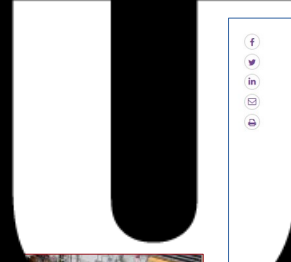


First result from FNAL

Muon g-2



Gavin Hesketh, 12th May 2021

Le Monde

ACTUALITÉS · ÉCONOMIE · VIDÉOS · OPINIONS · CULTURE · SCIENCES · PHYSIQUE

Les rotations du muon électrisent la communauté des physiciens

Une anomalie dans le comportement magnétique du muon, une particule élémentaire, a été observée. Le signe que le cadre conceptuel servant à décrire l'infiniment petit est en train de se fracturer. Des résultats contradictoires relancent le débat.

Par David Larousse

Publié aujourd'hui à 17h00, mis à jour à 18h54 · Lecture 7 min.

Article réservé aux abonnés

The New York Times

L'anomalia del muone: l'esperimento che suggerisce l'esistenza di nuove forze della natura

Un esperimento a Fermilab ha rivelato che il muone, una particella elementare, si comporta in modo diverso da quanto previsto dalla fisica attuale. Questo potrebbe indicare l'esistenza di nuove forze della natura.

DER SPIEGEL

Neue Erkenntnisse in der Teilchenphysik: Kundschafter ins Unbekannte

Seit 50 Jahren erschauen Forscher Einblicke in die Welt jenseits der bekannten Naturgesetze. Mit den Erkenntnissen aus dem Myon-g-2-Experiment könnte sich das Tor zu einer neuen Physik öffnen.

Von Johann Großle

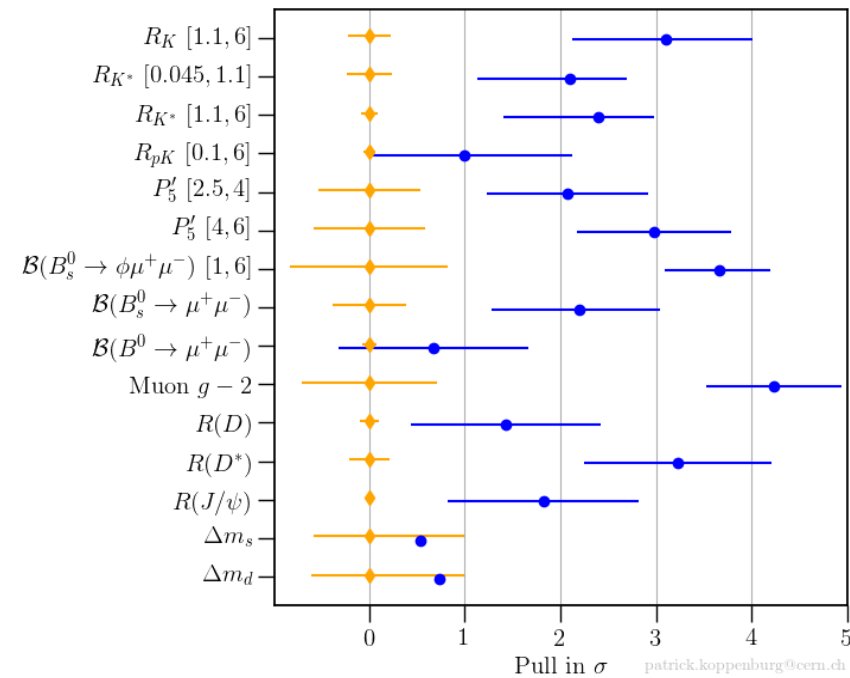
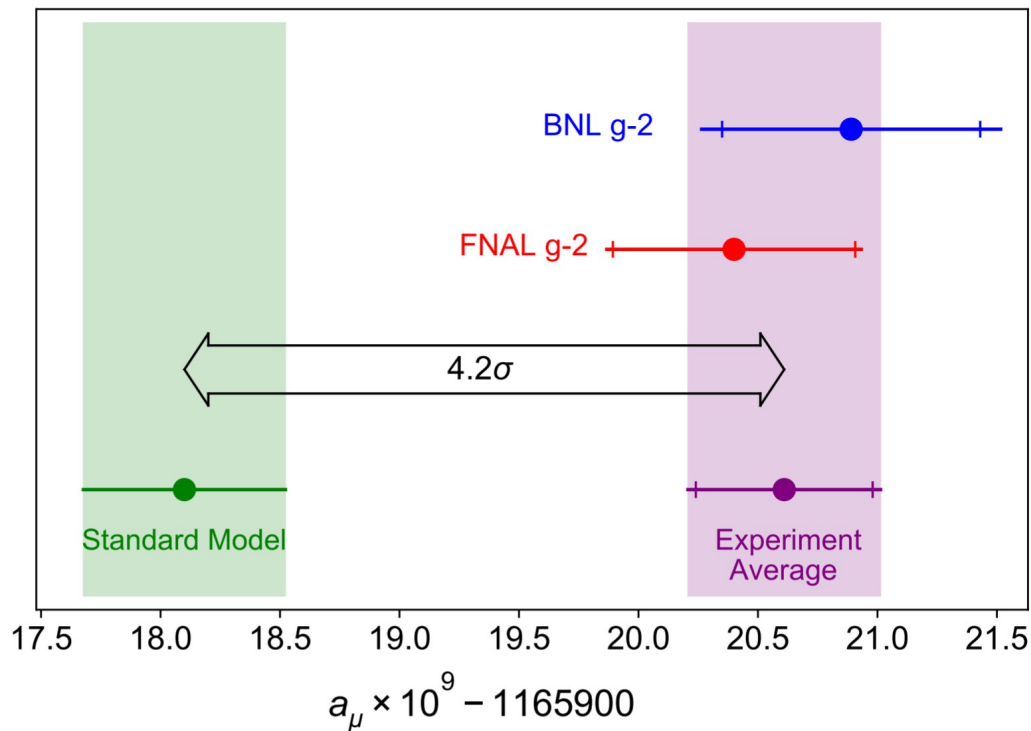
আনন্দবাজার পত্রিকা অনলাইন

ব্রহ্মাণ্ডে রয়েছে আরও অজ্ঞাত কণা, জানাল চেনা কণার আজব আচরণ

মিউন গ-২ পরীক্ষার ফলাফল থেকে জানা গেল যে মৌলিক কণাগুলির আচরণ আমাদের বর্তমান তত্ত্বের বিরুদ্ধে।

Fermilab Muon $g-2$ (E989) confirms the Brookhaven (E821) results

- measured a_μ to 0.46 ppm
- 4.2σ tension with consensus theory prediction



Some background

All fermions have an intrinsic magnetic moment:

- spin + g-factor

$$\vec{\mu} = g \frac{q}{2m} \vec{s}$$

For all fundamental fermions:

$$g = 2 \quad (\text{Dirac, 1927})$$



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$$\vec{\mu} = g \frac{q}{2m} \vec{s}$$

For all fundamental fermions:

$$g = 2 \quad (\text{Dirac, 1927})$$



Electron anomaly!

$$g = 2.00238 \pm 6 \quad (\text{Kusch \& Foley, 1948})$$

All fermions have an intrinsic magnetic moment:

- spin + g-factor

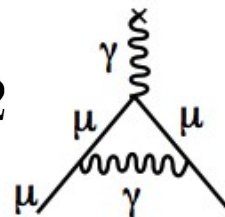
$$\vec{\mu} = g \frac{q}{2m} \vec{s}$$

For all fundamental fermions:

$$g = 2.00232$$

(Schwinger, 1948)

$$\frac{\alpha}{2\pi} = 0.00232$$



Electron anomaly!

$$g = 2.00238 \pm 6 \quad (\text{Kusch \& Foley, 1948})$$

All fermions have an intrinsic magnetic moment:

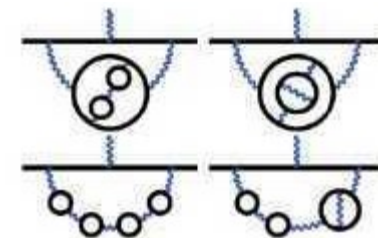
- spin + g-factor

$$\vec{\mu} = g \frac{q}{2m} \vec{s}$$

Electrons today:

$$g = 2.00231930436328 \pm 153$$

Phys. Rev. D **97**, 036001 (2018)



$$g = 2.00231930436146 \pm 56$$

Phys. Rev. A **83**, 052122 (2011)

"the most precisely determined quantity in physics"

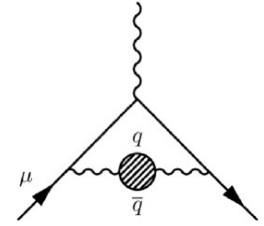
“Anomalous” term: $a = \frac{g - 2}{2}$

- all the loop corrections
- contributions scale as $(m_i/M)^2 \rightarrow$ muons x43,000 more sensitive to higher mass effects

$$a_{\mu}^{\text{SM}} = a_{\mu}^{\text{QED}} + a_{\mu}^{\text{EW}} + a_{\mu}^{\text{HVP}} + a_{\mu}^{\text{HLbL}}$$

$$a_{\mu}^{\text{SM}} = a_{\mu}^{\text{QED}} + a_{\mu}^{\text{EW}} + a_{\mu}^{\text{HVP}} + a_{\mu}^{\text{HLbL}}$$

			a_{μ}^{SM} portion	$\delta a_{\mu}^{\text{SM}}$ portion
QED	<p>1-loop + 2-loop + ...</p>	Perturbative (Known to five-loop)	$\sim 99.99\%$	$\sim 0.001\%$
EW	<p>γ, W, ν_{μ}, Z, H</p>	Perturbative (Known to two-loop)	~ 1 ppm	$\sim 0.2\%$
HVP	<p>had.</p>	Non-perturbative (Data-driven & lattice)	~ 59 ppm	$\sim 84\%$
HLbL	<p>had.</p>	Non-perturbative (Data-driven & lattice)	~ 1 ppm	$\sim 16\%$

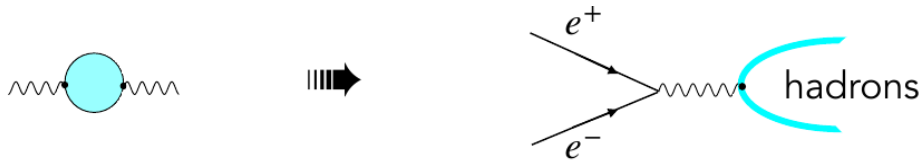


g-2 Theory Initiative determination of SM value

T. Aoyama et al, arXiv:2006.04822, Phys. Repts. 887 (2020) 1-166

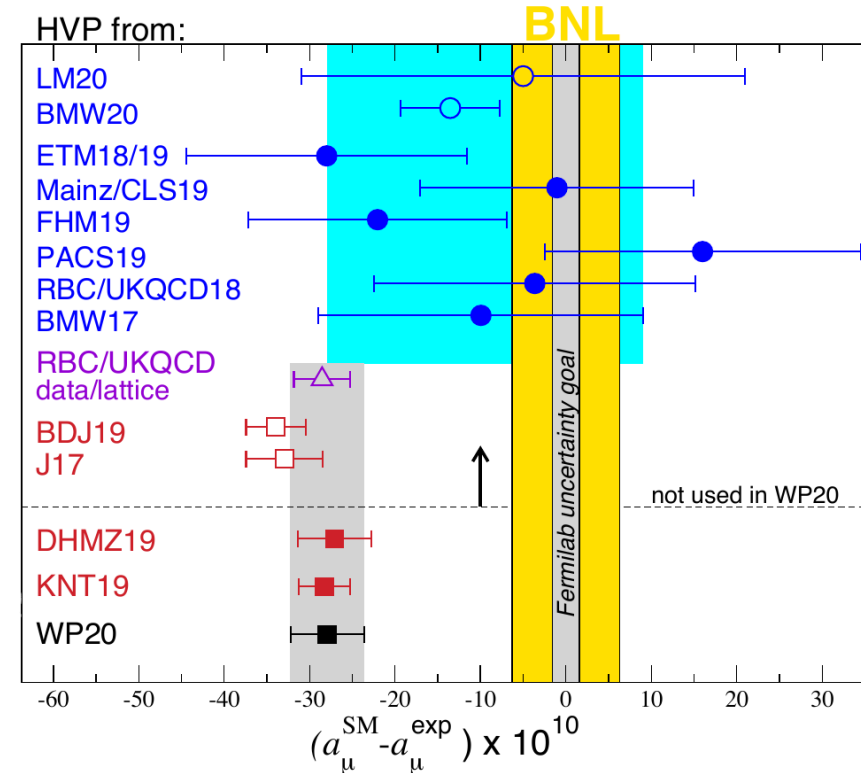
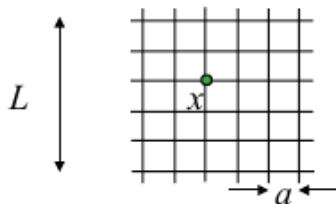
Consensus approach to HVP:

