# New physics and CP violation measurements at LHCb

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#### Contents

- The CKM Matrix and CP violation
  - The need for a clean measurement of  $\boldsymbol{\gamma}$
- Flavour oscillations and CP violation
  - Where are we now and where will LHCb take us?
  - $B_{d}$ ,  $B_{s}$  and charm systems
- Alternative  $\gamma,\,\beta$  and  $\beta_{_{\!\varsigma}}$  measurements



#### The Cabibbo-Kobayashi-Maskawa Quark Mixing Matrix





$$V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

- A 3x3 unitary matrix
- Described by 4 real parameters allows CP violation
  - PDG (Chau-Keung) parametrisation:  $\theta_{p}$ ,  $\theta_{g}$ ,  $\theta_{g}$ ,  $\theta_{g}$ ,  $\delta_{g}$ ,  $\delta$
  - Wolfenstein parametrisation:  $\lambda$ , A,  $\rho$ ,  $\eta$
- Highly predictive

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### **CP** violation and New Physics

- CKM ansatze is highly predictive
- Measurements of flavour phenomena overconstrain the SM parameters
- Consistency of different measurements tests for presence of new physics
- For example, compare
  - Clean SM reference points (from theory or measurement)
  - New physics sensitive measurements
- Recall past success of CP violation in discovering new physics



## CKM Matrix – Phases

P.Harrison *et al.*, arXiv:0904.3077 [hep-ph]

- Can form a matrix of angles between pairs of CKM matrix elements
  - $\Phi_{ij}$  = phase between remaining elements when row i and column j removed

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• unitarity implies sum of phases in any row or column = 180°

## CKM Matrix – Phases

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• unitarity implies sum of phases in any row or column = 180°

$$\Phi = \begin{pmatrix} d & s & b & \beta_s \approx \varphi_s/2 & d & s & b & \beta \equiv \varphi_1 \\ \Phi_{ud} & \Phi_{ud} & \Phi_{ub} & \Phi_{ub} \\ \Phi_{cd} & \Phi_{cs} & \Phi_{cb} \\ \Phi_{td} & \Phi_{ts} & \Phi_{tb} \end{pmatrix} \simeq \begin{pmatrix} 1^{\circ} & 22^{\circ} & 157^{\circ} \\ 67^{\circ} & 90^{\circ} & 23^{\circ} \\ 112^{\circ} & 68^{\circ} & 0^{\circ} \end{pmatrix} \qquad \alpha \equiv \varphi_2 \\ \varphi \equiv \varphi_3 & \varphi_2 & \varphi_3 &$$

#### The Unitarity Triangle





#### Today's Constraints on the Unitarity Triangle



#### Importance of $\gamma$

• γ plays a unique role in flavour physics

the only CP violating parameter that can be measured through tree decays \*

(\*) more-or-less

- A benchmark Standard Model reference point
  - doubly important after New Physics is observed
- How precise is precise enough?

0.1%

Gerch

(X)

- 10% 🛞 At 3 sigma hardly exclude anything
- 1%  $\Rightarrow$  Seems the right level to test NP
  - Good luck if you can get the funding ...

#### How To Measure y

- Focus on theoretically pristine measurement
  - Interference between



- colour allowed



- colour suppressed
- final state contains  $D^0$  final state contains  $\overline{D}^0$
- Use D decay final states accessible to both amplitudes
  - KK (GLW), Kπ (ADS), Κ<sub>ε</sub>ππ (GGSZ)
- Use also B<sub>d</sub> decays to DK<sup>\*0</sup> (untagged)

• Use also B decays to D K (tagged, time-dependent) Tim Gersho Physics & CP Violatio

#### LHCb sensitivity to y

$\delta_{B^0}$ (°)	0	45	90	135	180
$\sigma_{\gamma}$ for 0.5 fb <sup>-1</sup> (°)	8.1	10.1	9.3	9.5	7.8
$\sigma_{\gamma}$ for 2 fb <sup>-1</sup> (°)	4.1	5.1	4.8	5.1	3.9

