\left(\ell_{n}^{-1}\left(x-p_{n}\right)\right) e^{-\xi\left(x-p_{n}\right)}
$$

we note that $\phi_{n} \in C^{4}\left(\mathbb{R}, \mathbb{R}_{\geq 0}\right)$ and is also compactly supported. Similarly, set

$$
\psi_{n}(x):=\psi\left(\frac{\kappa}{\kappa_{n}}\left(x-\vartheta_{n}\right)\right)
$$

In this notation we have that

$$
\tilde{\pi}_{\phi, \psi}=\frac{1}{n} \sum_{f^{n} x=x} \phi_{n}\left(R^{n}(x)\right) \psi_{n}\left(\theta^{n}(x)\right) e^{\xi\left(R^{n}(x)-p_{n}\right)}
$$

## Proposition 4.2.

$$
\tilde{\pi}_{\phi, \psi} \sim e^{-\xi p_{n}} \frac{\int \phi_{n} \int \psi_{n}}{\sigma_{\alpha} \sqrt{2 \pi}} \frac{e^{H(\alpha) n}}{n^{3 / 2}} \quad \text { as } n \rightarrow \infty
$$

To prove this proposition we consider

$$
A(n):=\left|\frac{e^{\xi p_{n}} \sigma_{\alpha} \sqrt{2 \pi n^{3}}}{e^{H(\alpha) n}} \tilde{\pi}_{\phi, \psi}-\int_{\mathbb{R}} \phi_{n} \int_{\mathbb{S}^{1}} \psi_{n}\right|
$$

and show that $A(n) \rightarrow 0$ as $n \rightarrow \infty$. The following proposition provides us with an initial bound. Using Fourier inversion and Fourier expansion we get,

$$
\begin{align*}
\phi_{n}(x) e^{\xi\left(x-p_{n}\right)} & =e^{-\xi p_{n}} \int_{-\infty}^{\infty} \hat{\phi}_{n}(t) e^{(\xi+2 \pi i t) x} d t \quad \text { and }  \tag{4.2}\\
\psi_{n}(x) & =\sum_{k \in \mathbb{Z}} c_{n, k} e^{2 \pi i k x} \tag{4.3}
\end{align*}
$$

## Proposition 4.3.

$$
A(n) \leq \frac{1}{\sqrt{2 \pi}} \int_{-\infty}^{\infty}\left|\sum_{k \in \mathbb{Z}} \frac{c_{n, k}}{e^{H(\alpha) n}} \hat{\phi}_{n}\left(\frac{t}{2 \pi \sigma_{\alpha} \sqrt{n}}\right) Z_{n}\left(\xi+\frac{i t}{\sigma_{\alpha} \sqrt{n}}, k\right)-e^{-\frac{t^{2}}{2}} \int_{\mathbb{R}} \phi_{n} \int_{\mathbb{S}^{1}} \psi_{n}\right| d t
$$

Proof. Using (4.2) and (4.3) we can get

$$
\begin{aligned}
\frac{e^{\xi p_{n}} \sigma_{\alpha} \sqrt{2 \pi n^{3}}}{e^{H(\alpha) n}} \tilde{\pi}_{\phi, \psi}(n) & =\frac{\sigma_{\alpha} \sqrt{2 \pi n}}{e^{H(\alpha) n}} \sum_{f^{n} x=x} \int_{-\infty}^{\infty} \hat{\phi}_{n}(t) e^{(\xi+2 \pi i t) R^{n}(x)} d t \sum_{k \in \mathbb{Z}} c_{n, k} e^{2 \pi i k \theta^{n}(x)} \\
& =\frac{1}{\sqrt{2 \pi}} \int_{-\infty}^{\infty} \sum_{k \in \mathbb{Z}} \frac{c_{n, k}}{e^{H(\alpha) n}} \hat{\phi}_{n}\left(\frac{t}{2 \pi \sigma_{\alpha} \sqrt{n}}\right) \sum_{f^{n} x=x} e^{\left(\xi+\frac{i t}{\sigma_{\alpha} \sqrt{n}}\right) R^{n}(x)+2 \pi i k \theta^{n}(x)} d t \\
& =\frac{1}{\sqrt{2 \pi}} \int_{-\infty}^{\infty} \sum_{k \in \mathbb{Z}} \frac{c_{n, k}}{e^{H(\alpha) n}} \hat{\phi}_{n}\left(\frac{t}{2 \pi \sigma_{\alpha} \sqrt{n}}\right) Z_{n}\left(\xi+\frac{i t}{\sigma_{\alpha} \sqrt{n}}, k\right) d t .
\end{aligned}
$$