- $e^+e^- \rightarrow$ hadrons data + dispersion theory
- many data-sets + analyses, long-standing approach

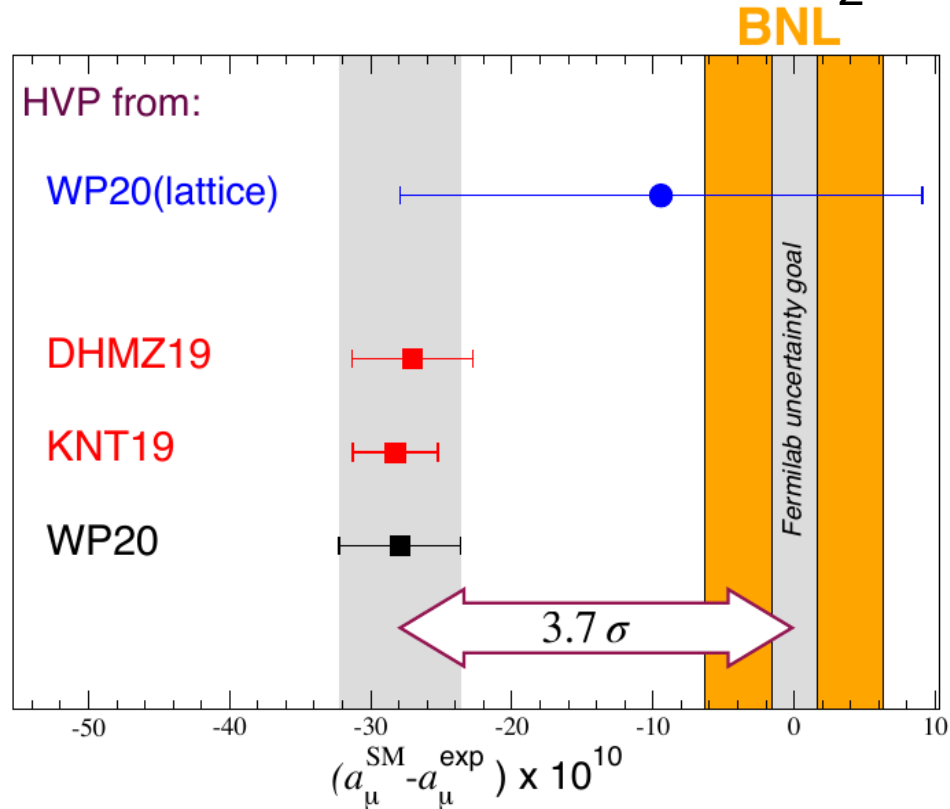


Lattice QCD:

- theory-based evaluation on super-computers
- huge recent progress, several groups



Results for anomalous term: $a = \frac{g - 2}{2}$



Brookhaven Experiment (E821)

- 540 ppm precision
- 3.7σ from Standard Model

Fermilab Experiment (E989)

- 540ppb \rightarrow 140 ppb
- better muon beam
 - lower inst rate, higher int rate
 - higher purity
- new detectors

BNL \rightarrow FNAL

[50 (stat) + 33 (syst) \rightarrow 11 (stat) + 11 (syst)] $\times 10^{-11}$

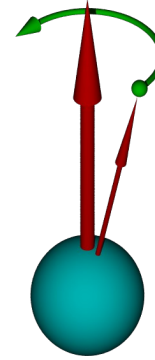
$$a_{\mu}^{\text{SM}} = 0.00116591810 (43) \quad \text{Theory Initiative}$$

$$a_{\mu}^{\text{Exp}} = 0.00116592089 (63) \quad \text{BNL E821 (2004)}$$

Measuring $g-2$

Spin precession in a magnetic field:

$$\begin{aligned}\omega_s &= g \frac{eB}{2m} + (1 - \gamma) \frac{eB}{\gamma m} \\ &= a_\mu \frac{eB}{m} + \frac{eB}{\gamma m}\end{aligned}$$

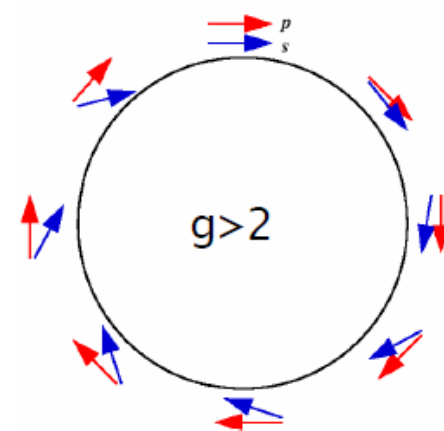
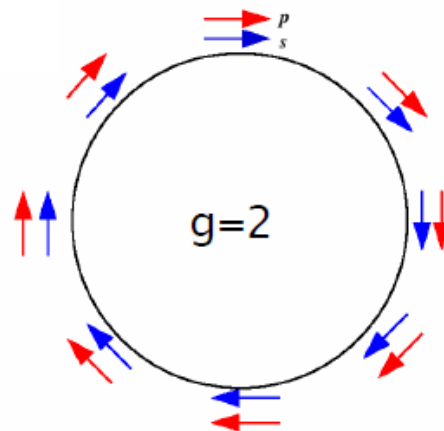


Spin precession in a magnetic field:

$$\omega_s = g \frac{eB}{2m} + (1 - \gamma) \frac{eB}{\gamma m}$$

$$= a_\mu \frac{eB}{m} + \frac{eB}{\gamma m}$$

$$\omega_s - \omega_c \equiv \omega_a = a_\mu \frac{eB}{m}$$

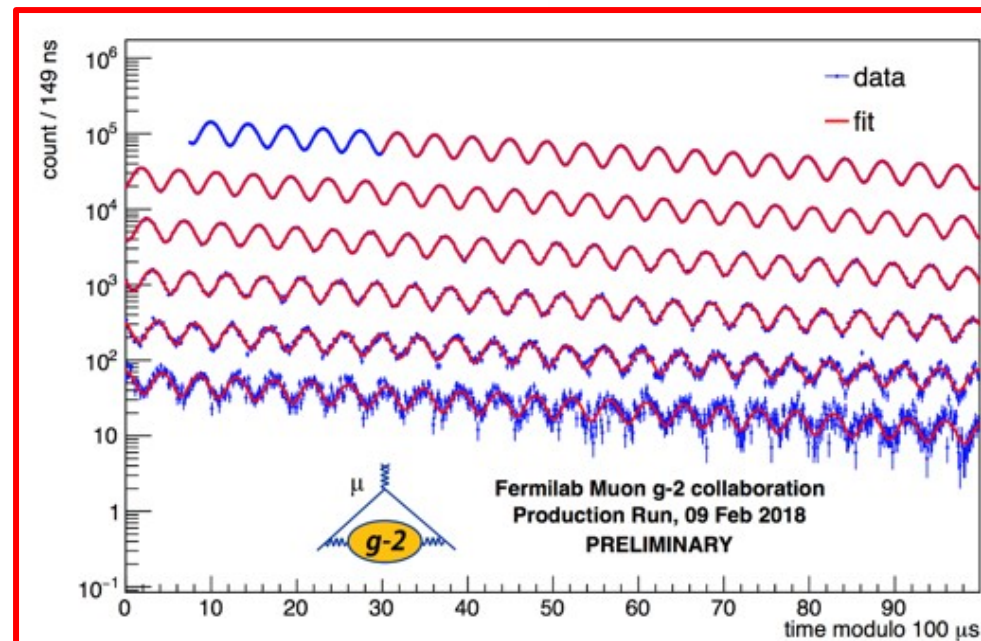
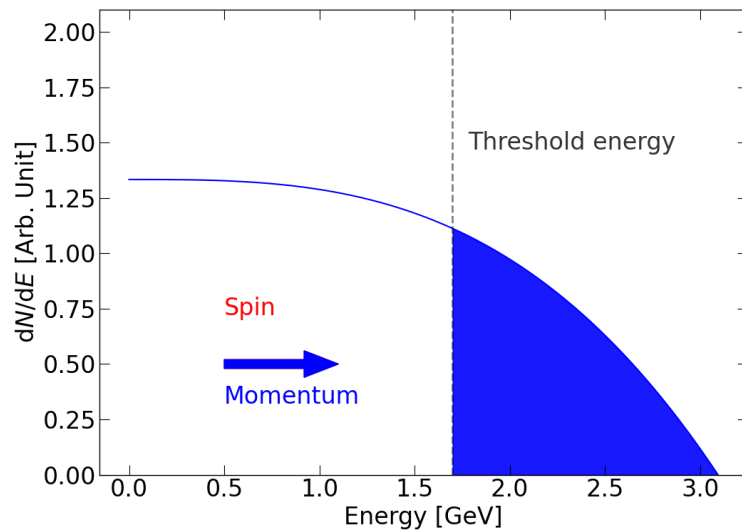
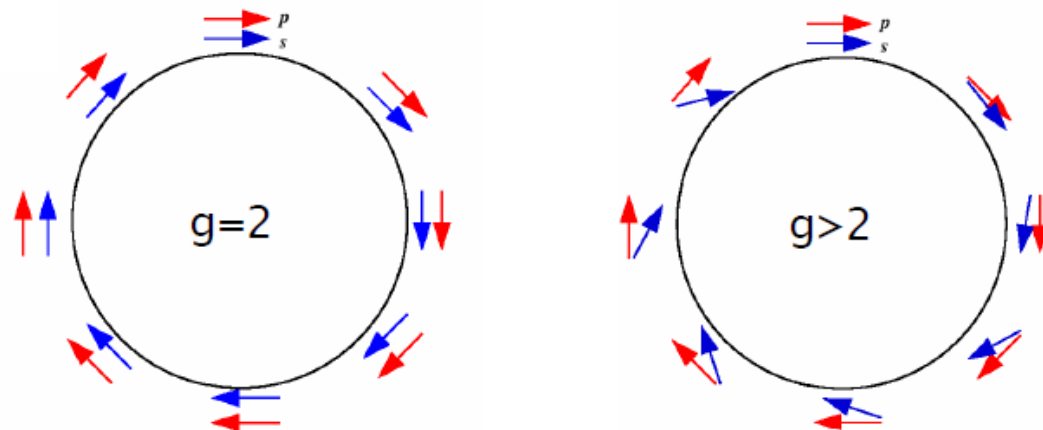