- Numbers assume nominal LHC performance
- Sensitivity to  $\delta_m$  inherent to  $B^0 \rightarrow DK^0$  ("quasi-two-body") analysis
- Precision can be further improved:
  - CLEOc results on  $D \rightarrow K\pi\pi^0$  allow it to be used in ADS analysis
  - $B^{\scriptscriptstyle 0} \to DK\pi$  Dalitz plot analysis gives improved sensitivity to  $\gamma$  with reduced dependence on  $\delta_{_{\rm B}}$



#### Flavour oscillations, CP violation and Nobel Prizes

- 1964 Discovery of CP violation in K<sup>0</sup> system
- 1980 Nobel Prize to Cronin and Fitch









- 2001 Discovery of CP violation in B<sub>d</sub> system
- 2008 Nobel Prize to Kobayashi and Maskawa



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Prog.Theor.Phys. 49 (1973) 652





#### Flavour oscillations, CP violation and Nobel Prizes

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- Generic (but shown for  $\mathsf{B}_{\mathrm{s}}$ ) decays to CP eigenstates

$$\begin{split} \Gamma(B_s(t) \to f) &= \mathcal{N}_f \, |A_f|^2 \, \frac{1 + |\lambda_f|^2}{2} \, e^{-\Gamma t} \\ &\times \left[ \cosh \frac{\Delta \Gamma t}{2} + \mathcal{A}_{\rm CP}^{\rm dir} \, \cos(\Delta m \, t) + \mathcal{A}_{\Delta \Gamma} \, \sinh \frac{\Delta \Gamma t}{2} + \mathcal{A}_{\rm CP}^{\rm mix} \, \sin(\Delta m \, t) \right] \\ \Gamma(\overline{B}_s(t) \to f) &= \mathcal{N}_f \, |A_f|^2 \, \frac{1 + |\lambda_f|^2}{2} \, (1 + a) \, e^{-\Gamma t} \\ &\times \left[ \cosh \frac{\Delta \Gamma t}{2} - \mathcal{A}_{\rm CP}^{\rm dir} \, \cos(\Delta m \, t) + \mathcal{A}_{\Delta \Gamma} \, \sinh \frac{\Delta \Gamma t}{2} - \mathcal{A}_{\rm CP}^{\rm mix} \, \sin(\Delta m \, t) \right]. \end{split}$$



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CP violating asymmetries
CP conserving parameter

$$(A_{CP}^{dir})^2 + (A_{\Delta\Gamma})^2 + (A_{CP}^{mix})^2 = 1$$



• Generic (but shown for  $B_0$ ) decays to CP eigenstates



Untagged analyses still sensitive to some interesting physics



- Generic (but shown for  $\mathsf{B}_{\mathrm{s}}$ ) decays to CP eigenstates

$$\begin{split} \Gamma(B_s(t) \to f) &= \mathcal{N}_f \, |A_f|^2 \, \frac{1 + |\lambda_f|^2}{2} \, e^{-\Gamma t} \\ &\times \left[ \cosh \frac{\Delta \Gamma t}{2} + \underbrace{\mathbf{0}}_{2} + \mathcal{A}_{\Delta \Gamma} \sinh \frac{\Delta \Gamma t}{2} + \mathcal{A}_{\mathrm{CP}}^{\mathrm{mix}} \sin \left( \Delta m t \right) \right] \\ \Gamma(\overline{B}_s(t) \to f) &= \mathcal{N}_f \, |A_f|^2 \, \frac{1 + |\lambda_f|^2}{2} \left( 1 + \underbrace{\mathbf{0}}_{2} e^{-\Gamma t} \right) \\ &\times \left[ \cosh \frac{\Delta \Gamma t}{2} - \underbrace{\mathbf{0}}_{2} + \mathcal{A}_{\Delta \Gamma} \sinh \frac{\Delta \Gamma t}{2} - \mathcal{A}_{\mathrm{CP}}^{\mathrm{mix}} \sin \left( \Delta m t \right) \right]. \end{split}$$

