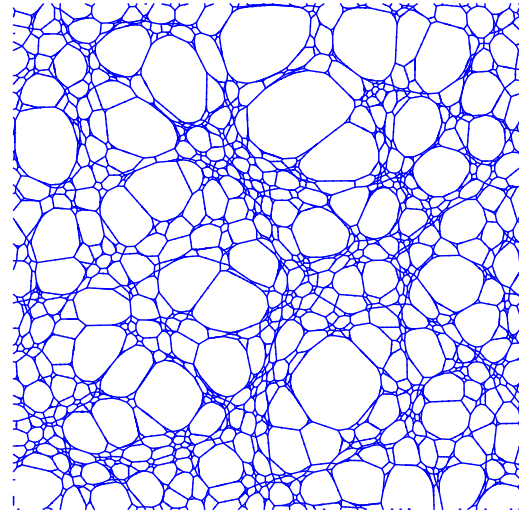


# Micromechanical modelling of near-ideal polymer networks

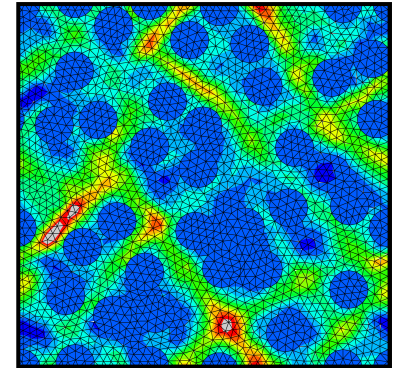
Laurence Brassart

18 May 2020



# Research interests

- Continuum mechanics
- Constitutive modelling
- Micromechanics, scale transition methods
- Chemo-mechanical couplings in materials



Materials:

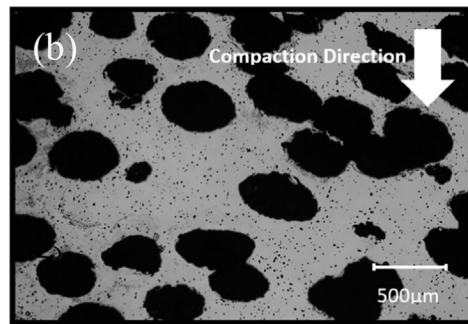
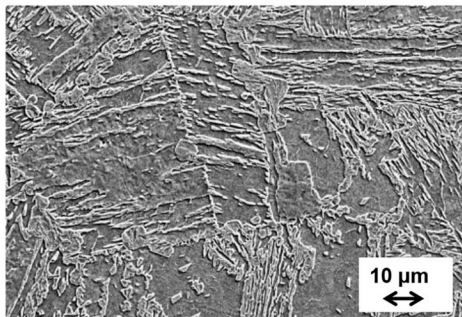
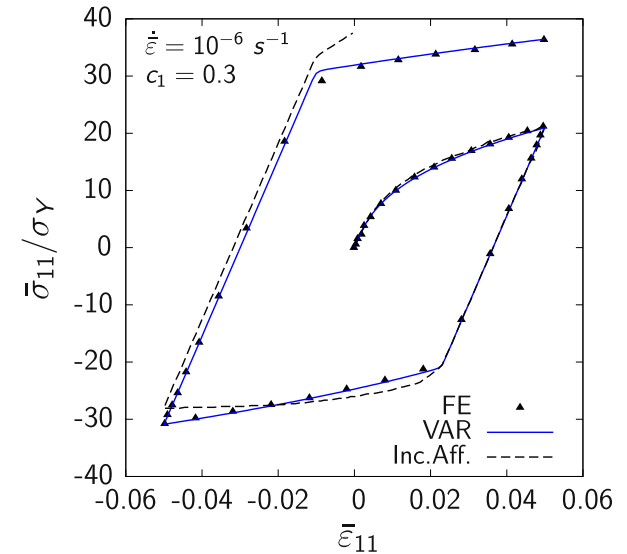
Composite materials; Polymers; Energy materials; Soft materials;  
Biomaterials

Collaborations:

Université catholique de Louvain, Harvard University, Monash  
University, Ecole Centrale de Nantes

# Micromechanics of composites

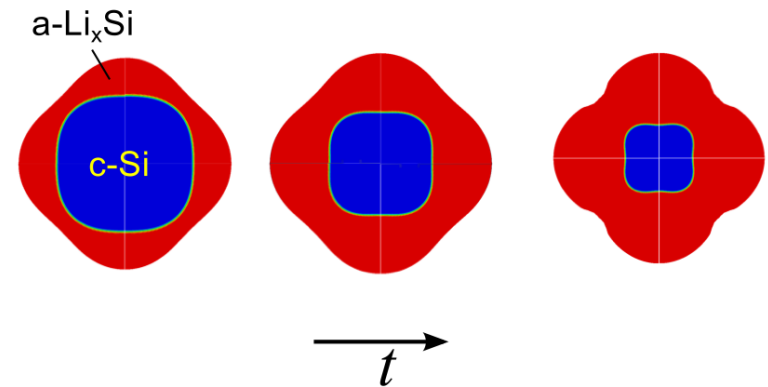
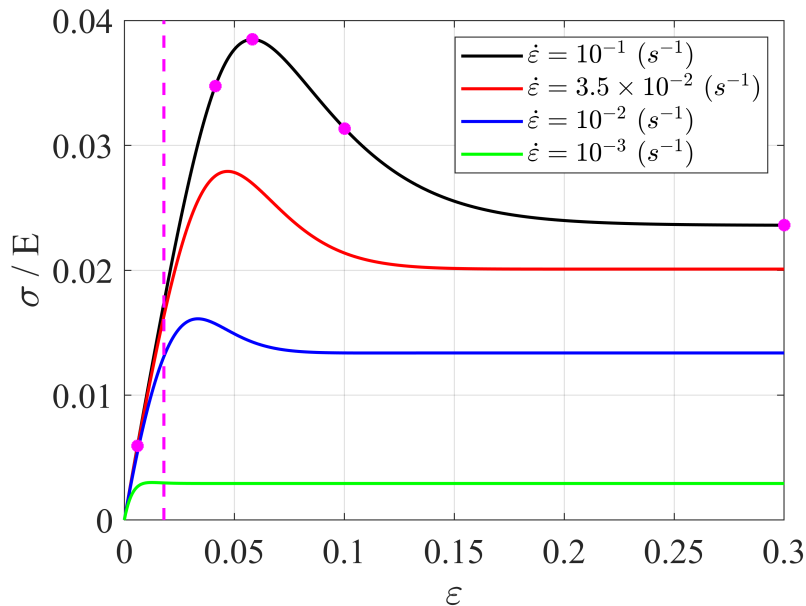
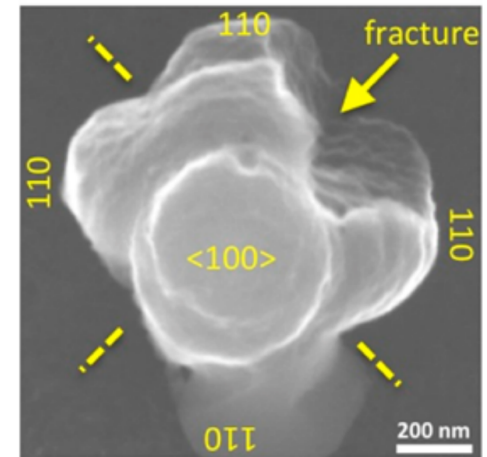
- Mean-field modelling in elasto-viscoplasticity
- FEA at micro and macro scales
- Applications to fibre-reinforced composites, dual-phase steels, porous materials, 3D-printed architected materials



Brassart et al., *IJP* (2012); Pierman et al, *Acta Mat*, (2014); Soro et al., *Mater. Sci. Eng. A* (2018), Ismail et al, *IJP* (2019);

# Mechanics of Li-ion batteries

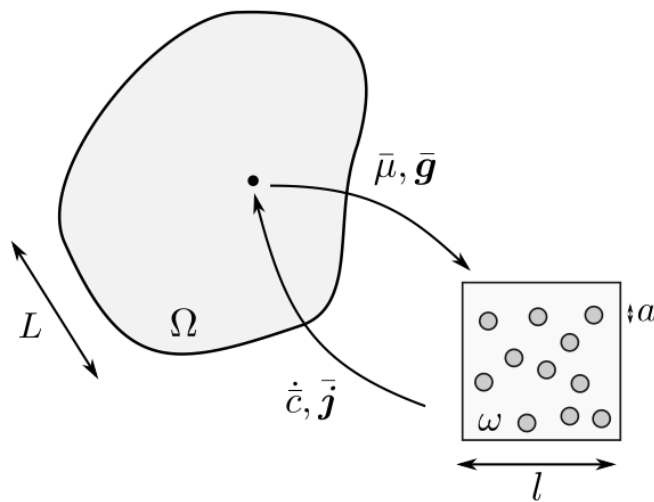
- Constitutive modelling of a-LiSi
- Coupled diffusion-plasticity analysis
- Anisotropic swelling and fracture in c-Si



Brassart and Suo, *JMPS* (2013); Brassart et al., *IJSS* (2013); Sandu et al., *ACS Nano*, (2014); Bagheri and Brassart, *in preparation*.