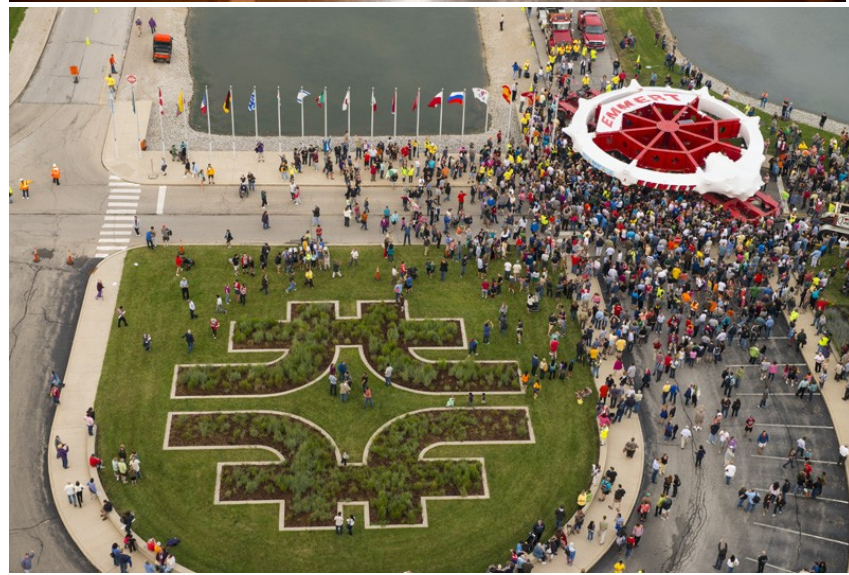
Spin precession in a magnetic field:

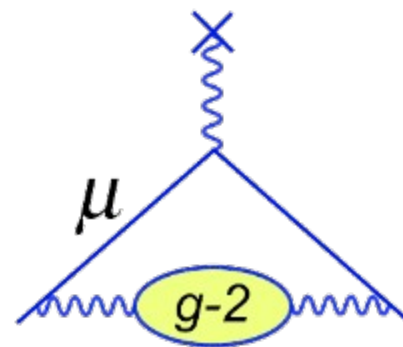
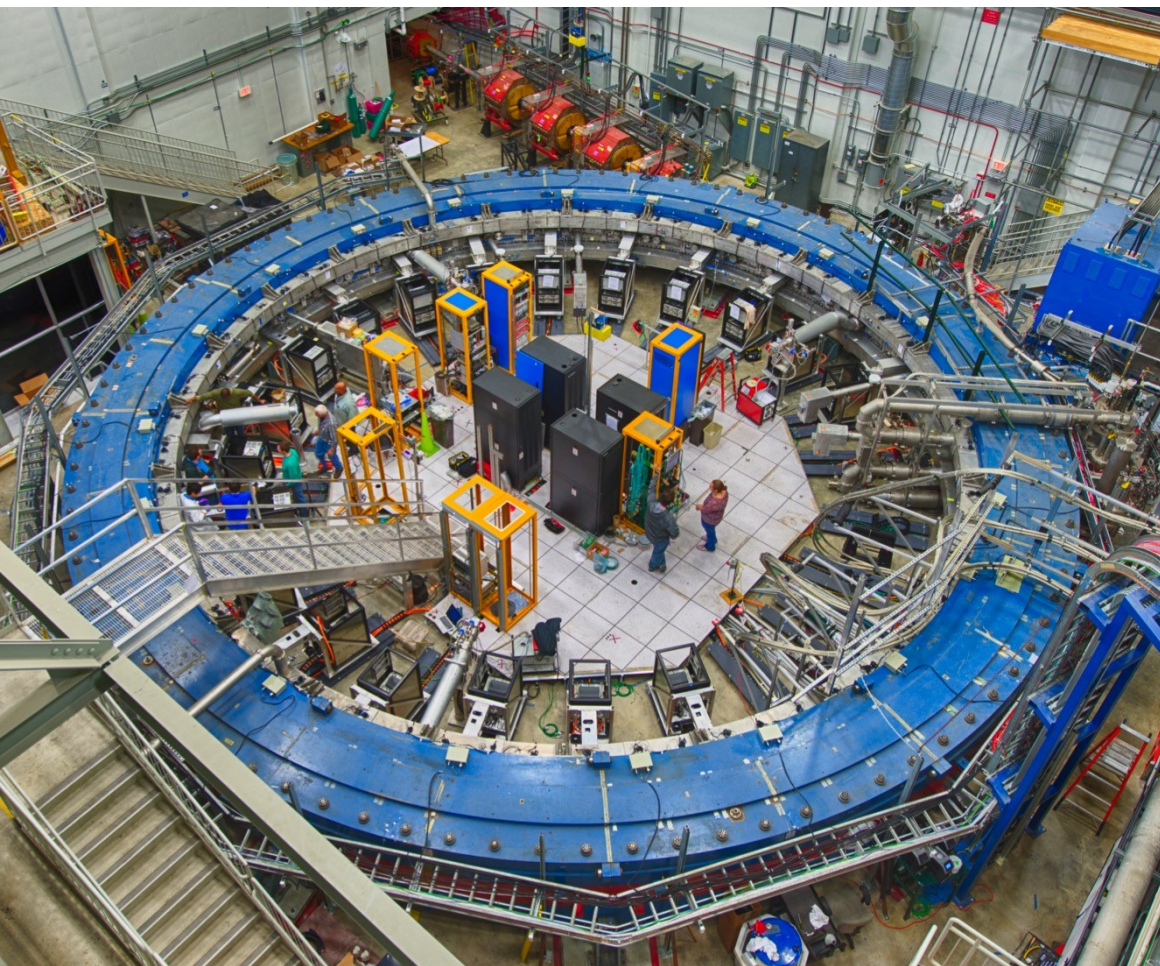
$$\omega_s = g \frac{eB}{2m} + (1 - \gamma) \frac{eB}{\gamma m}$$

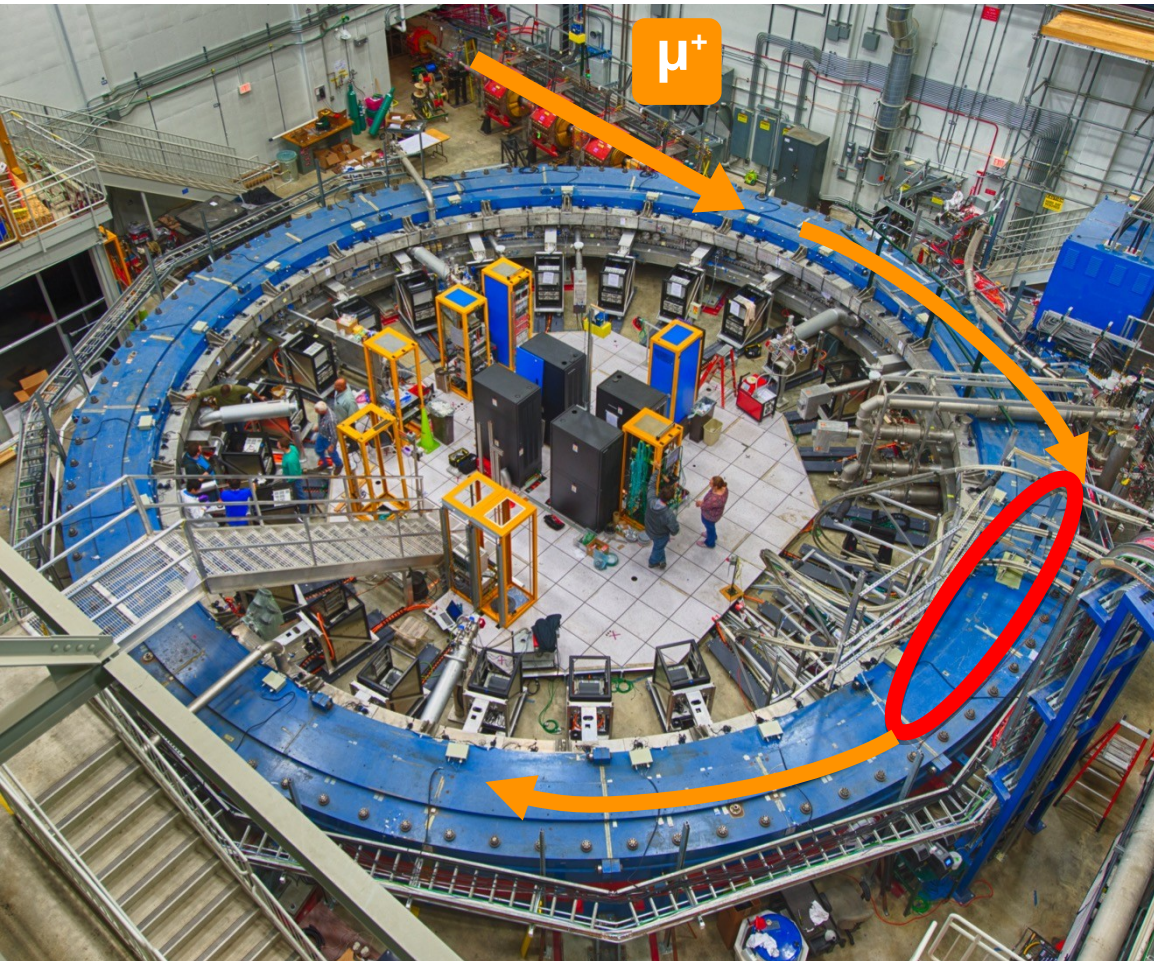
$$= a_\mu \frac{eB}{m} + \frac{eB}{\gamma m}$$