- In some channels, expect no direct CP violation
- and/or no CP violation in mixing



- Generic (but shown for  $\mathsf{B}_{\underline{0}}$ ) decays to CP eigenstates



- In some channels, expect no direct CP violation
- $B_d$  case:  $\Delta\Gamma$  negligible



• Generic (but shown for  $B_0$ ) decays to CP eigenstates



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•  $D^0$  case: both x =  $\Delta m/\Gamma$  and y= $\Delta \Gamma/2\Gamma$  small

#### (Almost) latest measurements of $\beta$ $\rightarrow J/\psi K^0$ BABAR BELLE



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PRL 98 (2007) 031802

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#### (Almost) latest measurements of β



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## LHCb prospects for $\beta$

• "Yesterday's sensation is tomorrow's calibration"



• Sensitivity  $\sigma(sin(2\beta)) \approx 0.06$  with 200/pb of data

• Validate selection, vertexing, resolution, tagging Tim Gershon New Physics & CP Violation

#### Charm mixing and CP violation

Including results from BABAR, Belle, CDF, CLEO(c), FOCUS

Latest new results Belle arXiv:0905.4185 [hep-ex]

BABAR arXiv:0908.0761 [hep-ex]



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- Mixing established (though still no single measurement >  $5\sigma$ ) <sub>23</sub>
  - No indication of CP violation

#### LHCb Prospects for Charm

- Most sensitive channels
  - $D \rightarrow K^{\scriptscriptstyle +} K^{\scriptscriptstyle -}, \ \pi^{\scriptscriptstyle +} K^{\scriptscriptstyle -}, \ K_{\scriptscriptstyle e} \pi^{\scriptscriptstyle +} \pi^{\scriptscriptstyle -}$
  - All well-suited to LHCb
- Tagging provided by  $D^{\star^{\pm}} \to D\pi^{\pm}$
- Huge cross-section for prompt charm production
  - plus large rates of secondary charm
- With 100/pb accumulate ~ 2.5M tagged  $D \to K^{^{+}}K^{^{-}}$



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# ${\rm B}_{\rm s}$ oscillations and CP violation

- Tevatron measurements using tagged  $B_{_{c}} \rightarrow J/\psi\phi$
- Angular analyses of vector-vector final state
- Results depend on  $\Delta\Gamma$

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# ${\rm B}_{\rm s}$ oscillations and CP violation



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## Complications of $B_s \rightarrow J/\psi \phi$



- VV final state
  - three helicity amplitudes mixture of CP-even and CP-odd
  - can be disentangled using angular & time-dependent distributions → additional sensitivity
- Φ width not negligible
  - contribution from KK S-wave can be handled in the analysis
  - $B_{s} \rightarrow J/\psi f_{0}$ ,  $f_{0} \rightarrow \pi\pi$  can also be studied (CP-eigenstate)

#### Aspects of $B_s \rightarrow J/\psi\phi$ analysis Full MC: inclusive $J/\psi(\mu\mu)$

Dunh

5300 5400 5500 5600 5700

u+u\*K KInvariant Mass (MeV/c2)

Norm Gauss

Mean Gauss Signa Gauss 0.07106

9 14±4.1 1409±1279.2 0.0000133±0.0001605

4 02 + 4 39

signal

bkg prompt

5100 5200

vents/10 MeV/c

- Selection
  - O(10<sup>5</sup>) signal events/2/fb
  - High efficiency for dimuon trigger
  - Largest background from prompt J/ψ
- Use of control samples







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If true  $\beta_s$  value is at current Tevatron central value, can measure it with 200/pb (even with reduced LHC energy)