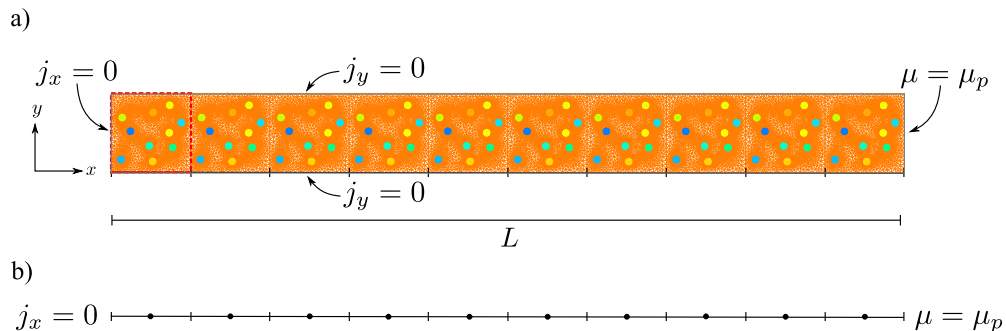
# Mean-field homogenisation for transient diffusion problems

- Diffusion in heterogeneous media with high diffusivity contrast
- Non-classical effective behaviour



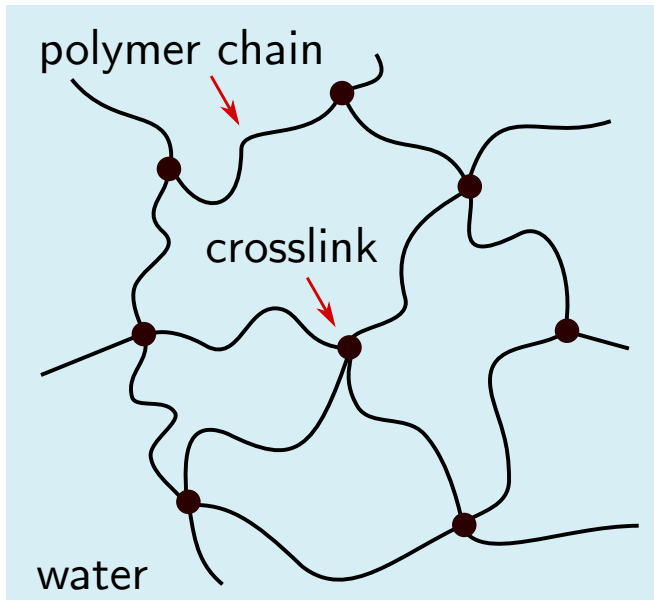
$$\bar{\mu} \dot{\bar{c}} - \bar{\mathbf{j}} \cdot \bar{\mathbf{g}} = -\frac{1}{V} \int_{\partial\omega} \mu \mathbf{j} \cdot \mathbf{n} dS$$

$$= \frac{1}{V} \int_{\omega} \mu \dot{c} dV - \frac{1}{V} \int_{\omega} \mathbf{j} \cdot \nabla \mu dV$$

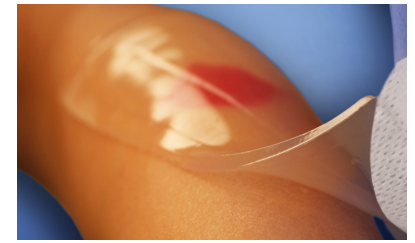


# Hydrogels in everyday life

Hydrogels are crosslinked polymer networks swollen in water



Contact lenses



Wound dressing



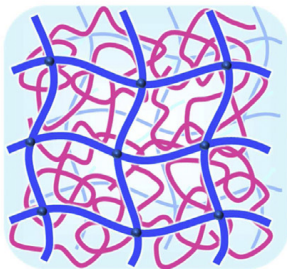
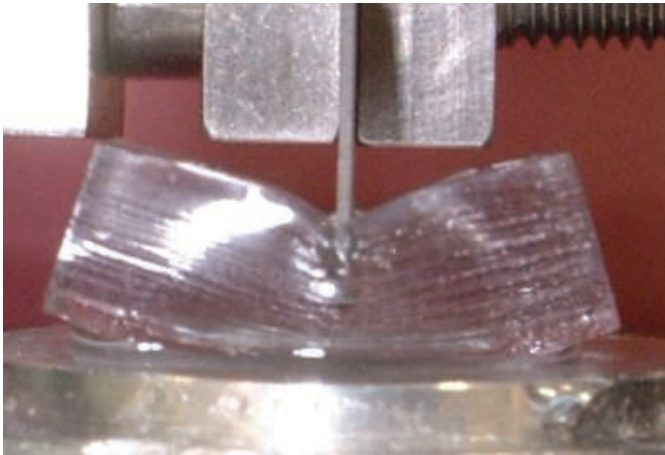
Superabsorbent  
diapers



Water beads

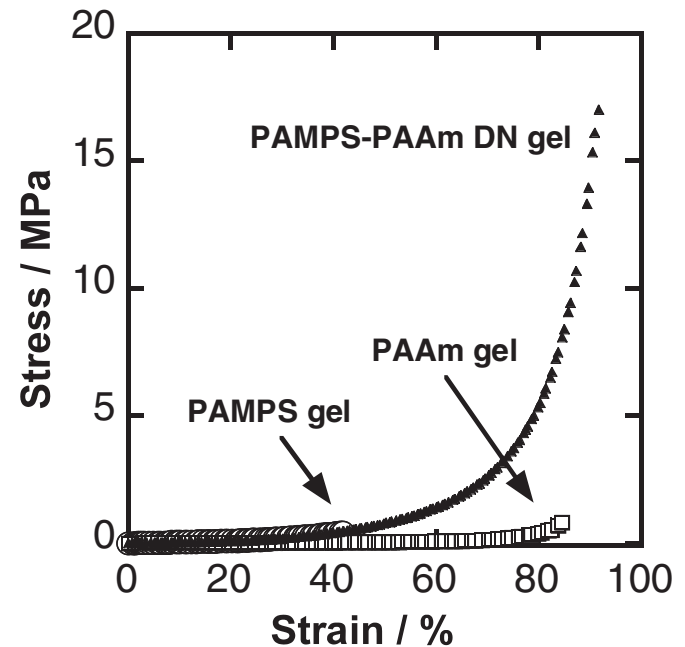
# Network design for strength and toughness

Example: Double-network gels



*1<sup>st</sup> network:*  
Rigid skeleton, dilute,  
short chains

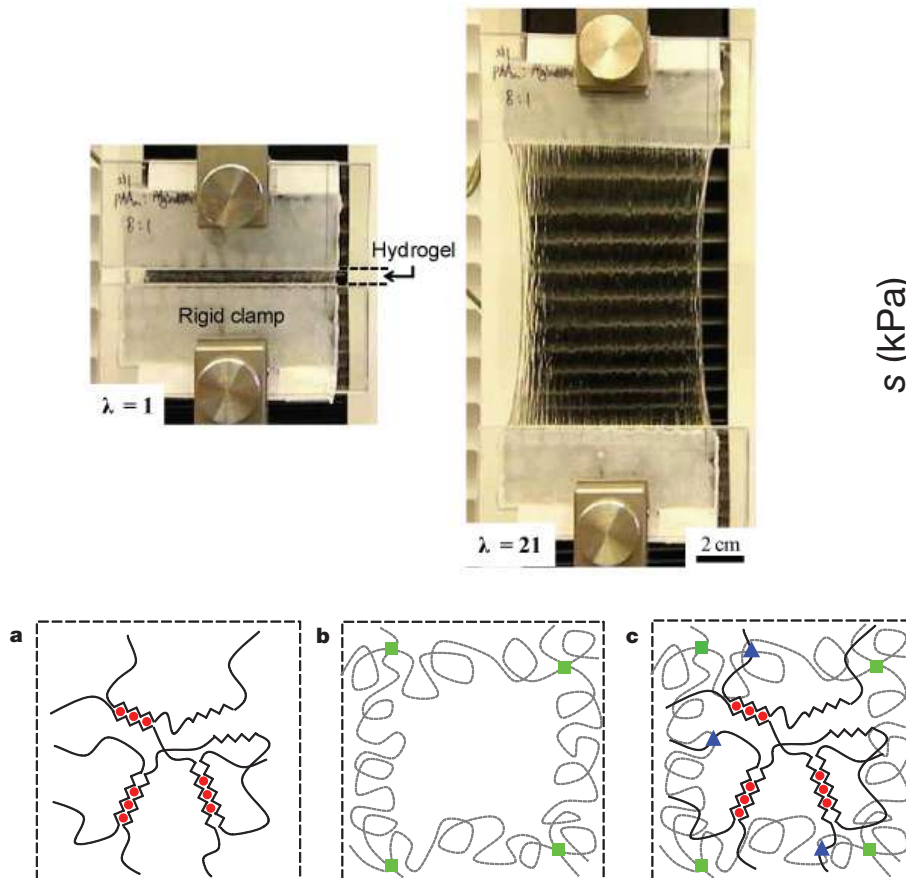
*2<sup>nd</sup> network:*  
Flexible matrix,  
concentrated, long chains



$$\Gamma = 1000 \text{ J/m}^2$$

# Network design for strength and toughness

Example: hybrid ionic-covalent gels



$$\Gamma = 8700 \text{ J/m}^2$$

$$\lambda_{\max} > 20$$



# Continuum modelling

Kinematics:

$$\boldsymbol{x} = \boldsymbol{\chi}(\boldsymbol{X}, t)$$

$$\boldsymbol{F} = \nabla \boldsymbol{\chi}$$

$$\det(\boldsymbol{F}) = 1 + \Omega C$$

Conservation equations:

$$\nabla \cdot \boldsymbol{P} = \mathbf{0}$$

$$\frac{\partial C}{\partial t} = -\nabla \cdot \boldsymbol{J}$$

Constitutive models:

$$\mu \dot{C} + \boldsymbol{P} : \dot{\boldsymbol{F}} - \boldsymbol{J} \cdot \nabla \mu - \dot{W} \geq 0$$

