# The Frenkel-Kontorova model for quasi-periodic environments of Fibonacci type

Philippe Thieullen (joint work with E. Garibaldi and S. Petite)

Université Bordeaux 1, Institut de Mathématiques

Ergodic Theory and Dynamical Systems: Perspectives and Prospects Warwick, 16-20 April 2012

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# Outline

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# Outline

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- I. Minimizing configurations and minimizing measures in Aubry-Mather theory in the periodic case
- II. Quasi-periodic environments of Fibonacci type
- III. A few results on Aubry-Mather theory in the quasi-periodic case

# I. Minimizing configurations

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# I. Minimizing configurations



- Consider a chain of atoms in  $\mathbb{R}$ :  $x_n$  position of the nth atom
- Each atom is in interaction with its nearest neighbours and with an external potential
- The energy at each site is  $E(x_n, x_{n+1}) = W(x_{n+1} x_n) + V(x_n)$

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# The main problem

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## The main problem

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Problem: Describe the set of configurations with the lowest total energy

$$E_{tot} = \sum_{n \in \mathbb{Z}} E(x_n, x_{n+1})$$
 (the total sum is infinite)

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**Definition:** A configuration  $\{x_n\}_{n\in\mathbb{Z}}$  is minimizing in the Aubry sense if

$$E(x_n, \dots, x_{n+k}) := \sum_{i=0}^{k-1} E(x_{n+i}, x_{n+i+1})$$
  
\$\le E(y\_n, \dots, y\_{n+k})\$

whenever  $x_n=y_n$  and  $x_{n+k}=y_{n+k},$  for all  $n\in\mathbb{Z}$  and  $k\geq 1$ 

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- The environment V(x) is periodic: (of period 1)

$$V(x) = \frac{K}{(2\pi)^2} \Big( 1 - \cos(2\pi x) \Big)$$

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- The elastic interaction W is quadratic:

$$W(y-x)=\frac{1}{2}|y-x-\lambda|^2\quad (-\frac{1}{2}\lambda^2)$$

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- A more general framework:  $E_\lambda(x,y)=E_0(x,y)-\lambda(y-x)$ 

$$E_0(x, y) \text{ is of class } C^2$$

$$E_0 \text{ is periodic:} \quad E_0(x+1, y+1) = E_0(x, y)$$

$$E_0 \text{ is superlinerar:} \quad \lim_{\|y-x\| \to +\infty} \frac{E_0(x, y)}{\|y-x\|} = +\infty$$

$$E_0 \text{ is twist:} \quad \frac{\partial^2 E_0}{\partial x \partial y} < -\alpha < 0$$

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#### Theorem (Aubry 1986): (The periodic case)

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1) There exist minimizing configurations with any prescribed rotation number  $\rho$ 

$$\sup_{n\in\mathbb{Z}} |x_n - x_0 - n\rho| < +\infty$$

2) All recurrent minimizing configuration admits a rotation number

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#### Theorem (Aubry 1986): (The periodic case)

1) There exist minimizing configurations with any prescribed rotation number  $\rho$ 

$$\sup_{n\in\mathbb{Z}} |x_n - x_0 - n\rho| < +\infty$$

- 2) All recurrent minimizing configuration admits a rotation number
- 3) The main idea in the proof: a translation by an integer of a minimizing configuration is still minimizing and cannot cross itself



- It is not any more true in the quasi-periodic case

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**Remark:** If  $\{x_n\}_{n\in\mathbb{Z}}$  is minimizing in the Aubry sense then

$$\frac{\partial E}{\partial y}(x_{n-1}, x_n) + \frac{\partial E}{\partial x}(x_n, x_{n+1}) = 0, \quad \forall \ n$$

 $(x_n, x_{n+1})$  can be computed from  $(x_{n-1}, x_n)$ 

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$$x_n, x_{n+1}) \quad \text{can be computed from} \quad (x_{n-1}, x_n) = 0$$

**Definition:** A minimizing configuration can be seen as a particular orbit of a dynamical system called *Euler-Lagrange dynamics*.

Let  $v_n = x_{n+1} - x_n$ ,  $\Phi_{EL} = \begin{cases} \mathbb{T}^1 \times \mathbb{R} & \to & \mathbb{T}^1 \times \mathbb{R} \\ (x_n, v_n) & \to & (x_{n+1} = x_n + v_n, v_{n+1} = v_n + V'(x_{n+1})) \end{cases}$ 

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**Definition:** A configuration  $\{x_n\}_n$  is minimizing in the Mather sense if

$$(x_n, v_n) \in \operatorname{Supp}(\mu_{\min}), \quad \forall \ n \in \mathbb{Z}, \qquad \text{where } \mu_{\min} \text{ is minimizing}$$
$$\mu_{\min} = \arg\min\left\{\int_{\mathbb{T}^1 \times \mathbb{R}} E(x, x + v) \, d\mu(x, v) \ : \ \mu \text{ is a } \Phi_{EL} \text{-inv prob}\right\}$$

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# A small part of Mather theory

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## A small part of Mather theory

**Theorem (Mather 1991):** (the periodic case) Recall  $E_{\lambda}(x,y) = E_0(x,y) - \lambda(y-x)$ 

- 1) Any configuration minimizing  $E_{\lambda}$  in the Mather sense is minimizing in the Aubry sense
- 2) For configurations minimizing in the Aubry sense, minimizing  $E_\lambda$  is equivalent to minimizing  $E_0$
- 3) Any recurrent minimizing configuration in the Aubry sense is minimizing  $E_{\lambda}$  in the Mather sense for any  $\lambda$  related to the rotation number  $\omega$

$$\begin{split} & \omega = -\frac{d\bar{E}}{d\lambda}(\lambda) \\ & \bar{E} := \min \Big\{ \int E(x,x+v) \, d\mu(x,v) \; : \; \mu \; \Phi_{EL}\text{-inv} \Big\} \end{split}$$