In addition, recalling that $\int_{-\infty}^{\infty} e^{-t^{2} / 2} d t=\sqrt{2 \pi}$ we get,

$$
\begin{aligned}
\sqrt{2 \pi} A(n) & =\left|\int_{-\infty}^{\infty} \sum_{k \in \mathbb{Z}} \frac{c_{n, k}}{e^{H(\alpha) n}} \hat{\phi}_{n}\left(\frac{t}{2 \pi \sigma_{\alpha} \sqrt{n}}\right) Z_{n}\left(\xi+\frac{i t}{\sigma_{\alpha} \sqrt{n}}, k\right)-e^{-\frac{t^{2}}{2}} \int_{\mathbb{R}} \phi_{n} \int_{\mathbb{S}^{1}} \psi_{n} d t\right| \\
& \leq \int_{-\infty}^{\infty}\left|\sum_{k \in \mathbb{Z}} \frac{c_{n, k}}{e^{H(\alpha) n}} \hat{\phi}_{n}\left(\frac{t}{2 \pi \sigma_{\alpha} \sqrt{n}}\right) Z_{n}\left(\xi+\frac{i t}{\sigma_{\alpha} \sqrt{n}}, k\right)-e^{-\frac{t^{2}}{2}} \int_{\mathbb{R}} \phi_{n} \int_{\mathbb{S}^{1}} \psi_{n}\right| d t .
\end{aligned}
$$

Consider now the following

$$
\begin{aligned}
& A_{1}(n):=\int_{-\varepsilon \sigma_{\alpha} \sqrt{n}}^{\varepsilon \sigma_{\alpha} \sqrt{n}}\left|\sum_{k \in \mathbb{Z}} \frac{c_{n, k}}{e^{H(\alpha) n}} \hat{\phi}_{n}\left(\frac{t}{2 \pi \sigma_{\alpha} \sqrt{n}}\right) Z_{n}\left(\xi+\frac{i t}{\sigma_{\alpha} \sqrt{n}}, k\right)-e^{-\frac{t^{2}}{2}} \int_{\mathbb{R}} \phi_{n} \int_{\mathbb{S}^{1}} \psi_{n}\right| d t, \\
& A_{2}(n):=\int_{|t| \geq \varepsilon \sigma_{\alpha} \sqrt{n}}\left|\sum_{k \in \mathbb{Z}} \frac{c_{n, k}}{e^{H(\alpha) n}} \hat{\phi}_{n}\left(\frac{t}{2 \pi \sigma_{\alpha} \sqrt{n}}\right) Z_{n}\left(\xi+\frac{i t}{\sigma_{\alpha} \sqrt{n}}, k\right)\right| d t, \\
& A_{3}(n):=\int_{|t| \geq \varepsilon \sigma_{\alpha} \sqrt{n}}\left|e^{-\frac{t^{2}}{2}} \int_{\mathbb{R}} \phi_{n} \int_{\mathbb{S}^{1}} \psi_{n}\right| d t,
\end{aligned}
$$

with $\varepsilon>0$ small enough as in Lemmas 3.5 and 3.11. It then follows from Proposition 4.3 that

$$
A(n) \leq \frac{1}{\sqrt{2 \pi}}\left[A_{1}(n)+A_{2}(n)+A_{3}(n)\right] .
$$

We hence bound these three quantities separately to show that $\lim _{n \rightarrow \infty} A(n)=0$. To obtain these bounds we first recall a standard result from Fourier Analysis.