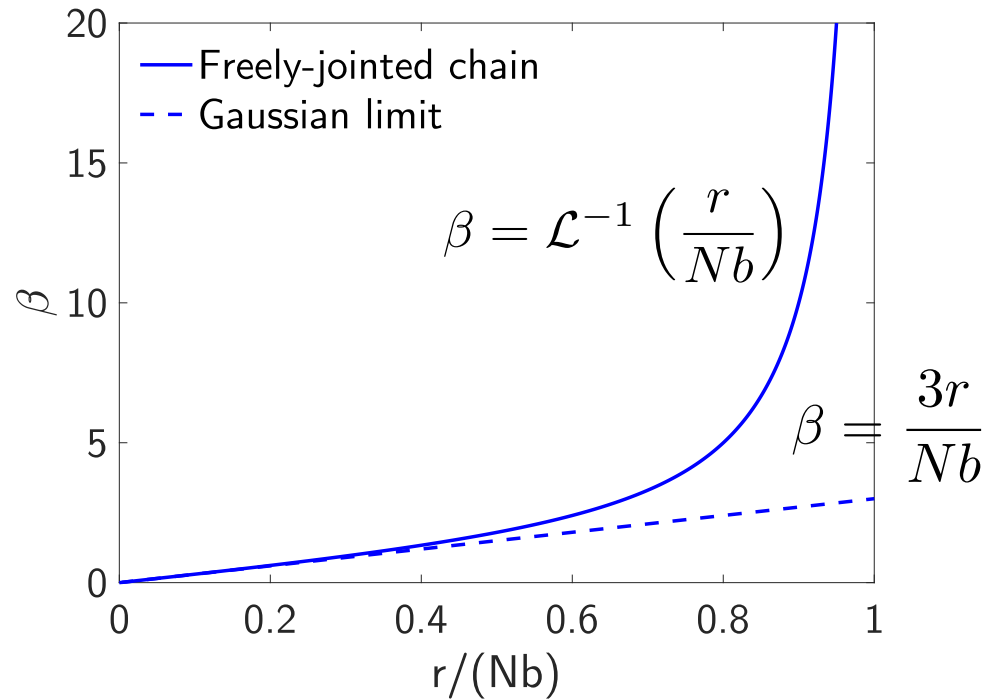
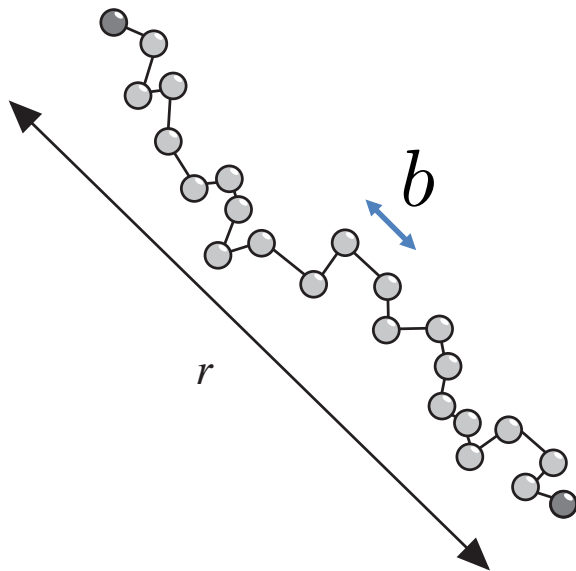
$$W = W^e(\boldsymbol{F}) + W^c(C)$$

+ kinetic relations

← Micromechanical  
modelling

# Single chain behaviour

Freely-jointed chain with  $N$  Kuhn segments of length  $b$



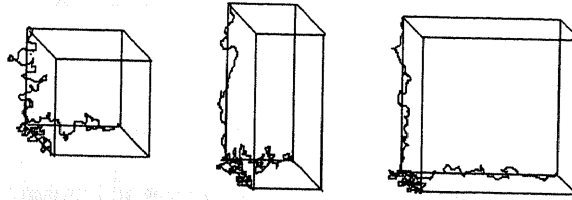
$$w = NkT \left( \frac{\beta}{\tanh \beta} + \log \frac{\beta}{\sinh \beta} \right)$$