#### Alternative approaches to $\gamma$ , $\beta$ and $\beta_s$

- All methods discussed so far use  $b \rightarrow c$  decay transitions
- Search for new physics by testing if b → s loop dominated decays give consistent results
- $\gamma: B_{g} \rightarrow hh' decays$

• eg. 
$$B_{s} \rightarrow K^{+}K^{-}$$
 cf.  $B_{d} \rightarrow \pi^{+}\pi^{-}$  via U-spin





## Alternative approaches to $\gamma,\,\beta$ and $\beta_{_{S}}$

- All methods discussed so far use  $b \rightarrow c$  decay transitions
- Search for new physics by testing if  $b \rightarrow s$  loop dominated decays give consistent results  $sin(2\beta^{eff}) \equiv sin(2\phi_1^{eff}) \bigoplus_{FPCP 2009} sin(2\phi_1^{eff}) \oplus_{FPCP 2009} sin(2\phi_1^{eff})$
- $\beta$  and  $\beta_s$

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• Promising sensitivity for

$$- B_{d} \rightarrow \phi K_{s}$$
$$- B_{s} \rightarrow \phi \phi$$



#### Conclusion

LHCb well placed to further our understanding of CP violation ...

... and, perhaps, to find new phyiscs!



# $B_{_S} \to J/\psi \phi \ formalism$

Differential decay rate	:	$\frac{d^4\Gamma(\mathbf{B}^0_{\mathrm{s}}\to \mathbf{J}/\psi\phi)}{dt\ d\cos\theta\ d\varphi\ d\cos\psi} \equiv \frac{d^4\Gamma}{dt\ d\Omega} \propto \sum_{k=1}^6 h_k(t) f_k(\Omega)$			
		h. (t)	Bs	Bs f (0, the co)	
$A_0(0) \rightarrow CP$ even $A_{\parallel}(0) \rightarrow CP$ even $A_{\perp}(0) \rightarrow CP$ odd	$\frac{\kappa}{2}$	$\frac{ A_{0}(t) ^{2}}{ A_{  }(t) ^{2}}$	$rac{ ar{A}_k(t) ^2}{ ar{A}_{  }(t) ^2}$	$\frac{f_k(\theta,\psi,\varphi)}{2\cos^2\psi(1-\sin^2\theta\cos^2\varphi)}$ $\frac{\sin^2\psi(1-\sin^2\theta\sin^2\varphi)}{\sin^2\psi(1-\sin^2\theta\sin^2\varphi)}$	
	3 4	$egin{array}{c}  A_{\perp}(t) ^2 \ \Im\{A_{\parallel}^*(t)A_{\perp}(t)\} \end{array}$	$\frac{ \bar{A}_{\perp}(t) ^2}{\Im\{\bar{A}_{  }^*(t)\bar{A}_{\perp}(t)\}}$	$\sin^2\psi\sin^2 heta \ -\sin^2\psi\sin2 heta\sinarphi$	
	5	$ \Re\{A_0^*(t)A_{  }(t)\} \\ \Im\{A_0^*(t)A_{\perp}(t)\} $	$ \Re\{A_0^*(t)A_{  }(t)\} \\ \Im\{\bar{A}_0^*(t)\bar{A}_{\perp}(t)\} $	$\frac{\frac{1}{\sqrt{2}}\sin 2\psi \sin^2 \theta \sin 2\varphi}{\frac{1}{\sqrt{2}}\sin 2\psi \sin 2\theta \cos \varphi}$	