Lemma 4.4. If $\psi \in C^{4}\left(\mathbb{S}^{1}, \mathbb{R}\right)$ has Fourier coefficients $\left(c_{k}\right)_{k \in \mathbb{Z}}$ then $c_{0}=\int_{\mathbb{S}^{1}} \psi$ and uniformly for $\psi \in C^{4}\left(\mathbb{S}^{1}, \mathbb{R}\right)$

$$
c_{k}=O\left(\|\psi\|_{C^{4}}|k|^{-4}\right) .
$$

If $\phi \in C^{4}(\mathbb{R}, \mathbb{R})$ is compactly supported and $\hat{\phi}$ is its Fourier transform then $\hat{\phi}(0)=\int_{\mathbb{R}} \phi$ and uniformly for $\phi \in C^{4}(\mathbb{R}, \mathbb{R})$ we have that

$$
\hat{\phi}(u)=O\left(\|\phi\|_{C^{4}}|u|^{-4}\right) .
$$

These bounds follow by repeated applications of integration by parts. Now since

$$
\psi_{n}^{(q)}(x)=\left(\frac{\kappa}{\kappa_{n}}\right)^{q} \psi^{(q)}\left(\frac{\kappa}{\kappa_{n}}\left(x-\vartheta_{n}\right)\right)
$$

there exists a constant $C>0$ such that for all $n,|k| \geq 1$

$$
\begin{equation*}
\left|c_{n, k}\right| \leq C \kappa_{n}^{-4}|k|^{-4}\|\psi\|_{C^{4}} \tag{4.4}
\end{equation*}
$$

Similarly, there exists $C>0$ such that for $n \in \mathbb{N}$ and $u \in \mathbb{R}$

$$
\begin{equation*}
\left|\hat{\phi}_{n}(u)\right| \leq C \ell_{n}^{-4}|u|^{-4}\|\phi\|_{C^{4}} . \tag{4.5}
\end{equation*}
$$

Proposition 4.5. $\lim _{n \rightarrow \infty} A_{1}(n)=0$.

Proof. We can use inequality (3.6) to bound $Z_{n}\left(\xi+\frac{i t}{\sigma_{\alpha} \sqrt{n}}, k\right)$ for $k \neq 0$. Therefore fixing $\eta>0$ and recalling the bounds for the Fourier coefficients from (4.4) we get

$$
\begin{aligned}
& \int_{-\varepsilon \sigma_{\alpha} \sqrt{n}}^{\varepsilon \sigma_{\alpha} \sqrt{n}}\left|\sum_{k \neq 0} \frac{c_{n, k}}{e^{H(\alpha) n}} \hat{\phi}_{n}\left(\frac{t}{2 \pi \sigma_{\alpha} \sqrt{n}}\right) Z_{n}\left(\xi+\frac{i t}{\sigma_{\alpha} \sqrt{n}}, k\right)\right| d t \\
= & O\left(\int_{-\varepsilon \sigma_{\alpha} \sqrt{n}}^{\varepsilon \sigma_{\alpha} \sqrt{n}} \sum_{k \neq 0} \kappa_{n}^{-4}|k|^{-4}\left\|\hat{\phi}_{n}\right\|_{\infty}\left(\frac{|t|}{\sigma_{\alpha} \sqrt{n}}+|k|\right)^{2+\eta} \rho_{\eta}^{n} d t\right)=O\left(\kappa_{n}^{-4}\left\|\hat{\phi}_{n}\right\|_{\infty} \rho_{\eta}^{n}\right)
\end{aligned}
$$

for some $\rho_{\eta} \in(0,1)$. Since $\phi$ is compactly supported and $p_{n} \in K$ we can uniformly bound $\hat{\phi}_{n}$ for all $n \in \mathbb{N}$. Further, recalling that the sequence $\left(\kappa_{n}\right)_{n=1}^{\infty}$ is of sub-exponential growth we get that this error tends to zero as $n \rightarrow \infty$. Therefore, we are now left to bound

$$
\int_{-\varepsilon \sigma_{\alpha} \sqrt{n}}^{\varepsilon \sigma_{\alpha} \sqrt{n}}\left|\frac{c_{n, 0}}{e^{H(\alpha) n}} \hat{\phi}_{n}\left(\frac{t}{2 \pi \sigma_{\alpha} \sqrt{n}}\right) Z_{n}\left(\xi+\frac{i t}{\sigma_{\alpha} \sqrt{n}}, 0\right)-e^{-\frac{t^{2}}{2}} \int_{\mathbb{R}} \phi_{n} \int_{\mathbb{S}^{1}} \psi_{n}\right| d t .
$$