$$\beta = \frac{fb}{kT}$$

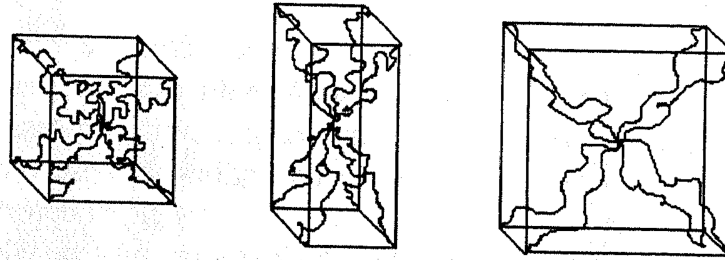
# Network models

Assembly of representative chains deforming affinely with the macroscopic strain

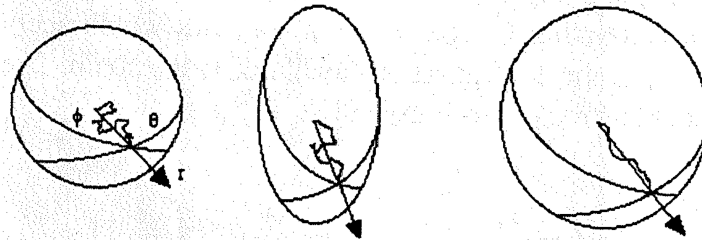
- 3-chain model



- 8-chain model

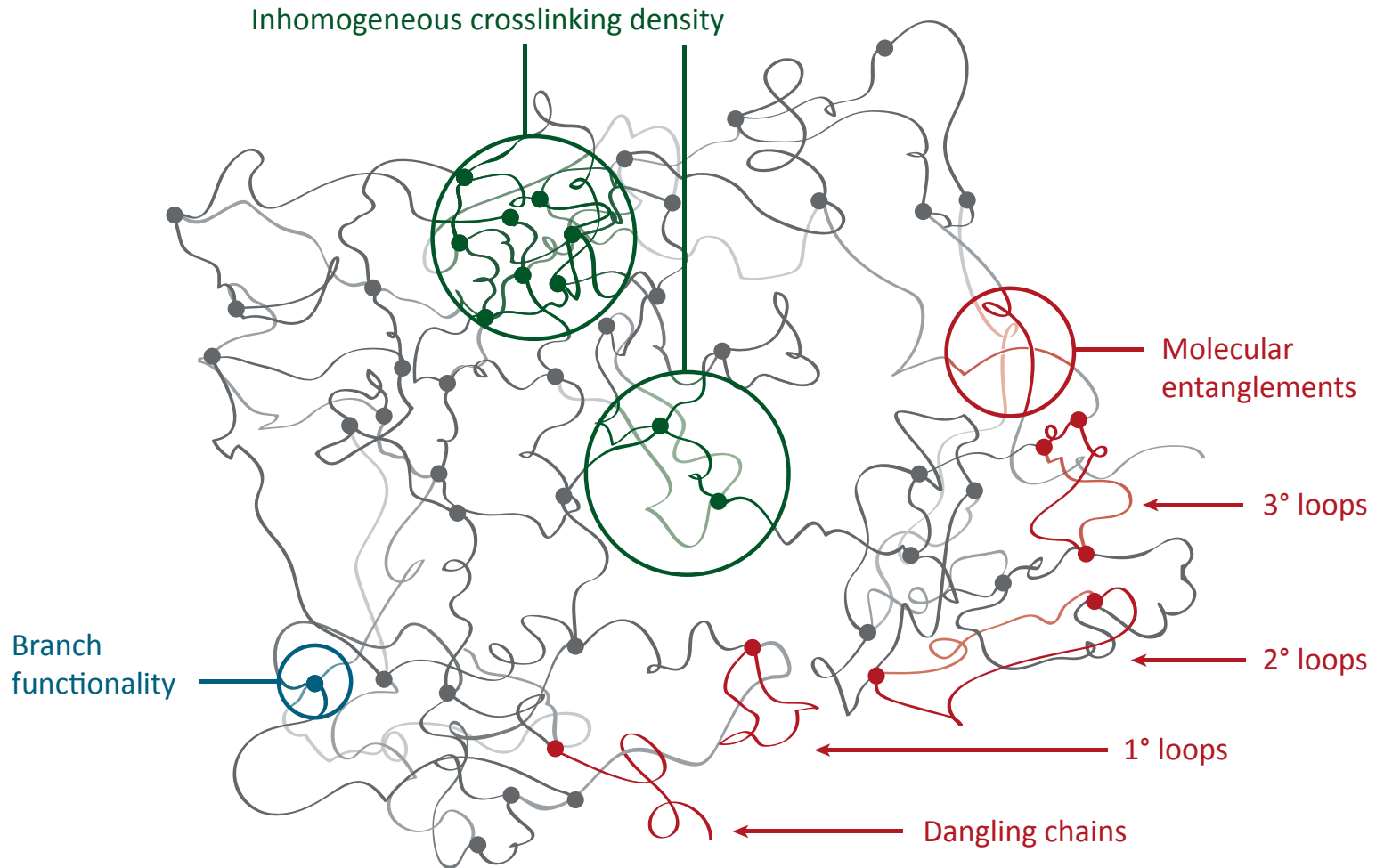


- Full-network model



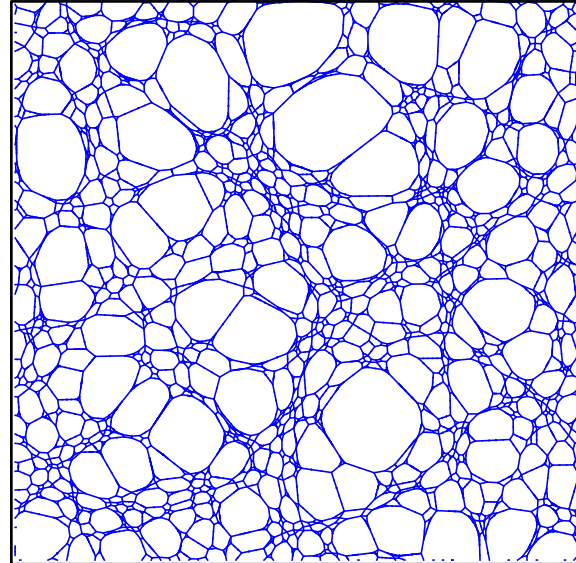
$$W^e = \frac{1}{V} \sum_n w$$

# Real networks are imperfect



# Micromechanical modelling

- Random assembly of springs (polymer chains) connecting at junction points (crosslinks)
- Spring behaviour described by entropic force-extension relation
- Includes topological defects



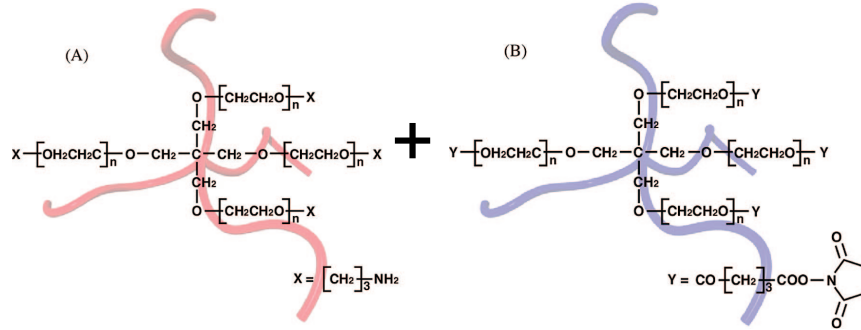
## Objectives

- Understand the relative contributions of network parameters to mechanical properties
- Generate reference results to validate constitutive models

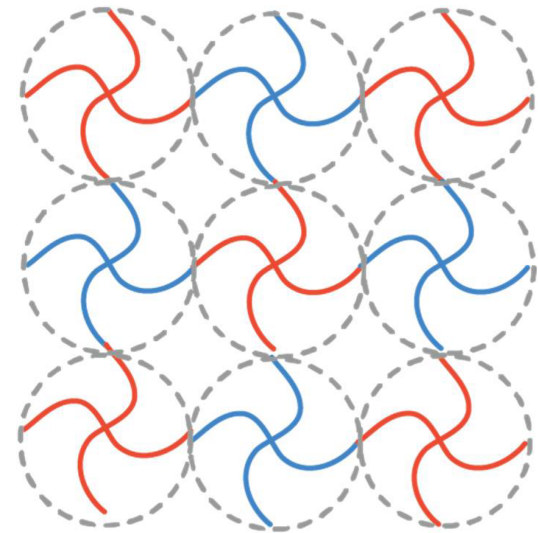
# Model system: near-ideal networks

## 4-arm PEG hydrogels

### Polymer precursors



### Gelation in solution



### Ideal-network gel

- Controlled arm length
- High conversion rate
- No first-order loops
- Few entanglements

# Network structures

## Input:

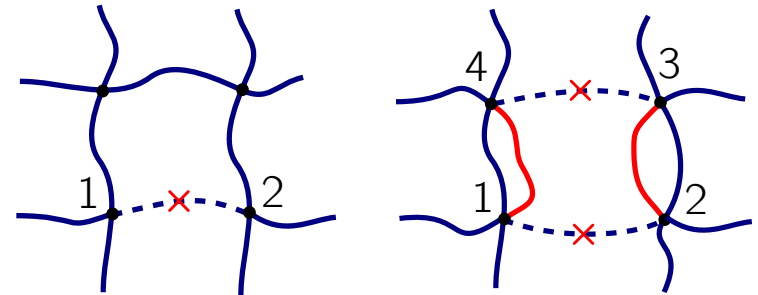
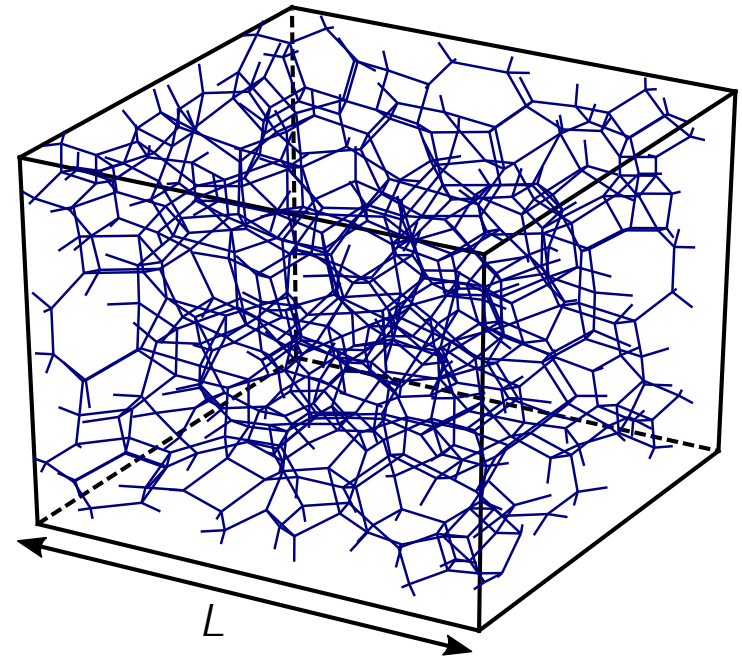
- Chain density  $\nu = n/L^3$
- Conversion probability  $p$   
→ average coordination  $\bar{Z}$
- Loop fraction  $\zeta$

$$\mathbf{x}_\alpha = \mathbf{F} \cdot \mathbf{X}_\alpha$$

$$\mathbf{P} = \frac{1}{V} \sum_{\alpha} (f_{\alpha}^e \otimes \mathbf{X}_{\alpha})$$