$$\omega_s - \omega_c \equiv \omega_a = a_\mu \frac{eB}{m}$$



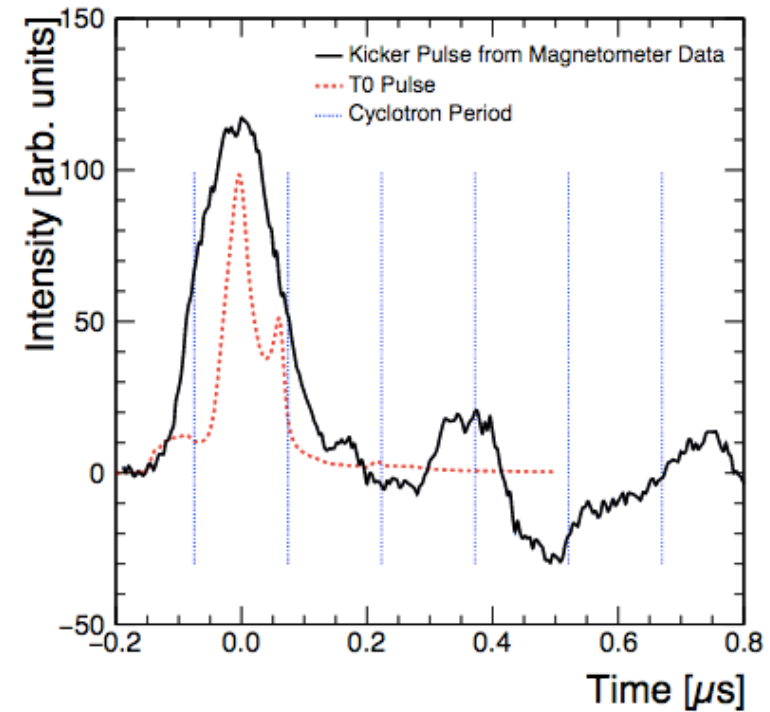


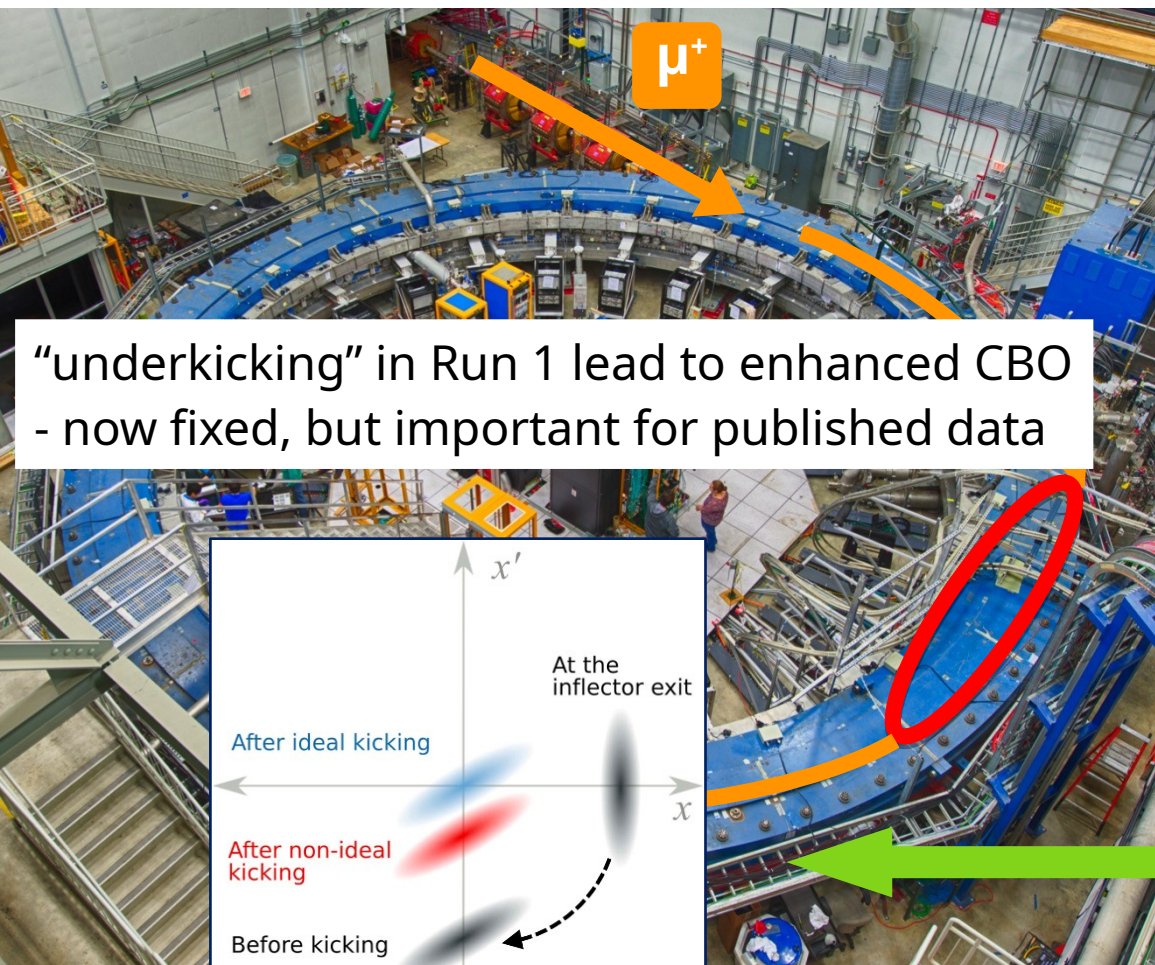




Beam pulse ~ 125 ns wide

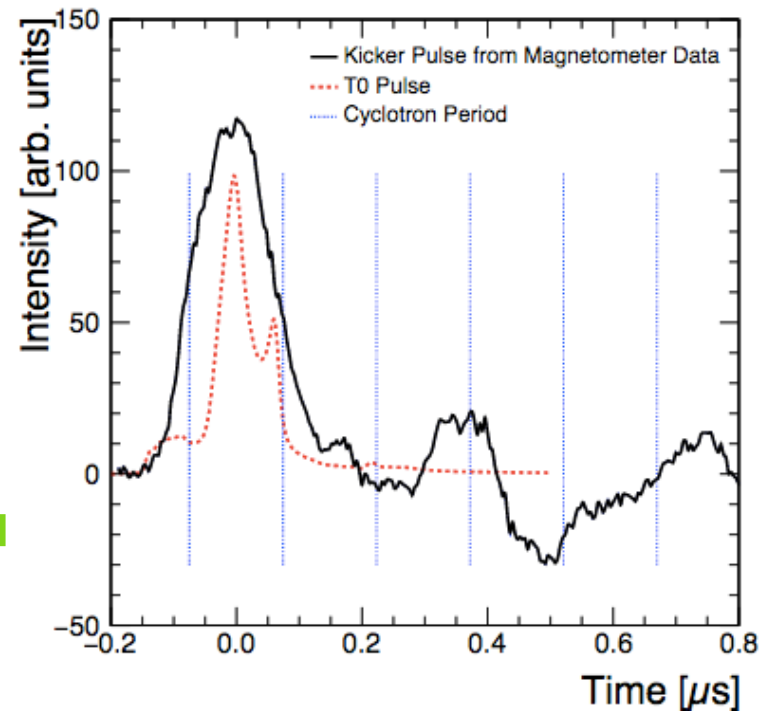
Muons ~ 77 mm away from ideal radius
- "kick" them into right orbit

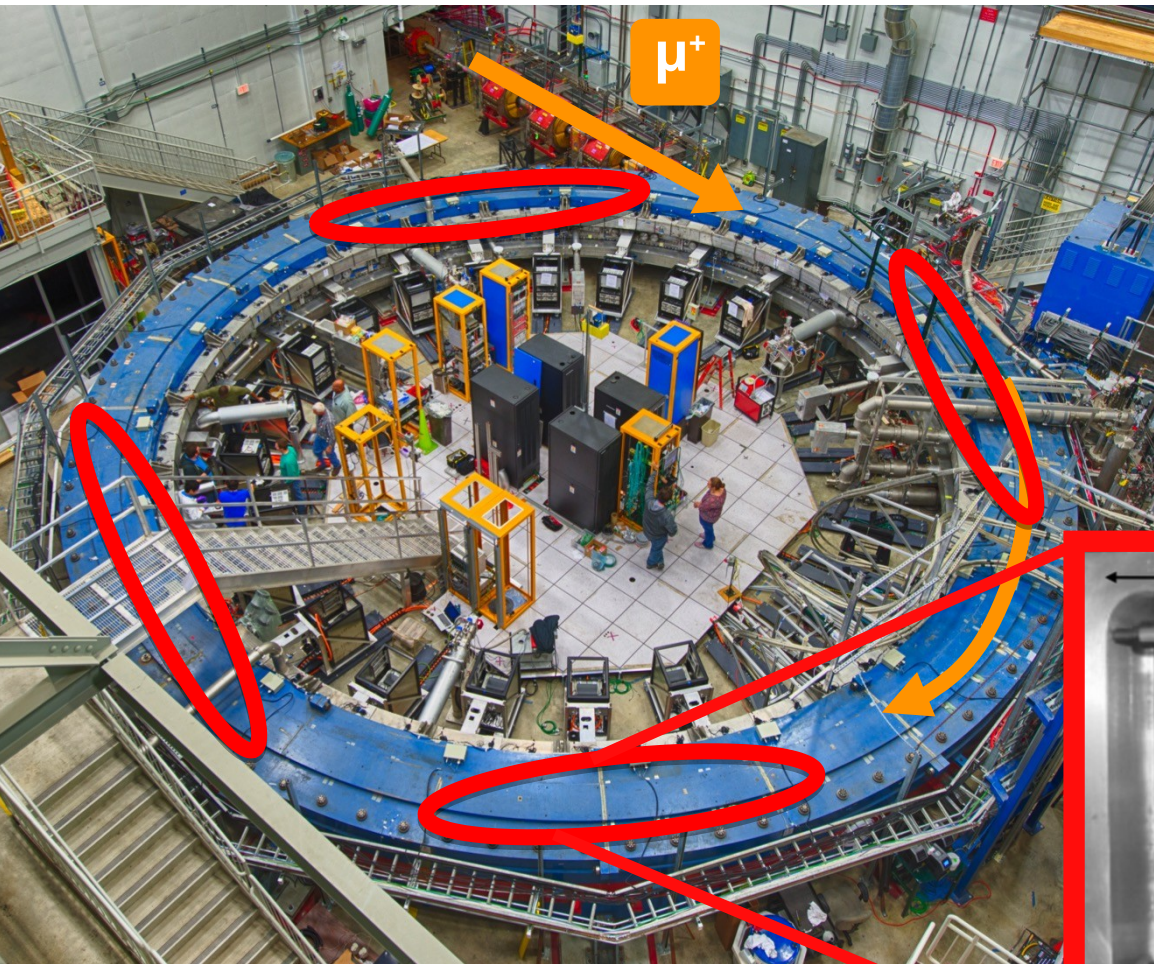




Beam pulse ~ 125 ns wide

Muons ~77 mm away from ideal radius
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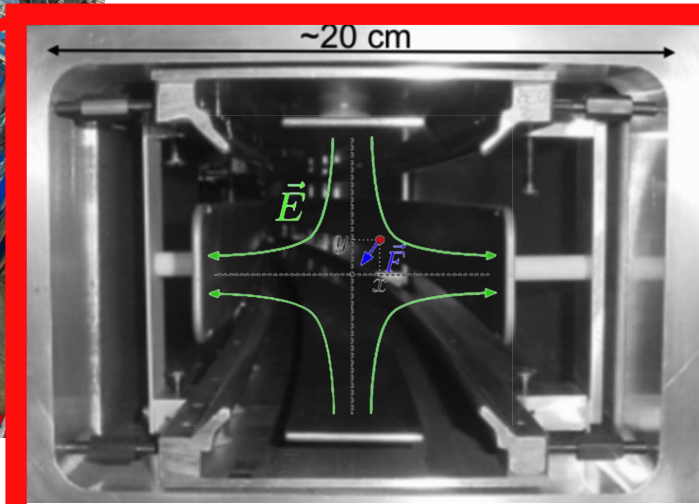


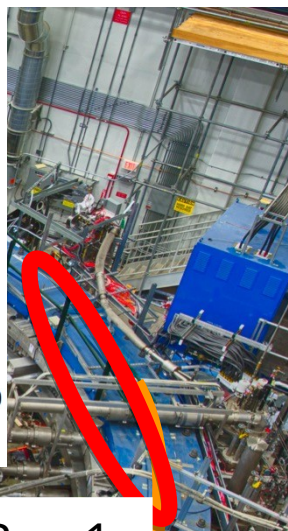
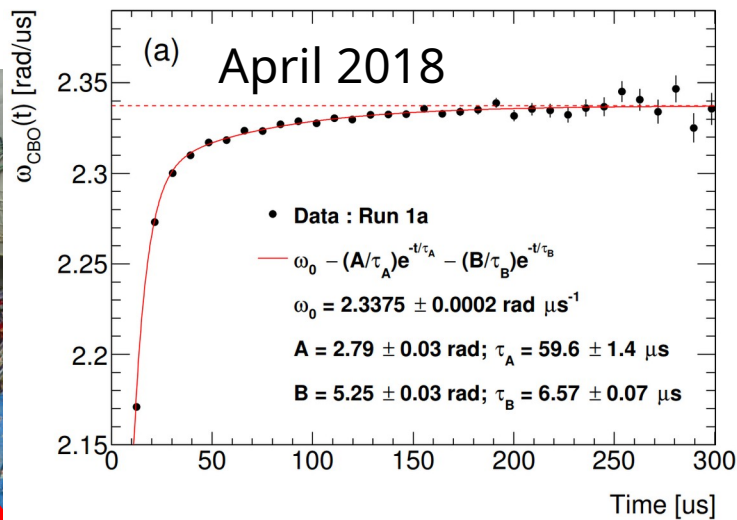
Scraping at the start of each fill:

- beam is moved to collimators
- momentum spread 0.15% \rightarrow 0.1%

Electric quadrupoles:

- focus the beam vertically





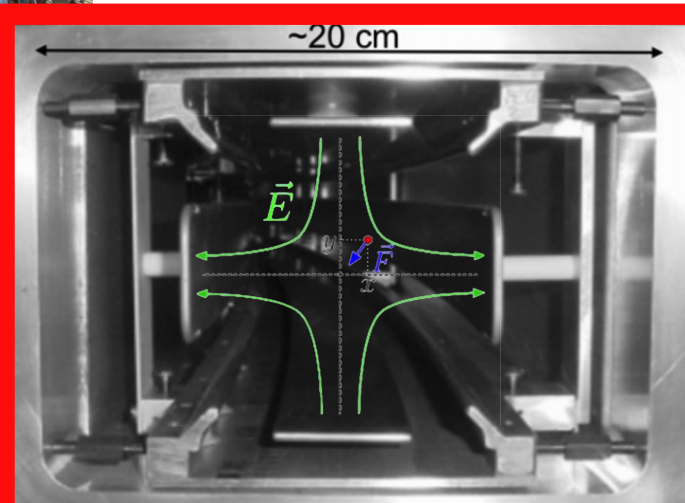
Scraping at the start of each fill:

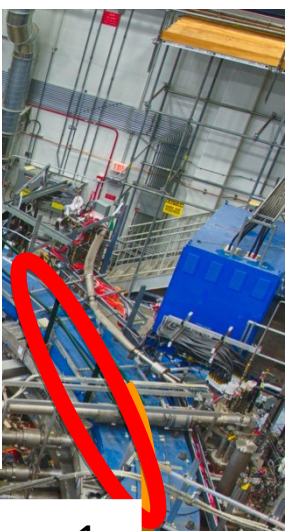
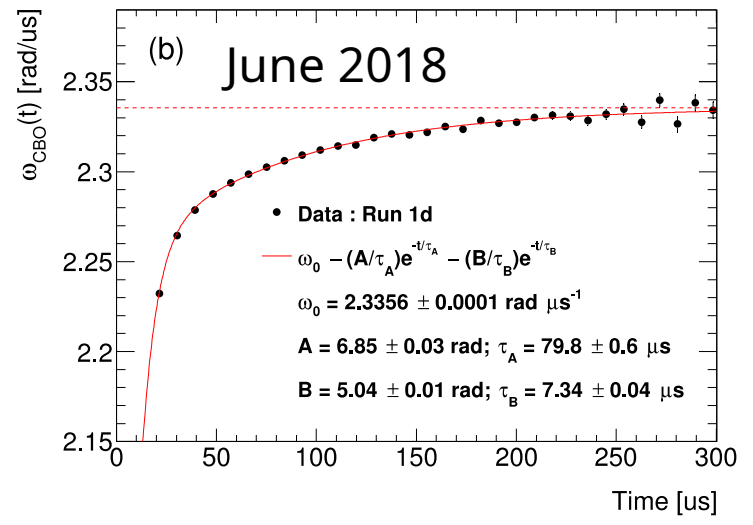
- beam is moved to collimators
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Electric quadrupoles:

- focus the beam vertically

- 2 (out of 32) damaged resistors in Run 1
- affected quadrupole charging time
 - drift in CBO
 - enhanced phase-acceptance systematic





Scraping at the start of each fill:

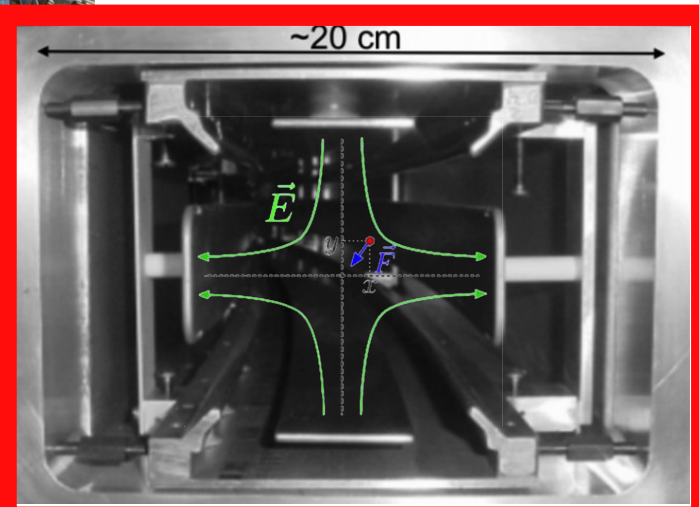
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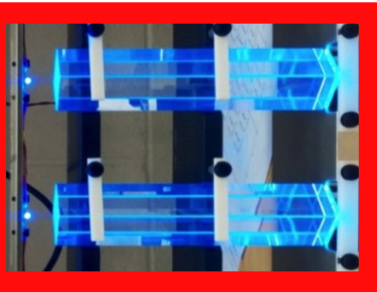
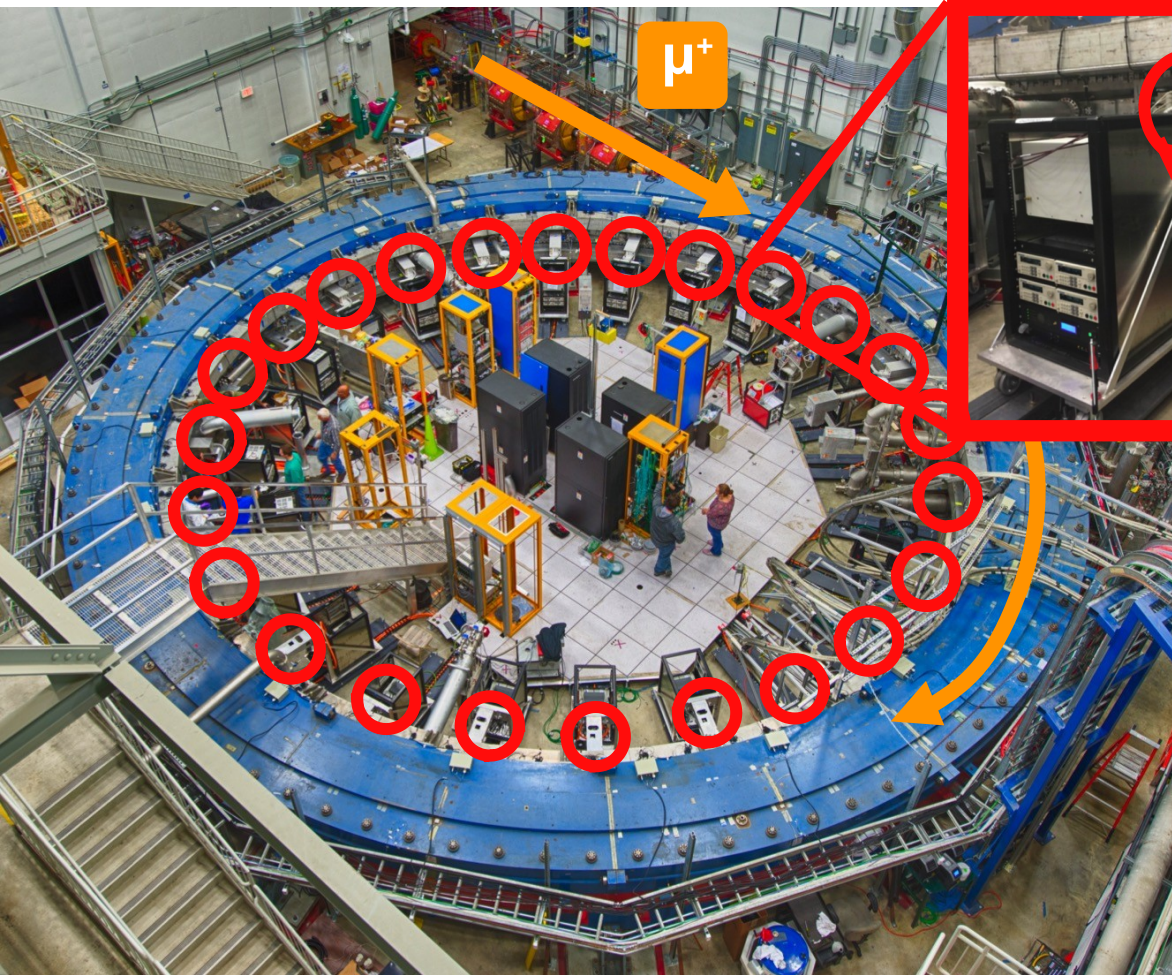
Electric quadrupoles:

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2 (out of 32) damaged resistors in Run 1

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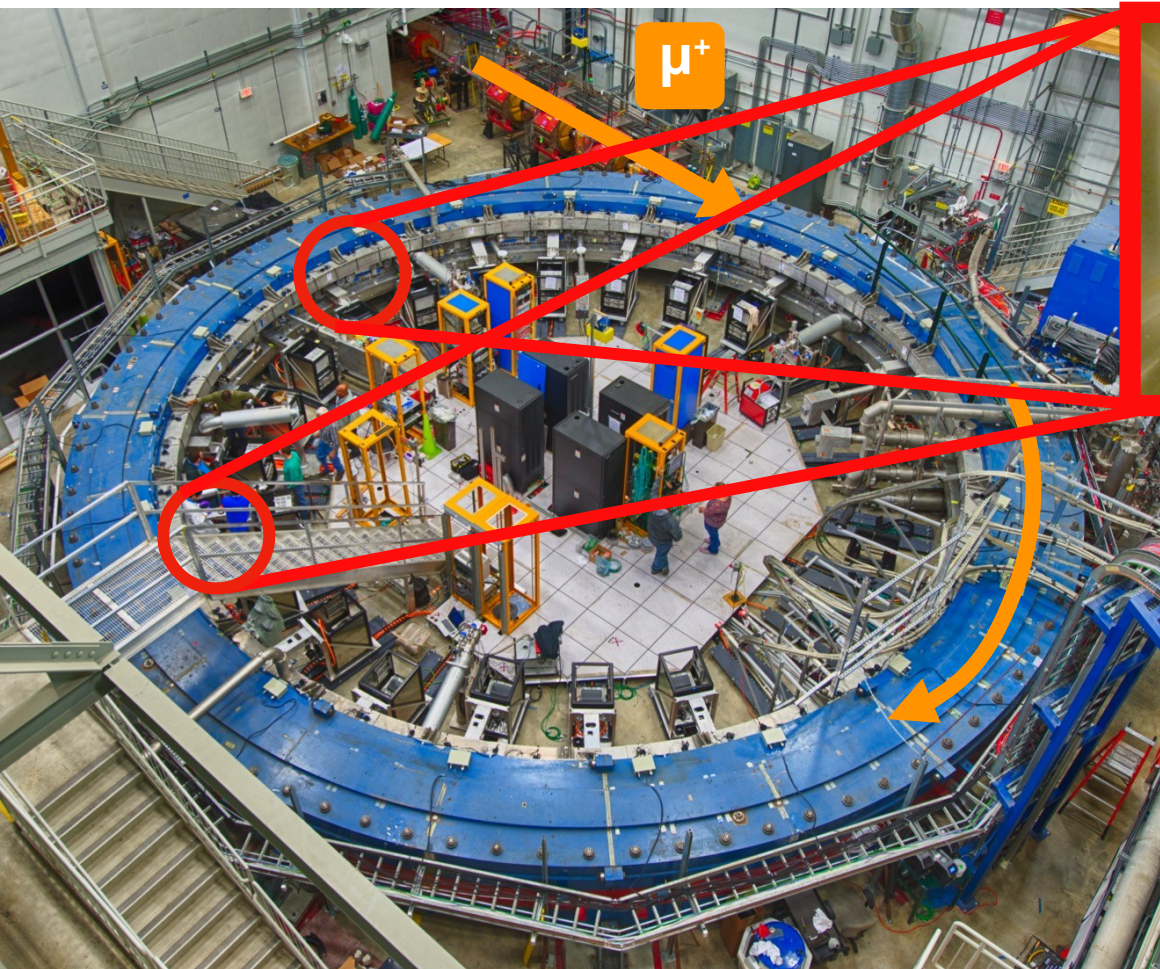