Using part (2) from Lemma 3.11 we get that for some $\beta \in(0,1)$

$$
\int_{\mathbb{S}^{1}} \psi_{n} \int_{-\varepsilon \sigma_{\alpha} \sqrt{n}}^{\varepsilon \sigma_{\alpha} \sqrt{n}}\left|\hat{\phi}_{n}\left(\frac{t}{2 \pi \sigma_{\alpha} \sqrt{n}}\right) e^{n\left(P\left(\left(\xi+\frac{i t}{\sigma_{\alpha} \sqrt{n}}\right) R\right)-H(\alpha)\right)}-e^{-\frac{t^{2}}{2}} \int_{\mathbb{R}} \phi_{n}\right| d t+O\left(\beta^{n}\right),
$$

On the domain of integration, we see that as $n \rightarrow \infty$
(1) $e^{n\left(P\left(\left(\xi+\frac{i t}{\sigma \alpha \sqrt{n}}\right) R\right)-H(\alpha)\right)} \rightarrow e^{-t^{2} / 2}$ by Lemma 3.5,
(2) $\hat{\phi}_{n}\left(\frac{t}{2 \pi \sigma_{\alpha} \sqrt{n}}\right) \rightarrow \hat{\phi}_{n}(0)=\int_{\mathbb{R}} \phi_{n}$ by continuity.

Furthermore, for large $n$ we have the bound $e^{n\left(P\left(\left(\xi+\frac{i t}{\sigma_{\alpha} \sqrt{n}}\right) R\right)-H(\alpha)\right)} \leq e^{-t^{2} / 4}$ and so

$$
\left|e^{n\left(P\left(\left(\xi+\frac{i t}{\sigma \alpha \sqrt{n}}\right) R\right)-H(\alpha)\right)}-e^{-t^{2} / 2}\right| \leq 2 e^{-t^{2} / 4}
$$

Finally, since $\hat{\phi}_{n}$ is uniformly bounded, we can apply the Dominated Convergence Theorem to get that $\lim _{n \rightarrow \infty} A_{1}(n)=0$.

Proposition 4.6. $\lim _{n \rightarrow \infty} A_{2}(n)=0$.
Proof.

$$
A_{2}(n) \leq \sum_{k \in \mathbb{Z}} \frac{\left|c_{n, k}\right|}{e^{H(\alpha) n}} \int_{|t| \geq \varepsilon \sigma_{\alpha} \sqrt{n}}\left|\hat{\phi}_{n}\left(\frac{t}{2 \pi \sigma_{\alpha} \sqrt{n}}\right) Z_{n}\left(\xi+\frac{i t}{\sigma_{\alpha} \sqrt{n}}, k\right)\right| d t
$$

Firstly, we use the bounds from (4.4) and (4.5). In addition, for $k \neq 0$ we use inequality (3.6) to get that a fixed $\eta \in(0,1)$ there exists $\rho_{\eta} \in(0,1)$ such that