## Output:

- Stress-strain curves
- Chain length and orientation distribution

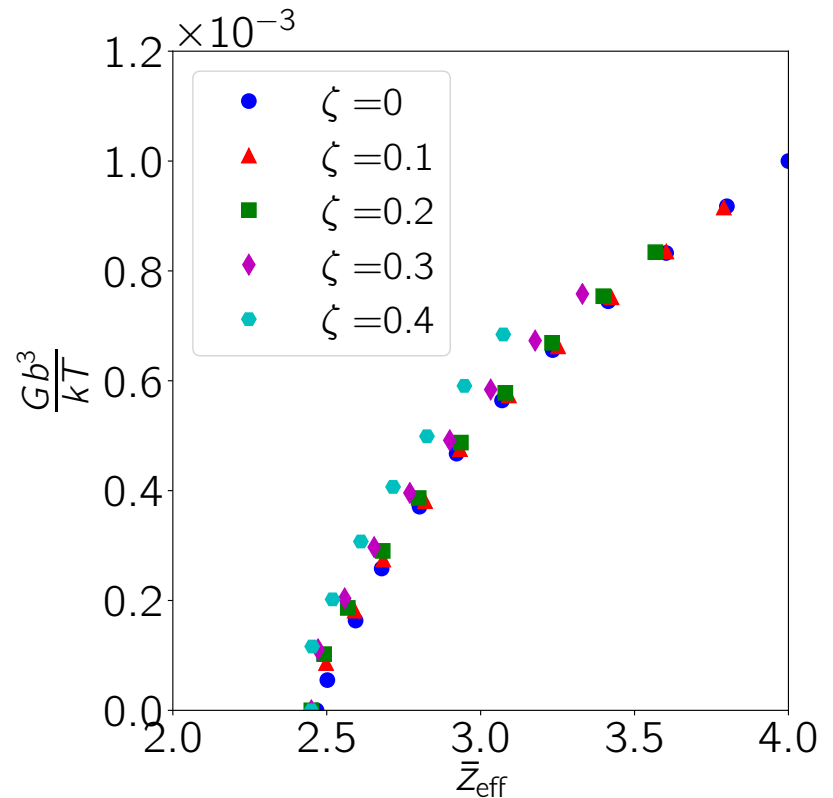
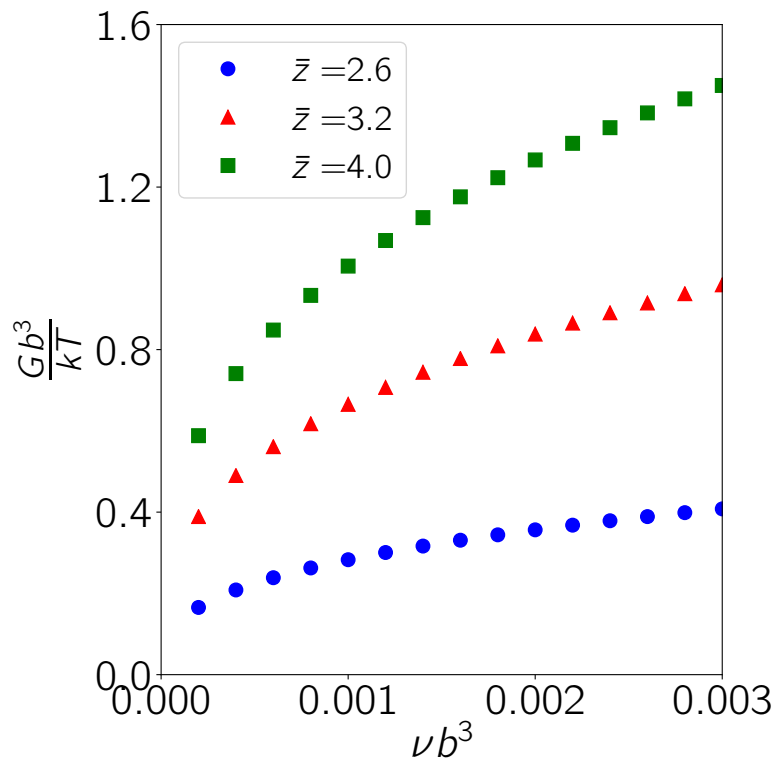


Energy minimisation (LAMMPS)

$$\bar{Z}_{\text{eff}} \approx \bar{Z} \left(1 - \frac{\zeta}{2}\right)$$

# Scaling of the shear modulus

$$G \sim \frac{1}{N} \nu^{1/3} (\bar{Z}_{\text{eff}} - \bar{Z}_{\text{eff},c})^{2/3}$$





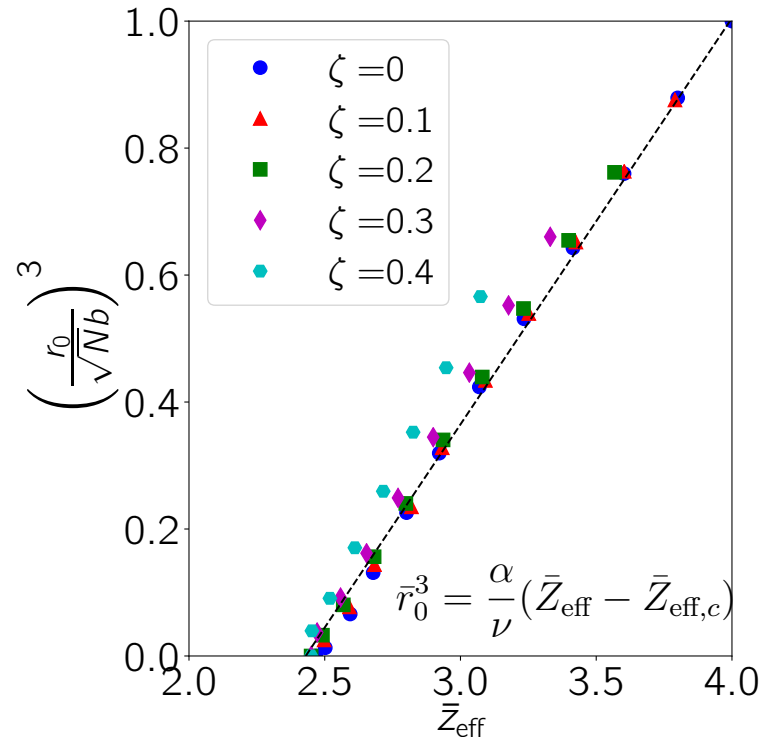
# Interpretation based on chain pre-stretch

Elastic modulus of a network of Gaussian chains:

$$G = kT\nu \frac{\bar{r}_0^2}{Nb^2} \quad \bar{r}_0^2 = \langle r_0^2 \rangle \quad (\text{exact result})$$

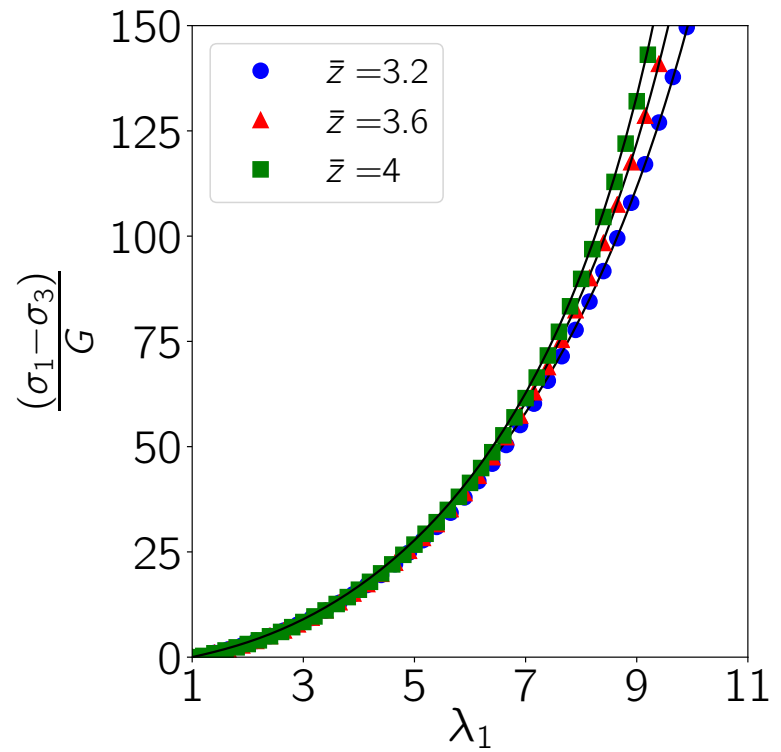
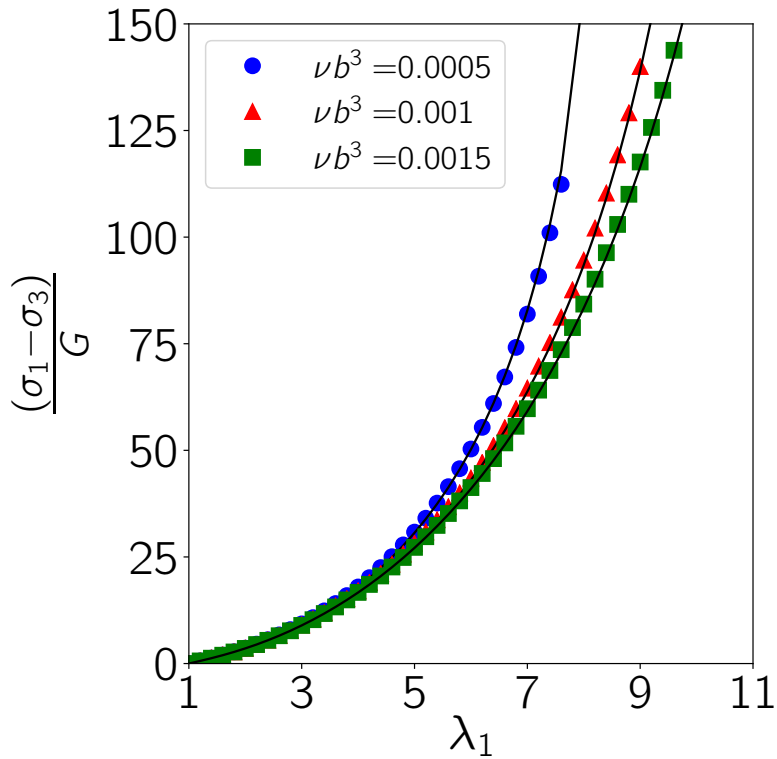
- Topology impacts the modulus via the chain pre-stretch
- Scaling of chain pre-stretch with topology largely explained by geometric arguments
- Coincides with the classical affine estimate only when  $\bar{r}_0 = \sqrt{N}b$ :