24 calorimeters:

- 6x9 PbF₂ crystals (2.5x2.5x14cm, 15X₀)

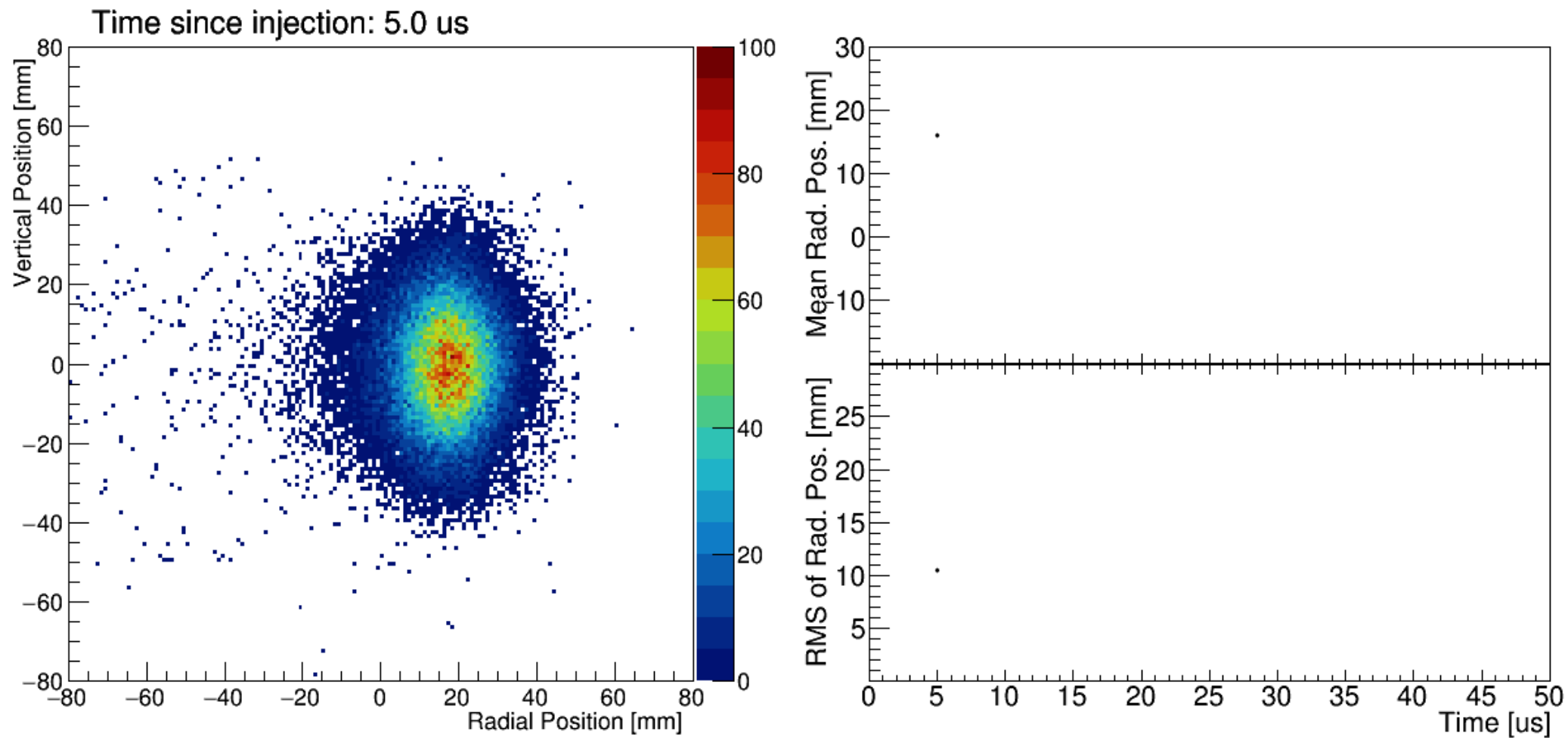
Read out by SiPMs, 1296 total channels

Dedicated laser calibration system



Two tracking stations:

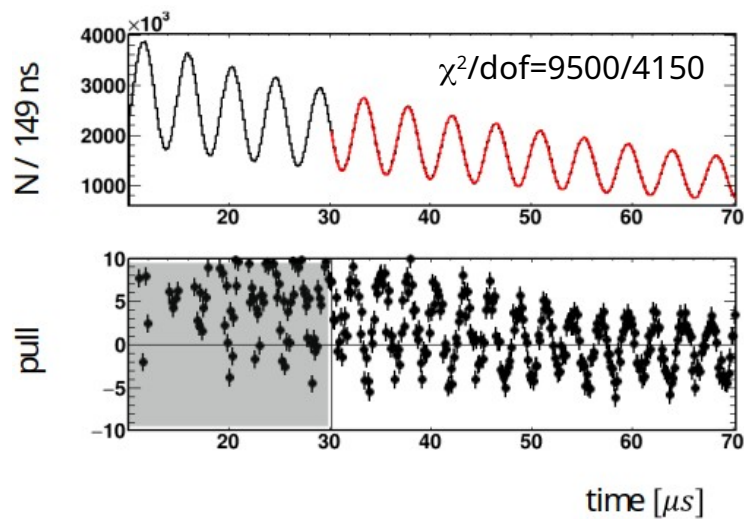
- 8 modules, 128 straws in each
- stereo angle for vertical position
- trace positrons back to decay point
- UK built



The basic method:

- count positrons over threshold
- 5-parameter fit to determine ω_a

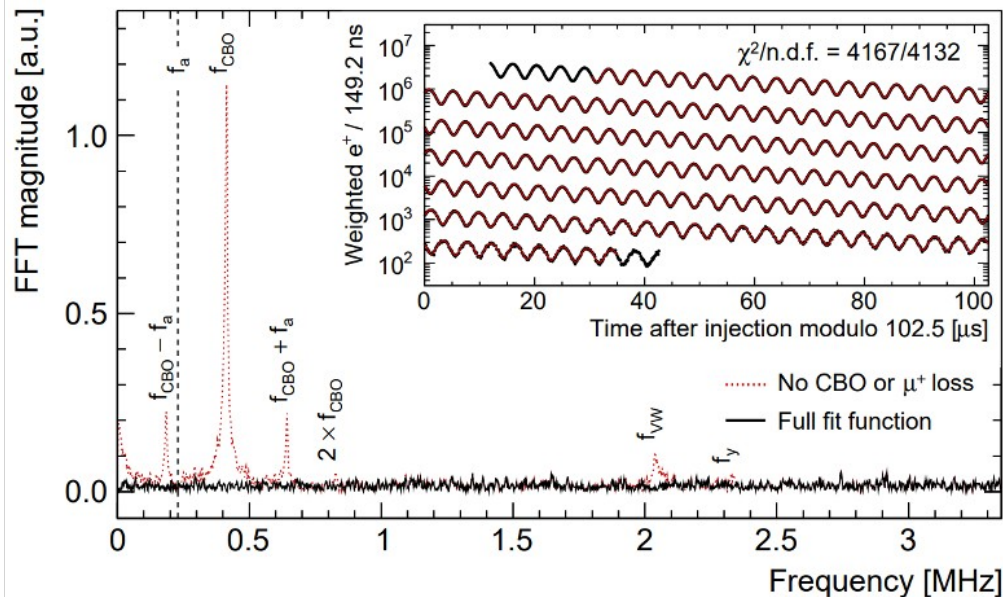
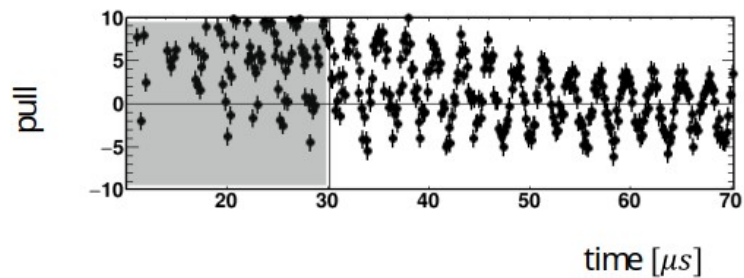
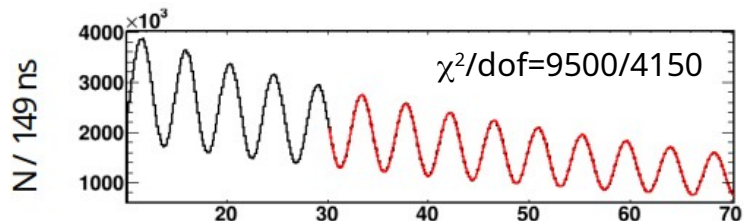
$$N(t) = N_0 \times e^{-t/\tau} \times [1 + A \cos(\omega_a t + \phi)]$$



The basic method:

- count positrons over threshold
- 5-parameter fit to determine ω_a

$$N(t) = N_0 \times e^{-t/\tau} \times [1 + A \cos(\omega_a t + \phi)]$$



Real fit function includes:

- calorimeter gain change (from laser data)
- positron pileup
- Muon losses
- Coherent betatron oscillations