$$
\begin{aligned}
& \sum_{k \neq 0} \frac{\left|c_{n, k}\right|}{e^{H(\alpha) n}} \int_{|t| \geq \varepsilon \sigma_{\alpha} \sqrt{n}}\left|\hat{\phi}_{n}\left(\frac{t}{2 \pi \sigma_{\alpha} \sqrt{n}}\right) Z_{n}\left(\xi+\frac{i t}{\sigma_{\alpha} \sqrt{n}}, k\right)\right| d t \\
= & O\left(\sum_{k \neq 0} \kappa_{n}^{-4}|k|^{-4} \int_{|t| \geq \varepsilon \sigma_{\alpha} \sqrt{n}}\left|\ell_{n}^{-4}\left(\frac{t}{2 \pi \sigma_{\alpha} \sqrt{n}}\right)^{-4}\left(\left|\frac{t}{\sigma_{\alpha} \sqrt{n}}\right|+|k|\right)^{2+\eta} \rho_{\eta}^{n}\right| d t\right) \\
= & O\left(\frac{n^{2} \rho_{\eta}^{n}}{\kappa_{n}^{4} \ell_{n}^{4}} \sum_{k \neq 0} \int_{|t| \geq \varepsilon \sigma_{\alpha} \sqrt{n}} \frac{\left(\left|t / \sigma_{\alpha} \sqrt{n}\right|+|k|\right)^{2+\eta}}{t^{4} k^{4}} d t\right)=O\left(\frac{n^{2}}{\kappa_{n}^{4} \ell_{n}^{4}} \rho_{\eta}^{n}\right) .
\end{aligned}
$$

On the other hand, for $k=0$ we get using part (2) of Lemma 3.11 that for some $\beta \in(0,1)$,

$$
\frac{\left|c_{n, 0}\right|}{e^{H(\alpha) n}} \int_{|t| \geq \varepsilon \sigma_{\alpha} \sqrt{n}}^{|t| \leq \sigma_{\alpha} \sqrt{n}}\left|\hat{\phi}_{n}\left(\frac{t}{2 \pi \sigma_{\alpha} \sqrt{n}}\right) Z_{n}\left(\xi+\frac{i t}{\sigma_{\alpha} \sqrt{n}}, 0\right)\right| d t=O\left(\left|c_{n, 0}\right|\left\|\hat{\phi}_{n}\right\|_{\infty} \beta^{n}\right)=O\left(\left|c_{n, 0}\right| \beta^{n}\right)
$$

since $\hat{\phi}_{n}$ is uniformly bounded across all $n \in \mathbb{N}$. We can also uniformly bound $\left|c_{n, 0}\right|$ since

$$
c_{n, 0}=\int_{\mathbb{S}^{1}} \psi_{n} \leq\|\psi\|_{\infty}
$$

Finally, as above we can use inequality (3.6) to bound the rest by the following,

$$
\begin{aligned}
& \frac{\left|c_{n, 0}\right|}{e^{H(\alpha) n}} \int_{|t| \geq \sigma_{\alpha} \sqrt{n}}\left|\hat{\phi}_{n}\left(\frac{t}{2 \pi \sigma_{\alpha} \sqrt{n}}\right) Z_{n}\left(\xi+\frac{i t}{\sigma_{\alpha} \sqrt{n}}, 0\right)\right| d t \\
= & O\left(\int_{|t| \geq \sigma_{\alpha} \sqrt{n}} \ell_{n}^{-4}\left(\frac{t}{2 \pi \sigma_{\alpha} \sqrt{n}}\right)^{-4}\left|\frac{t}{\sigma_{\alpha} \sqrt{n}}\right|^{2+\eta} \rho_{\eta}^{n} d t\right) \\
= & O\left(\frac{n^{2} \rho_{\eta}^{n}}{\ell_{n}^{4}} \int_{|t| \geq \sigma_{\alpha} \sqrt{n}}|t|^{\eta-2} d t\right)=O\left(\frac{n^{2}}{\ell_{n}^{4}} \rho_{\eta}^{n}\right) .
\end{aligned}
$$

Combining the three bounds obtained above and recalling that the sequences $\left(\ell_{n}\right)_{n=1}^{\infty}$ and $\left(\kappa_{n}\right)_{n=1}^{\infty}$ are of sub-exponential growth we obtain that $\lim _{n \rightarrow \infty} A_{2}(n)=0$.