$$G = \nu kT$$



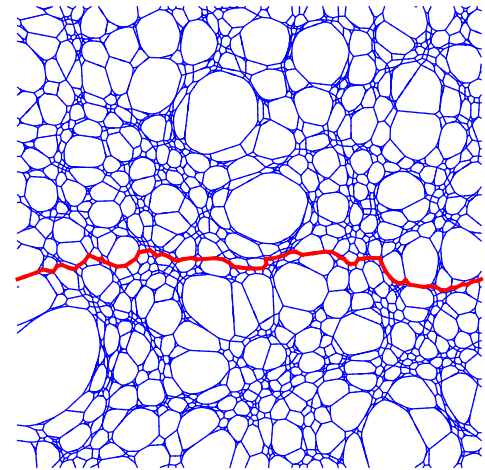
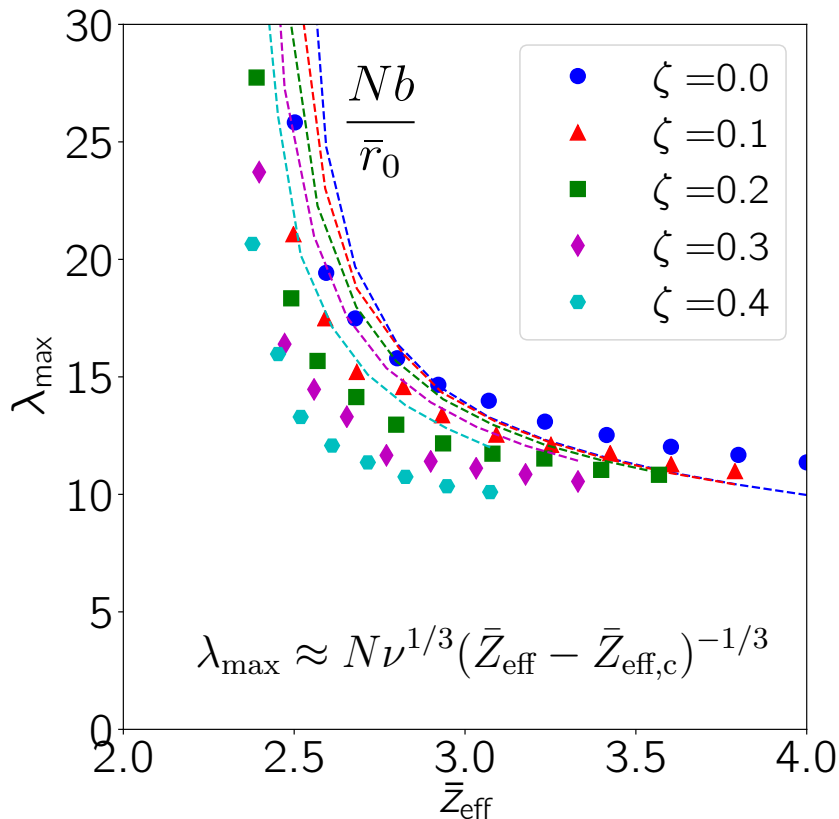
# Large-deformation behaviour

Stiffening rate depends on density and topology via the pre-stretch



# Limit extensibility

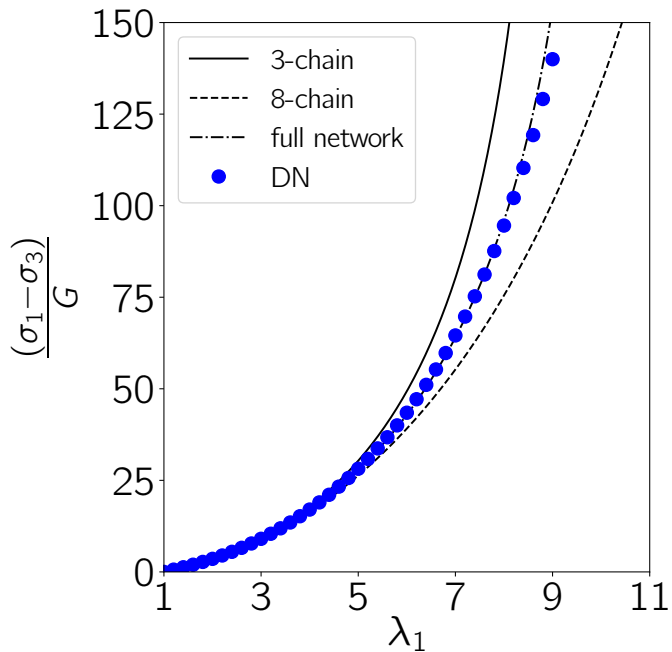
- Limit extensibility partly explained by the pre-stretch
- Loops reduce the shortest chain path



$$\lambda_{\max} = \frac{n^{SP} Nb}{L}$$

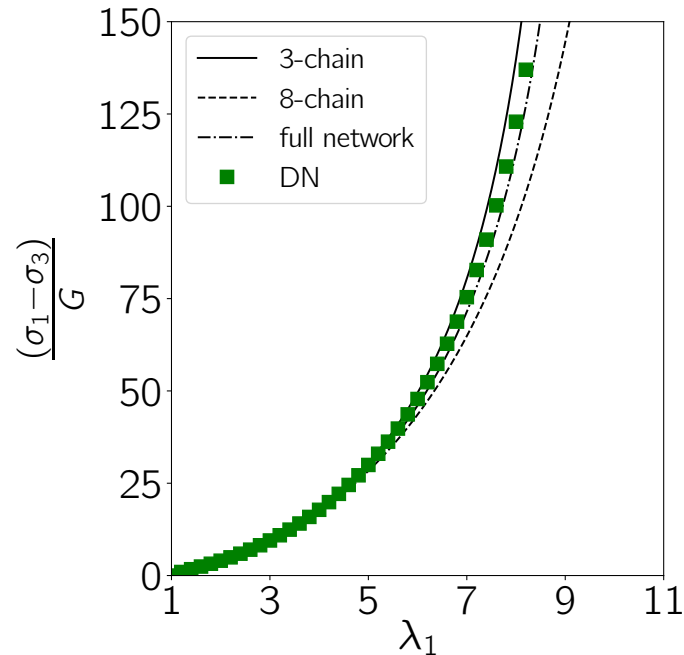
# Comparison to analytical models

- Overall, the full-network is the most accurate
- The 8-chain model consistently underestimates the response



Uniaxial extension

$$\lambda_1 \geq 1 \quad \lambda_2 = \lambda_3 = \frac{1}{\sqrt{\lambda_1}}$$



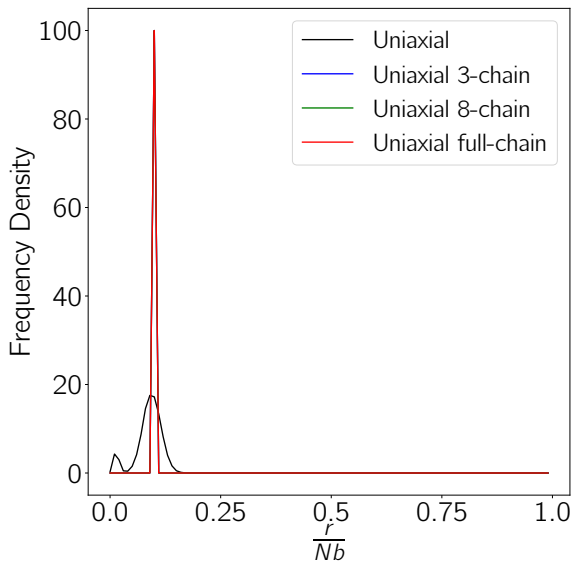
Biaxial extension

$$\lambda_1 = \lambda_2 \geq 1, \quad \lambda_3 = \frac{1}{\lambda_1^2}$$

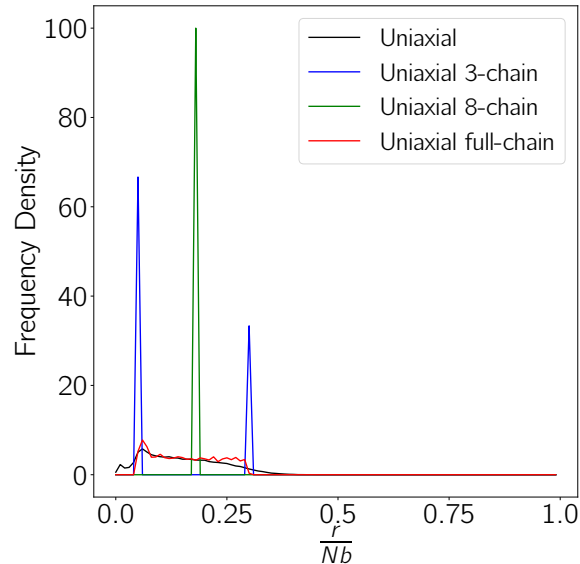
# Chain length distribution: Uniaxial extension