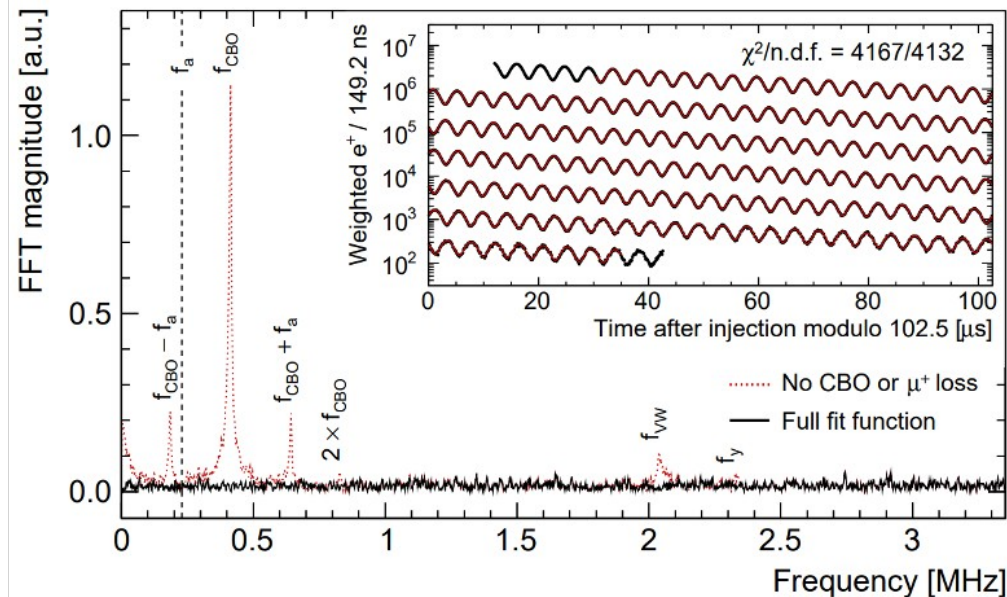
$$N_x(t) = 1 + e^{-t/\tau_{\text{CBO}}} A_{N,x,1} \cos(\omega_{\text{CBO}} t + \phi_{N,x,1}) + e^{-2t/\tau_{\text{CBO}}} A_{N,x,2} \cos(2\omega_{\text{CBO}} t + \phi_{N,x,2})$$

$$N_y(t) = 1 + e^{-t/\tau_y} A_{N,y,1} \cos(\omega_y t + \phi_{N,y,1}) + e^{-2t/\tau_y} A_{N,y,2} \cos(\omega_{\text{VW}} t + \phi_{N,y,2})$$

$$\Lambda(t) = 1 - K_{\text{loss}} \int_0^t e^{t'/\gamma\tau_\mu} L(t') dt'$$

$$A_x(t) = 1 + e^{-t/\tau_{\text{CBO}}} A_{A,x,1} \cos(\omega_{\text{CBO}} t + \phi_{A,x,1})$$

$$\phi_x(t) = 1 + e^{-t/\tau_{\text{CBO}}} A_{\phi,x,1} \cos(\omega_{\text{CBO}} t + \phi_{\phi,x,1})$$



- 2 positron reconstruction methods
- 4 different analysis methods (T,A,R,Q)
- 6 independent analysis teams, using different combinations
- Many cross-checks performed, varying sensitivities to systematics

Just a few details

$$a_{\mu} = \omega_a \frac{m}{eB}$$

$$a_\mu = \omega_a \frac{m}{eB} \longrightarrow a_\mu = \frac{\omega_a}{\tilde{\omega}_p(\text{Tr})} \frac{\mu_p(\text{Tr})}{\mu_e(\text{H})} \frac{\mu_e(\text{H})}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2}$$

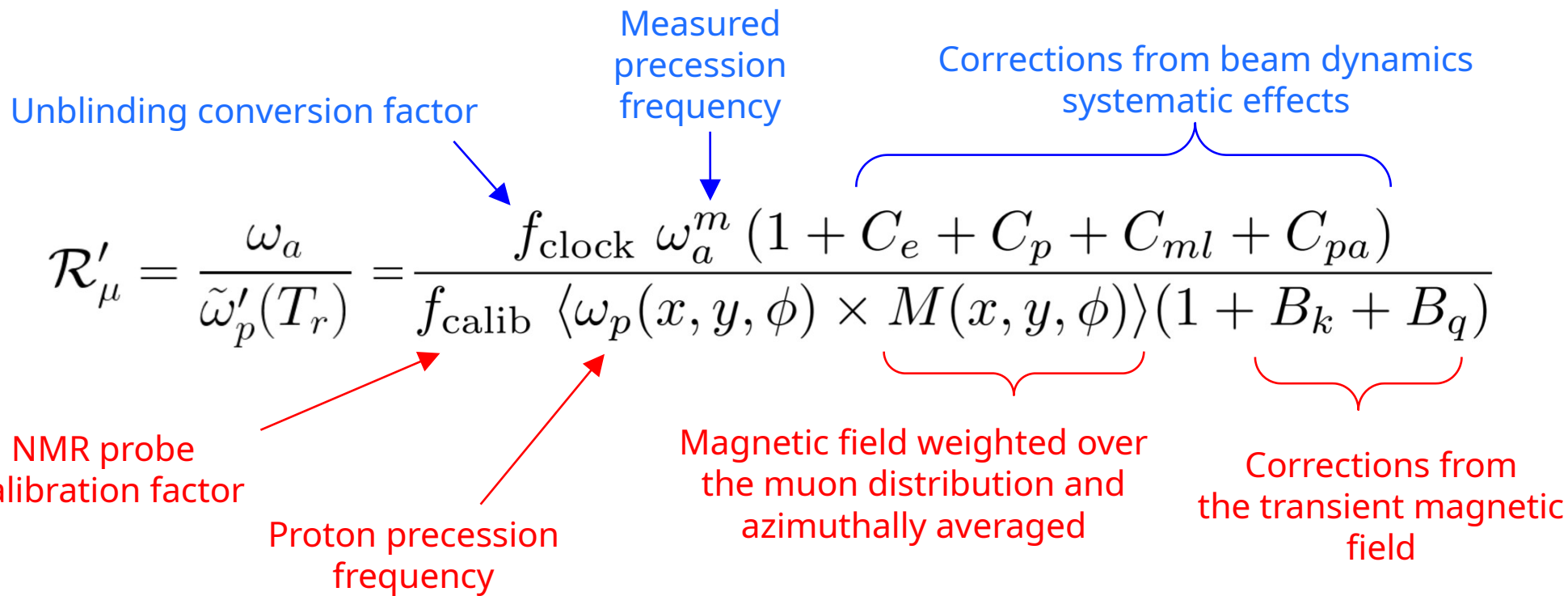
Total uncertainty from external inputs: 24 ppb

We determine the ratio:

$$\mathcal{R}'_\mu = \frac{\omega_a}{\tilde{\omega}_p(\text{Tr})}$$

$\frac{\mu'_p(T_r)}{\mu_e(H)}$	10.5 ppb uncertainty at $T_r = 34.7^\circ\text{C}$ <i>Metrologia</i> 13, 179 (1977)
$\frac{\mu_e(H)}{\mu_e}$	Bound state QED calculation, exact <i>Rev. Mod. Phys.</i> 88, 035009 (2016)
$\frac{m_\mu}{m_e}$	Muonium hyperfine splitting 22 ppb uncertainty <i>Phys. Rev. Lett.</i> 82, 11 (1999)
$\frac{g_e}{2}$	0.28 ppt uncertainty <i>Phys. Rev. A</i> 83, 052122 (2011)

$$a_\mu = \omega_a \frac{m}{eB} \quad \longrightarrow \quad a_\mu = \frac{\omega_a}{\tilde{\omega}_p(\text{Tr})} \frac{\mu_p(\text{Tr})}{\mu_e(\text{H})} \frac{\mu_e(\text{H})}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2}$$



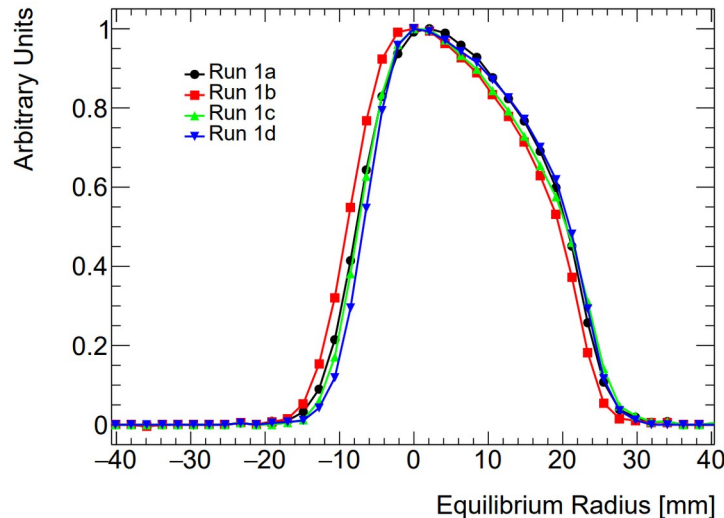
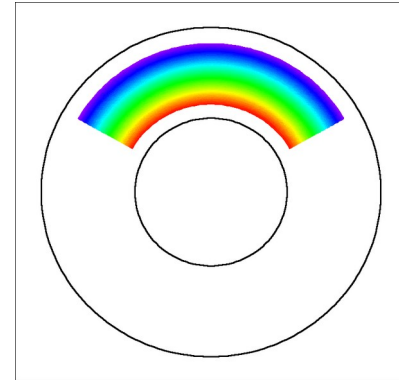
$$\omega_a = a_\mu \frac{eB}{m}$$

$$\vec{\omega}_a = \frac{e}{m} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \vec{\beta} \times \vec{E} - a_\mu \left(\frac{\gamma}{\gamma + 1} \right) (\vec{\beta} \cdot \vec{B}) \vec{\beta} \right]$$

$$\vec{\omega}_a = \frac{e}{m} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \vec{\beta} \times \vec{E} - a_\mu \left(\frac{\gamma}{\gamma + 1} \right) (\vec{\beta} \cdot \vec{B}) \vec{\beta} \right]$$

Coupling to electrostatic quadrupoles:

- use “magic momentum” of 3.1 GeV, but momentum spread $\sim 0.10\%$
 \rightarrow “E-field correction”
- from Fourier analysis of “fast-rotation”



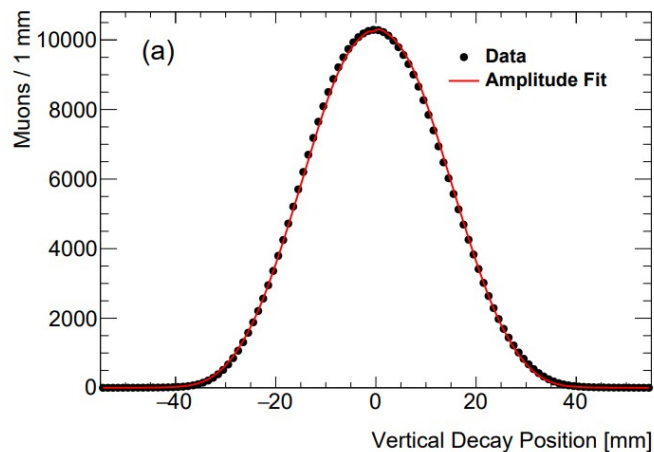
	Data Set	Run-1a	Run-1b	Run-1c	Run-1d
	C_e	471	464	534	475
	Stat. uncertainty	< 1	1	< 1	< 1
	Fourier method	8	13	14	4
	Momentum-time	52	52	52	52
	Quadrupole calibration	6	6	6	6
	Field index	2	2	2	4
	Syst. uncertainty	53	54	54	53

$$\vec{\omega}_a = \frac{e}{m} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \vec{\beta} \times \vec{E} - a_\mu \left(\frac{\gamma}{\gamma + 1} \right) (\vec{\beta} \cdot \vec{B}) \vec{\beta} \right]$$