Finally, it is clear that $\lim _{n \rightarrow \infty} A_{3}(n)=0$. This completes the proof of Proposition 4.2.
4.2. Approximation argument. Here we show how the previous auxiliary estimates provide us with the proof of Theorem 2.1 through an approximation argument. By Proposition 4.2 and Lemma 4.1 we have that for all compactly supported $\phi \in C^{4}(\mathbb{R}, \mathbb{R})$ and all $\psi \in C^{4}\left(\mathbb{S}^{1}, \mathbb{R}\right)$

$$
\begin{equation*}
\pi_{\phi, \psi}(n) \sim e^{-\xi p_{n} \int \phi_{n} \int \psi_{n}} \frac{e^{H(\alpha) n}}{\sigma_{\alpha} \sqrt{2 \pi}} \frac{n^{3 / 2}}{n^{3}} \tag{4.6}
\end{equation*}
$$

as $n \rightarrow \infty$.
Fixing $\eta>0$ we wish to construct compactly supported $\phi \in C^{4}(\mathbb{R}, \mathbb{R})$ and $\psi \in C^{4}\left(\mathbb{S}^{1}, \mathbb{R}\right)$ satisfying the following:

$$
\begin{aligned}
& \mathbb{1}_{\left[-\frac{1}{2}, \frac{1}{2}\right]} \leq \phi \leq 1+\eta, \quad \operatorname{supp}(\phi) \subset\left[-\frac{1+\eta}{2}, \frac{1+\eta}{2}\right] \text { and } \int_{\mathbb{R}} \phi \leq 1+\eta \\
& \mathbb{1}_{\left[-\frac{\kappa}{2}, \frac{\kappa}{2}\right]} \leq \psi \leq 1+\eta, \quad \operatorname{supp}(\psi) \subset\left[-\frac{\kappa+\eta}{2}, \frac{\kappa+\eta}{2}\right] \text { and } \int_{\mathbb{S}^{1}} \psi \leq \kappa+\eta
\end{aligned}
$$

A smooth function $\Phi: \mathbb{R} \rightarrow \mathbb{R}_{\geq 0}$ is called a positive mollifier, if it satisfies the following properties:
(1) it is compactly supported,
(2) $\int_{\mathbb{R}} \Phi=1$,
(3) $\lim _{\varepsilon \rightarrow 0} \Phi_{\varepsilon}(x)=\lim _{\varepsilon \rightarrow 0} \varepsilon^{-1} \Phi(x / \varepsilon)=\delta(x)$ where $\delta(x)$ is the Dirac delta function.

Let $\gamma_{1}, \ldots, \gamma_{4}>0$ and set $G=\left(1+\gamma_{1}\right) \mathbb{1}_{\left[-\frac{1}{2}-\gamma_{2}, \frac{1}{2}+\gamma_{2}\right]}$ and $H=\left(1+\gamma_{3}\right) \mathbb{1}_{\left[-\frac{\kappa}{2}-\gamma_{4}, \frac{\kappa}{2}+\gamma_{4}\right]}$. Then for sufficiently small $\varepsilon, \gamma_{1}, \ldots, \gamma_{4}>0$ the functions

$$
\phi=G * \Phi_{\varepsilon} \quad \text { and } \quad \psi=H * \Phi_{e}
$$

satisfy all the required properties. Note that since $\kappa<1$ and the constants $\varepsilon, \gamma_{4}$ were chosen sufficiently small it is harmless to assume that $\psi$ is defined on $\mathbb{R}$ rather than $\mathbb{S}^{1}$. Using (4.6) and the properties
above we can deduce that