The full-network model well captures the chain length distribution

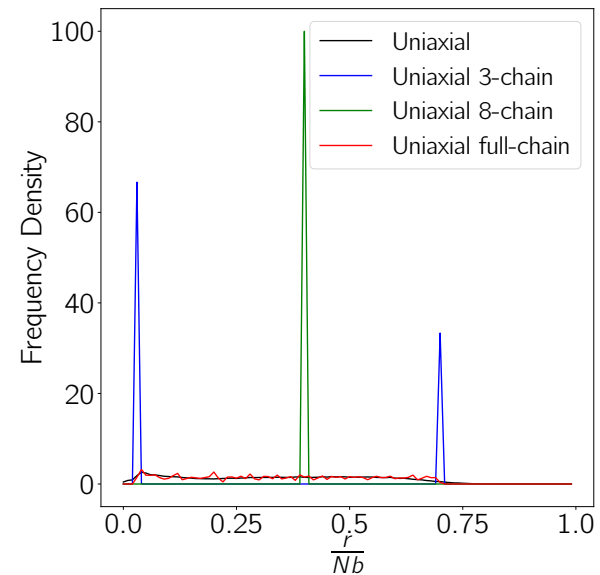
$\lambda_1 = 1$



$\lambda_1 = 3$

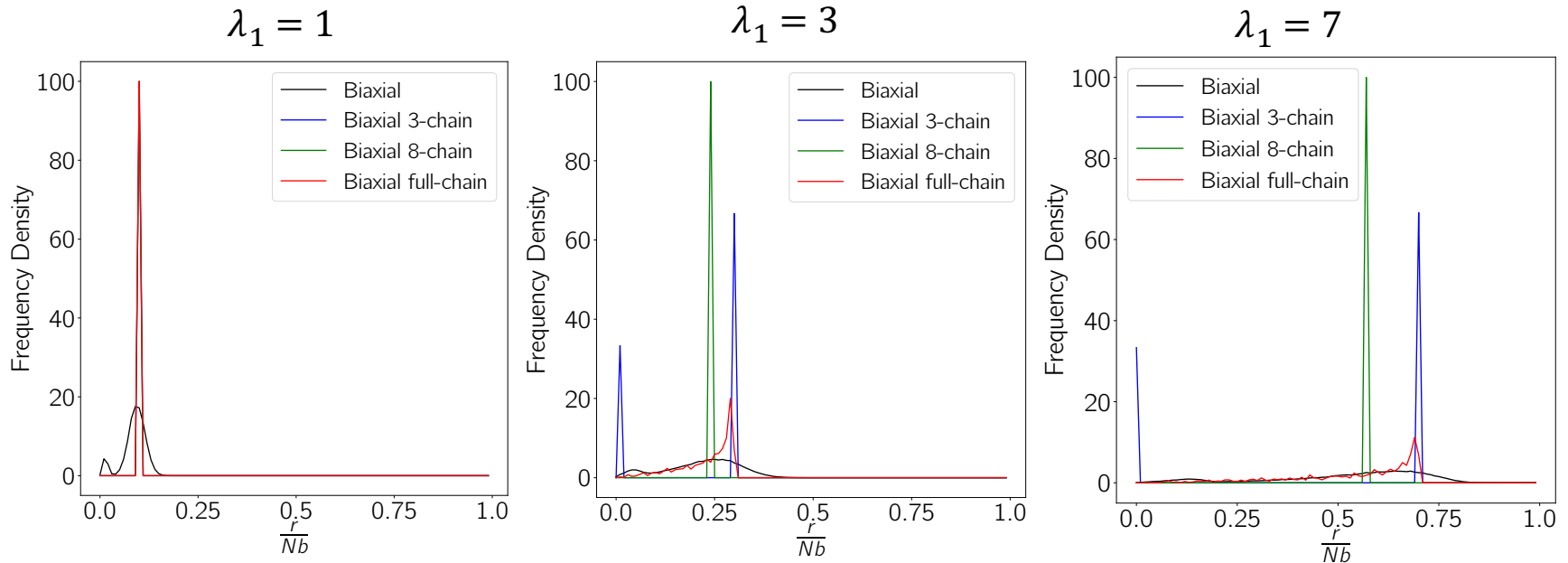


$\lambda_1 = 7$



# Chain length distribution: Biaxial extension

The full-network misses out a fraction of the highly extended chains.



# Application to PEG hydrogels



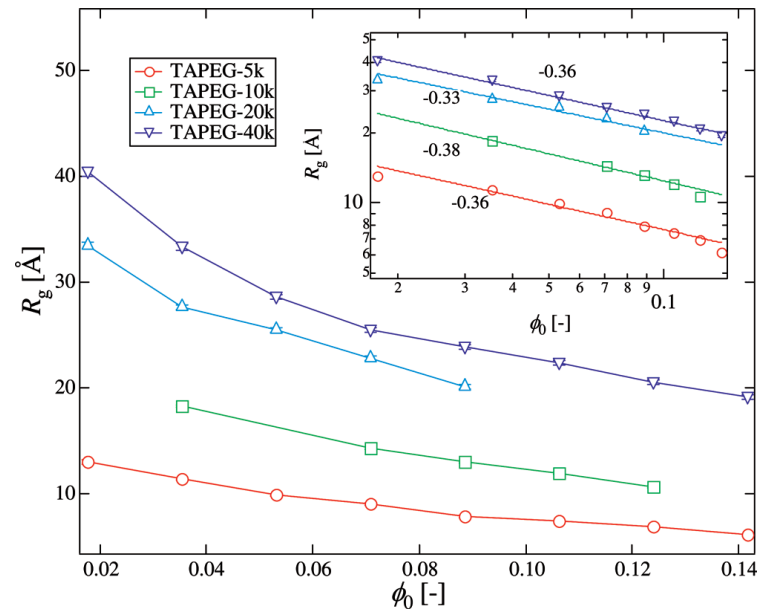
- Kuhn length:  $b = 1.1$  nm
- Contour length:

$$L_c = n_{mon} \times 0.36 \text{ nm} = Nb$$

- Modified FJC to account for **non-ideal chain behaviour**

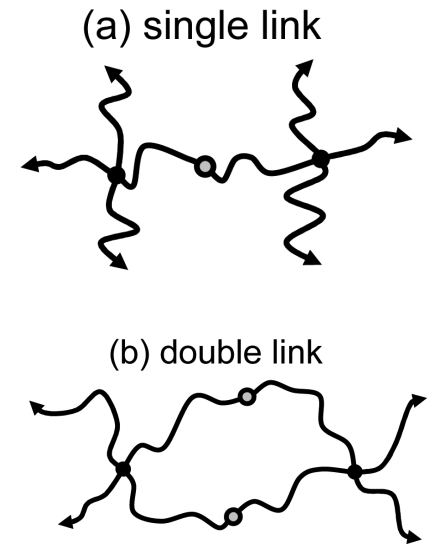
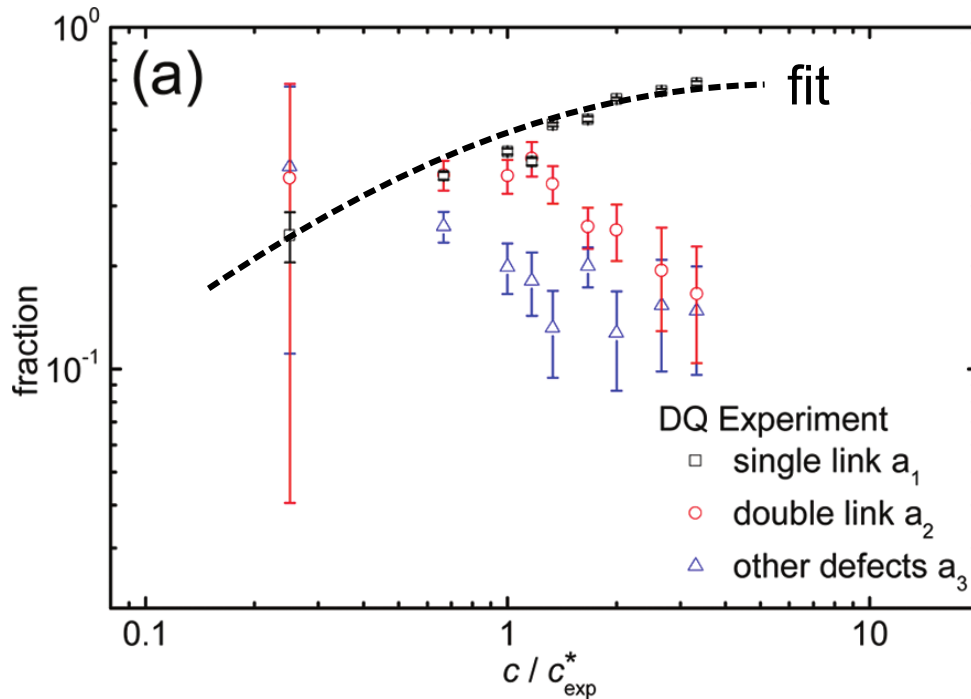
$$\frac{f\bar{b}}{kT} = \mathcal{L}^{-1} \left( \frac{r}{N\bar{b}} \right) \quad \bar{b} = \frac{R_F^2}{Nb} \quad \bar{N} = \frac{(Nb)^2}{R_F^2}$$