$$\begin{split} |\bar{A}_{0}(t)|^{2} &= |\bar{A}_{0}(0)|^{2} \mathrm{e}^{-\Gamma_{\mathrm{s}}t} \Big[ \cosh\left(\frac{\Delta\Gamma_{\mathrm{s}}t}{2}\right) - \cos\Phi \sinh\left(\frac{\Delta\Gamma_{\mathrm{s}}t}{2}\right) - \sin\Phi \sin(\Delta m_{\mathrm{s}}t) \Big], \\ |\bar{A}_{\parallel}(t)|^{2} &= |\bar{A}_{\parallel}(0)|^{2} \mathrm{e}^{-\Gamma_{\mathrm{s}}t} \Big[ \cosh\left(\frac{\Delta\Gamma_{\mathrm{s}}t}{2}\right) - \cos\Phi \sinh\left(\frac{\Delta\Gamma_{\mathrm{s}}t}{2}\right) - \sin\Phi \sin(\Delta m_{\mathrm{s}}t) \Big], \\ |\bar{A}_{\perp}(t)|^{2} &= |\bar{A}_{\perp}(0)|^{2} \mathrm{e}^{-\Gamma_{\mathrm{s}}t} \Big[ \cosh\left(\frac{\Delta\Gamma_{\mathrm{s}}t}{2}\right) + \cos\Phi \sinh\left(\frac{\Delta\Gamma_{\mathrm{s}}t}{2}\right) + \sin\Phi \sin(\Delta m_{\mathrm{s}}t) \Big], \\ \Im\{\bar{A}_{\parallel}^{*}(t)\bar{A}_{\perp}(t)\} &= |\bar{A}_{\parallel}(0)||\bar{A}_{\perp}(0)|\mathrm{e}^{-\Gamma_{\mathrm{s}}t} \Big[ -\cos(\delta_{\perp} - \delta_{\parallel})\sin\Phi \sinh\left(\frac{\Delta\Gamma_{\mathrm{s}}t}{2}\right) \\ &- \sin(\delta_{\perp} - \delta_{\parallel})\cos(\Delta m_{\mathrm{s}}t) + \cos(\delta_{\perp} - \delta_{\parallel})\cos\Phi \sin(\Delta m_{\mathrm{s}}t) \Big], \\ \Re\{\bar{A}_{0}^{*}(t)\bar{A}_{\parallel}(t)\} &= |\bar{A}_{0}(0)||\bar{A}_{\parallel}(0)|\mathrm{e}^{-\Gamma_{\mathrm{s}}t}\cos\delta_{\parallel} \Big[ \cosh\left(\frac{\Delta\Gamma_{\mathrm{s}}t}{2}\right) - \cos\Phi\sinh\left(\frac{\Delta\Gamma_{\mathrm{s}}t}{2}\right) \\ &- \sin\Phi\sin(\Delta m_{\mathrm{s}}t) \Big] and \\ \Im\{\bar{A}_{0}^{*}(t)\bar{A}_{\perp}(t)\} &= |\bar{A}_{0}(0)||\bar{A}_{\perp}(0)|\mathrm{e}^{-\Gamma_{\mathrm{s}}t} \Big[ -\cos\delta_{\perp}\sin\Phi\sinh\left(\frac{\Delta\Gamma_{\mathrm{s}}t}{2}\right) \\ &- \sin\delta_{\perp}\cos(\Delta m_{\mathrm{s}}t) \Big] and \\ \Im\{\bar{A}_{0}^{*}(t)\bar{A}_{\perp}(t)\} &= |\bar{A}_{0}(0)||\bar{A}_{\perp}(0)|\mathrm{e}^{-\Gamma_{\mathrm{s}}t} \Big[ -\cos\delta_{\perp}\sin\Phi\sinh\left(\frac{\Delta\Gamma_{\mathrm{s}}t}{2}\right) \\ &- \sin\delta_{\perp}\cos(\Delta m_{\mathrm{s}}t) \Big] + \cos\delta_{\perp}\cos\Phi\sin(\Delta m_{\mathrm{s}}t) \Big]. \end{split}$$

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 $\pm$  signs differ for  $B_{_{\rm S}}$  and  $\overline{\rm B}_{_{\rm S}}$ 

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