# II. Environment of Fibonacci type

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# II. Environment of Fibonacci type

**Problem:** Describe the set of minimizing configurations for quasi-periodic environments of Fibonacci type

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**Problem:** Describe the set of minimizing configurations for quasi-periodic environments of Fibonacci type



–  $\mathbb R$  is partitioned into segments of two kinds: long and short

- the external potential admits two forms:  $V_L(x)$  and  $V_S(x)$
- $-\underline{\Omega} =$  the closure of all the shifts of the Fibonacci word

$$\ldots, LSL, LS | LS, L, LS, LSL, LSLLS, \ldots$$

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**Notations:**  $\underline{\Omega}$  is the compact set of Fibonacci words.  $\underline{\Omega}$  is compact minimal and uniquely ergodic. Each  $\underline{\omega} \in \underline{\Omega}$  gives a quasi-periodic potential  $V_{\underline{\omega}}(x)$ . As before, the total energy per site is

$$E_{\underline{\omega}}(x,y) = W(y-x) + V_{\underline{\omega}}(x), \qquad W''(x) < -\alpha < 0$$

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Theorem (Gambaudo, Guiraud, Petite, 2006): We fixe an environment  $\underline{\omega} \in \underline{\Omega}$ .

- Any minimizing configuration in the Aubry sense has a rotation number

$$\rho = \lim_{m-n \to +\infty} \frac{x_m - x_n}{m - n}$$

- Any rotation number is achieved

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#### Question:

What about minimizing configurations in the Mather sense? What plays the role of  $\mathbb{T}^1\times\mathbb{R}$  in the periodic case?

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# The space of quasi-periodic environments

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# The space of quasi-periodic environments Extension of $\Omega$ :



- the origin does not play any role. We consider the set of all shifts

$$\Omega = \underline{\Omega} \times \mathbb{R} / \sim \quad \Leftrightarrow \quad \begin{cases} \text{different parametrizations but} \\ \text{same sequence of impurities} \end{cases}$$

- $\Omega$  is a suspension over  $\underline{\Omega}$  built with a return map of length L or S
- In the periodic case L = S and  $\Omega = \mathbb{T}^1$
- In the quasi-periodic case  $\Omega$  plays the role of  $\mathbb{T}^1$

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## Mather measures

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#### Mather measures

A global potential V: Recall  $(\Omega, \{\phi^t\})$  denotes the minimal Fibonacci flow

$$V_{\omega}(x) = V \circ \phi^{x}(\omega)$$
  

$$E_{\omega}(x, y) = W(y - x) + V_{\omega}(x)$$
  

$$= L(\phi^{x}(\omega), y - x)$$
  

$$L(\omega, v) = W(v) + V(\omega)$$



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#### Minimizing measures in the Mather sense:

- There is no way we can define an equivalent Euler-Lagrange map  $\Phi_{EL}$
- In the periodic case:  $\Phi_{EL}(x,v) = (x+v,\ldots)$ ,  $x \in \mathbb{T}^1$ ,  $v \in \mathbb{R}$
- A measure  $\mu$  is holonomic if

$$\int f(\omega) \, d\mu(\omega, v) = \int f \circ \phi^v(\omega) \, d\mu(\omega, v), \qquad \forall \ f \in C^0(\Omega)$$

– A measure  $\mu_{min}$  is minimizing in the Mather sense if

$$\mu_{min} = \arg\min\left\{\int_{\Omega\times\mathbb{R}} L(\omega, v) \, d\mu(\omega, v) : \mu \text{ is holonomic} \right\}_{\langle \Box \rangle \land \langle \Box \rangle \land \langle \Box \rangle \land \langle \Xi \land \langle \Xi \rangle \land \langle \Xi \land \Box \langle \Xi \rangle \land \langle \Xi \land \langle \Xi \rangle \land \langle \Xi \rangle \land \langle \Xi \land \Box \land \langle \Xi \land \Box \land$$

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## III. A few results in Aubry-Mather theory

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#### III. A few results in Aubry-Mather theory Mather set: $M := \bigcup \{ \text{Supp}(\mu) : \text{holonomic minimizing} \} \subset \Omega \times \mathbb{R}$

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#### III. A few results in Aubry-Mather theory Mather set: $M := \bigcup \{ \text{Supp}(\mu) : \text{holonomic minimizing} \} \subset \Omega \times \mathbb{R}$

#### **Results:**

– Optimal segments  $\{x_k\}_k$  (minimizing  $E(x_0, \ldots, x_n)$ ) have uniform bounded gaps

$$\exists C \text{ s.t. } \{x_k\}_{k=0}^n \text{ is optimal } \Rightarrow |x_{k+1} - x_k| < C$$

- The Mather set is compact and non empty

– The lowest mean energy can be computed using either minimizing configurations or minimizing measures  $(L(\omega, v) = W(v) + V(\omega))$ 

$$\bar{E} = \lim_{n \to +\infty} \min_{\omega, x_0, \dots, x_n} \frac{1}{n} E_{\omega}(x_0, \dots, x_n)$$
$$= \min_{\mu \text{ holonomic}} \int L(\omega, v) \, d\mu(\omega, v) = \int L \, d\mu_{min}$$

- If  $\tilde{M}$  denotes the projection of the Mather set on  $\Omega$ , then  $\tilde{M}$  has a non empty intersection with any orbit of length long enough of the Fibonacci flow.