$$
\begin{aligned}
& \limsup _{n \rightarrow \infty} \frac{\sigma_{\alpha} \sqrt{2 \pi n^{3}}}{e^{H(\alpha) n}} \pi\left(n, \alpha, I_{n}, S_{n}\right) \\
= & \limsup _{n \rightarrow \infty} \frac{\sigma_{\alpha} \sqrt{2 \pi n^{3}}}{e^{H(\alpha) n}} \sum_{\tau \in \mathcal{P}_{n}} \mathbb{1}_{\left[-\frac{1}{2}, \frac{1}{2}\right]}\left(\ell_{n}^{-1}\left(\log |\lambda(\tau)|-n \alpha-p_{n}\right)\right) \mathbb{1}_{\left[-\frac{\kappa}{2}, \frac{\kappa}{2}\right]}\left(\frac{\kappa}{\kappa_{n}}\left(\hat{\lambda}(\tau)-\vartheta_{n}\right)\right) \\
\leq & \limsup _{n \rightarrow \infty} \frac{\sigma_{\alpha} \sqrt{2 \pi n^{3}}}{e^{H(\alpha) n}} \sum_{\tau \in \mathcal{P}_{n}} \phi\left(\ell_{n}^{-1}\left(\log |\lambda(\tau)|-n \alpha-p_{n}\right)\right) \psi\left(\frac{\kappa}{\kappa_{n}}\left(\hat{\lambda}(\tau)-\vartheta_{n}\right)\right) \\
= & \limsup _{n \rightarrow \infty} e^{-\xi p_{n}} \int_{\mathbb{R}} \phi_{n} \int_{\mathbb{S}^{1}} \psi_{n} .
\end{aligned}
$$

We have

$$
\int_{\mathbb{R}} \psi_{n}=\int_{\mathbb{R}} \psi\left(\frac{\kappa}{\kappa_{n}}\left(y-\vartheta_{n}\right)\right) d y=\int_{\mathbb{R}} \psi\left(\frac{\kappa}{\kappa_{n}} y\right) d y=\frac{\kappa_{n}}{\kappa} \int_{\mathbb{R}} \psi \leq \kappa_{n}+\frac{\eta}{\kappa}=\nu\left(S_{n}\right)+O(\eta) .
$$

Similarly,

$$
\begin{aligned}
& \quad \int_{\mathbb{R}} \phi_{n}=\int_{\mathbb{R}} \phi\left(\ell_{n}^{-1}\left(x-p_{n}\right)\right) e^{-\xi\left(x-p_{n}\right)} d x=\ell_{n} \int_{\mathbb{R}} \phi(u) e^{-\xi \ell_{n} u} d u \\
& =\ell_{n} \int_{\left[-\frac{1+\eta}{2}, \frac{1+\eta}{2}\right]} \phi(u) e^{-\xi \ell_{n} u} d u \leq \ell_{n} \int_{\left[-\frac{1}{2}, \frac{1}{2}\right]} \phi(u) e^{-\xi \ell_{n} u} d u+\eta(1+\eta) e^{2(1+|\xi|)|K|}|K| \\
& \quad \ell_{n} \int_{\left[-\frac{1}{2}, \frac{1}{2}\right]} \phi(u) e^{-\xi \ell_{n} u} d u \leq \ell_{n} \int_{\left[-\frac{1}{2}, \frac{1}{2}\right]}(1+\eta) e^{-\xi \ell_{n} u} d u \leq e^{\xi p_{n}} \int_{I_{n}} e^{-\xi u} d u+\eta e^{(1+|\xi||K|}|K|
\end{aligned}
$$

Therefore,

$$
e^{-\xi p_{n}} \int_{\mathbb{R}} \phi_{n} \int_{\mathbb{S}^{1}} \psi_{n} \leq \nu\left(S_{n}\right) \int_{I_{n}} e^{-\xi u} d u+O(\eta) .
$$

Similarly, one can show that

$$
\liminf _{n \rightarrow \infty} \frac{\sigma_{\alpha} \sqrt{2 \pi n^{3}}}{e^{H(\alpha) n}} \pi\left(n, \alpha, I_{n}, S_{n}\right) \geq \liminf _{n \rightarrow \infty}\left(\nu\left(S_{n}\right) \int_{I_{n}} e^{-\xi x} d x\right)+O(\eta) .
$$

Since the choice of $\eta>0$ was arbitrary we get the result.
Assuming $\lim _{n \rightarrow \infty} \ell_{n}=0$ the derivation of the asymptotic formula (2.3) from (2.2) is straightforward. The asymptotic formula (2.4) corresponding to choosing the measure of maximal entropy follows in a similar manner. By the definition of the pressure function $\mu_{\max }$ is the equilibrium state of $\xi R$ for $\xi=0$. Then, the proof follows in the same way as above.

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