Matsunaga et al., *Macromol.* (2009)



$$R_F \approx v^{1/5} b^{3/5} N^{3/5}$$

# Quantifying Loops in 4-arm PEG gels

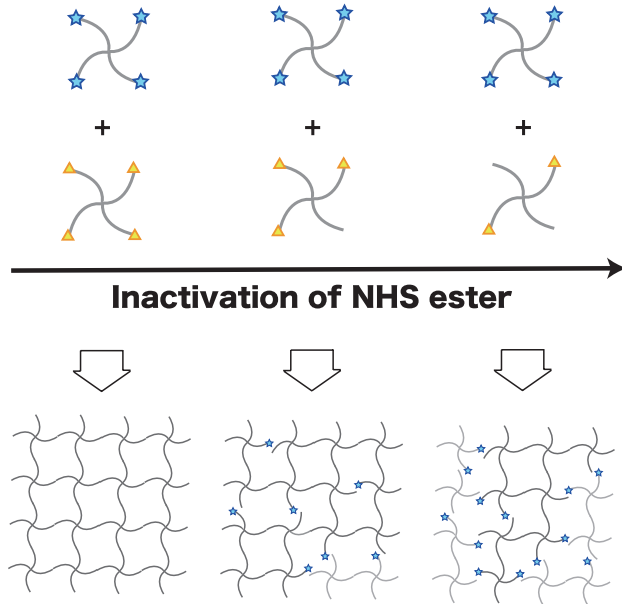


Lange et al., *Macromol.* (2011)

The fraction of loops increases as the pre-polymer volume fraction decreases

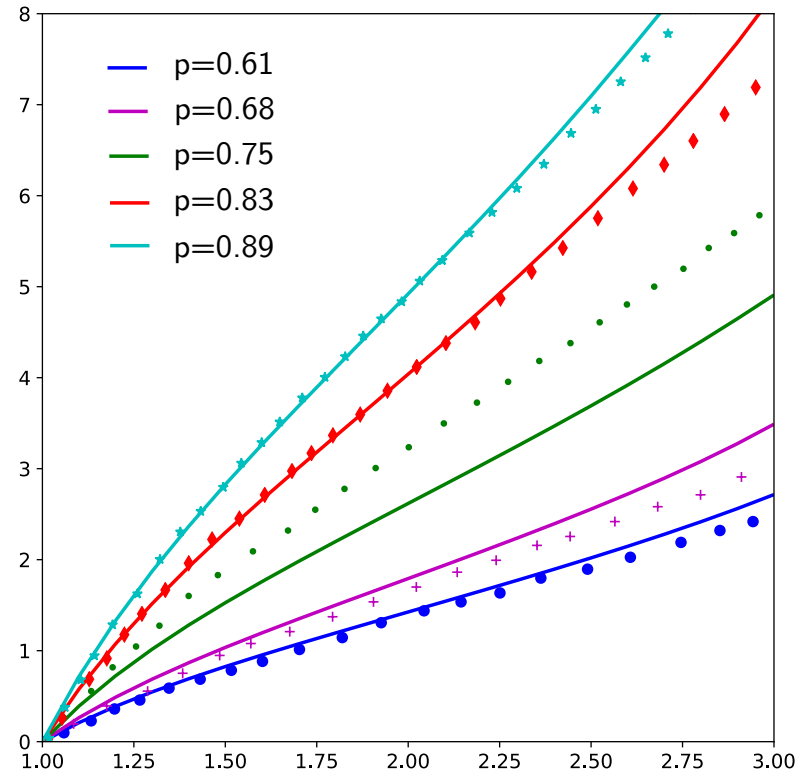


# 4-arm PEG with tunable connectivity probability



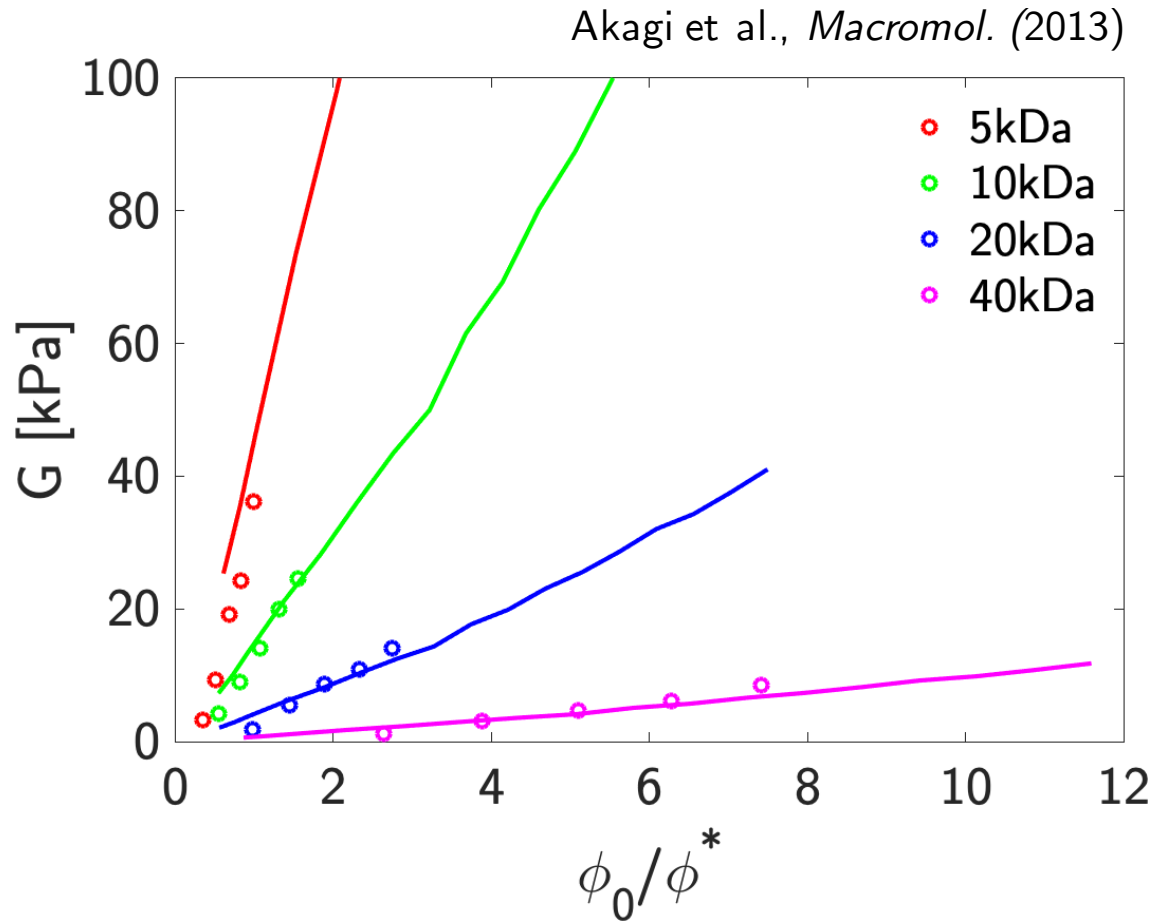
Nishi et al., *J. Chem. Phys.* (2012)

Akagi et al., *RSC Adv.* (2013)



Good prediction with only one fitting parameter  $\nu^*$

# Dependence of modulus on concentration



Modulus dependence on molecular weight results primarily from the non-ideality of chain behaviour

# Summary

- Discrete network models as a tool to investigate the role of network parameters on the mechanical response of hydrogels
- Network defects (dangling ends, loops) have a significant impact on mechanical properties
- Coupling between strain pre-stretch and topology gives scaling relations different from classical theories
- The model can explain experimental trends by considering non-ideal chain behaviour and the presence of network defects

## References:

G. Alamé and L. Brassart, *Soft Matter* 15, 5703 (2019)

G. Alamé and L. Brassart, *About to be submitted*.

# Future work

Extend the computational framework to describe:

- Chain scission
- Crosslink breaking and reforming
- Distribution of chain length
- Interpenetrating networks