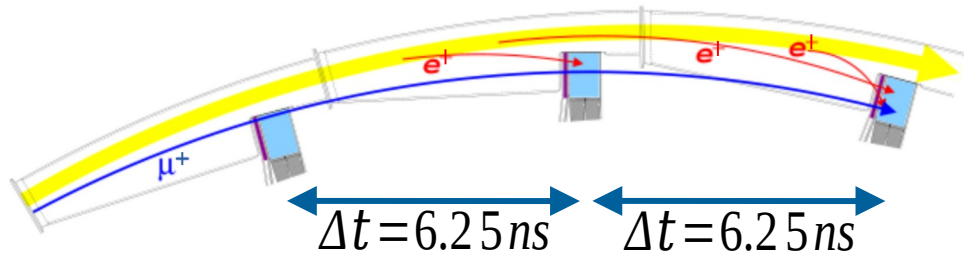
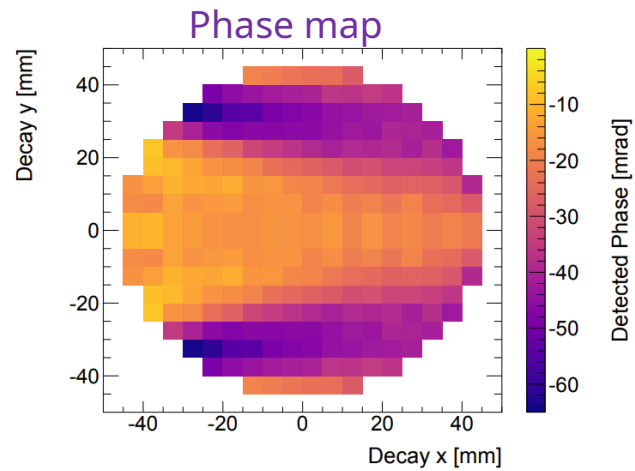
Small vertical momentum component:

→ “Pitch correction”

- use tracking detectors to measure vertical width of beam

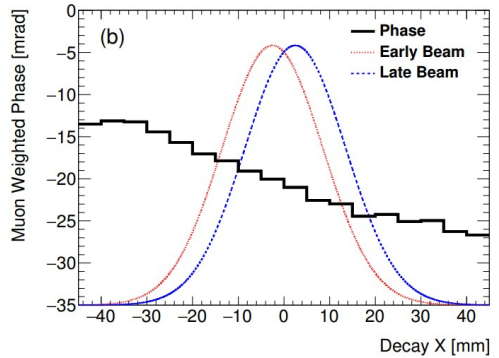
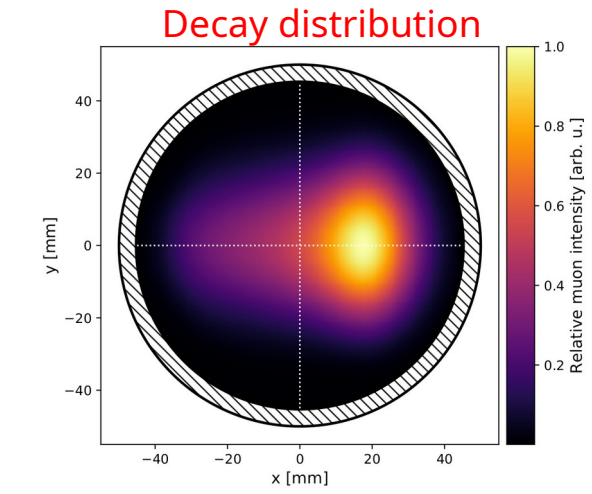


Data Set	Run-1a	Run-1b	Run-1c	Run-1d
C_p (ppb)	176	199	191	166
Stat. uncertainty	< 1	< 1	< 1	< 1
Tracker reco.	11	12	12	11
Tracker res. & acc	3	4	4	3
$\beta_y(\phi)$ & calo. acc.	1	1	2	1
Amplitude fit	1	< 1	1	3
Quad calibration	4	4	4	4
Syst. uncertainty	12	14	14	12



“Phase/Acceptance” correction:

- changes in beam distribution during fill
→ time-dependent phase
- enhanced by damaged resistors in Run-1

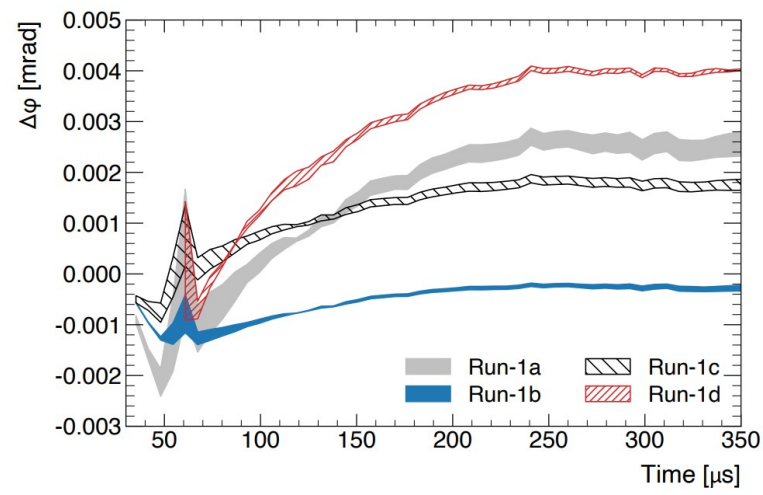
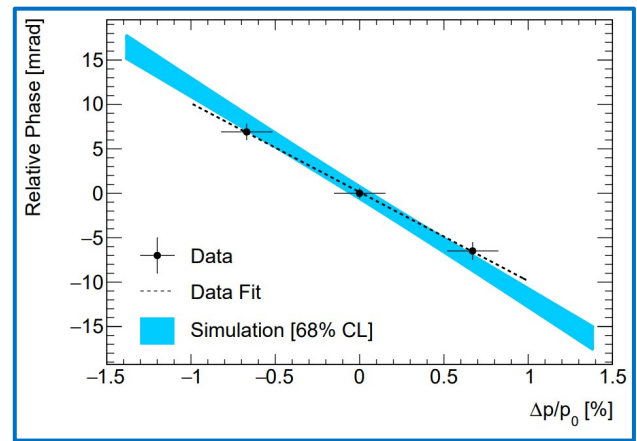


Data Set	Run-1a	Run-1b	Run-1c	Run-1d
C_{pa}	-184	-165	-117	-164
Stat. uncertainty	23	20	15	14
Tracker & CBO	73	43	41	44
Phase maps	52	49	35	46
Beam dynamics	27	30	22	45
Total uncertainty	96	74	60	80

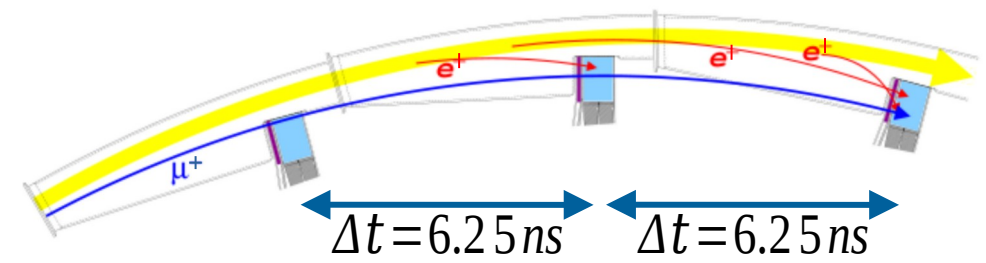
Lost Muons:

- low momentum muons are lost more quickly than high
- slightly different phase to ensemble
- causes change in overall phase vs time

$$\frac{d\phi_0}{dt} = \frac{d\phi_0}{d\langle p \rangle} \frac{d\langle p \rangle}{dt}$$



Data Set	Run-1a	Run-1b	Run-1c	Run-1d
C_{ml}	-14	-3	-7	-17
Phase-momentum	2	0	1	3
Form of $l(t)$	2	0	1	1
f_{loss} function	2	1	2	2
Linear sum ($\sigma_{C_{ml}}$)	6	2	4	6

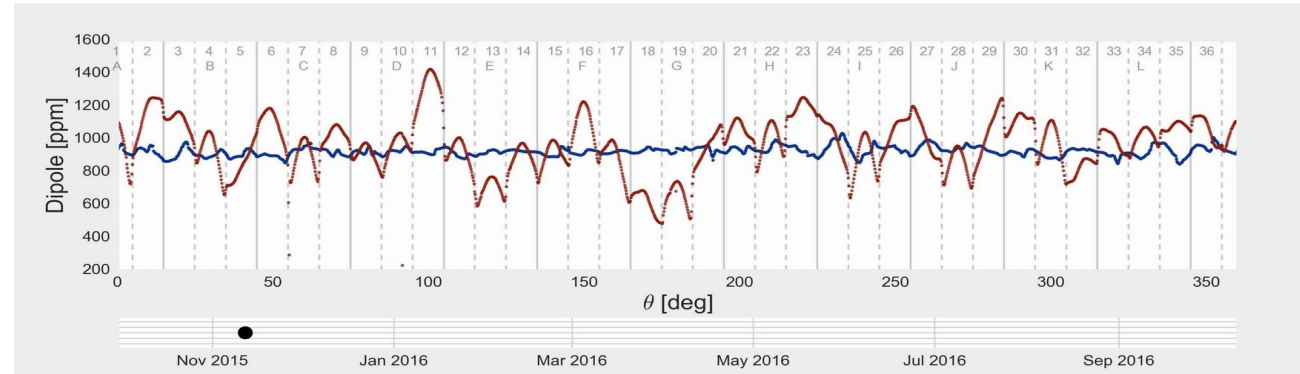
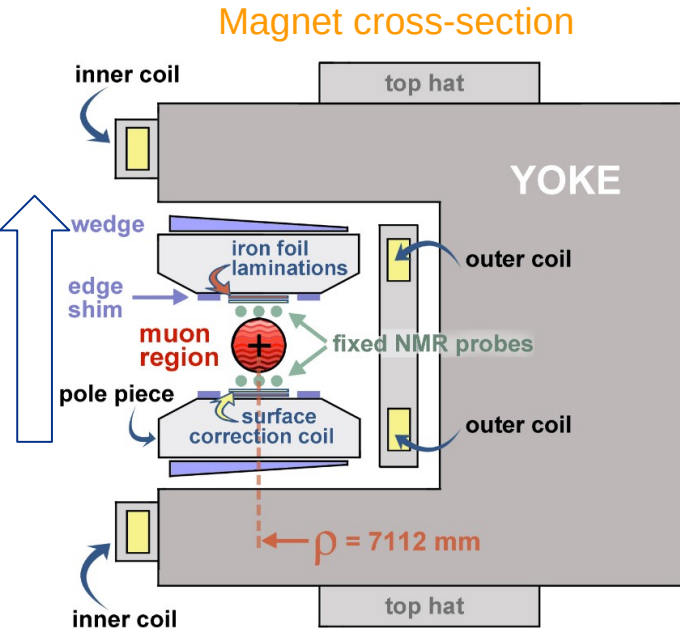


$$a_\mu = \omega_a \frac{m}{eB} \quad \longrightarrow \quad a_\mu = \frac{\omega_a}{\tilde{\omega}_p(\text{Tr})} \frac{\mu_p(\text{Tr})}{\mu_e(\text{H})} \frac{\mu_e(\text{H})}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2}$$

$$\mathcal{R}'_\mu = \frac{\omega_a}{\tilde{\omega}'_p(T_r)} = \frac{f_{\text{clock}} \omega_a^m (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{\text{calib}} \underbrace{\langle \omega_p(x, y, \phi) \times M(x, y, \phi) \rangle}_{\text{Magnetic field weighted over the muon distribution and azimuthally averaged}} \underbrace{(1 + B_k + B_q)}_{\text{Corrections from the transient magnetic field}}$$

Unblinding conversion factor → ω_a
Measured precession frequency → ω_a^m
Corrections from beam dynamics systematic effects → $(1 + C_e + C_p + C_{ml} + C_{pa})$
NMR probe calibration factor → f_{calib}
Proton precession frequency → f_{clock}
Magnetic field weighted over the muon distribution and azimuthally averaged → $\langle \omega_p(x, y, \phi) \times M(x, y, \phi) \rangle$
Corrections from the transient magnetic field → $(1 + B_k + B_q)$

7.112 m radius 'C'-shape magnet with vertically-aligned field $B = 1.45$ T, ppm-level uniformity



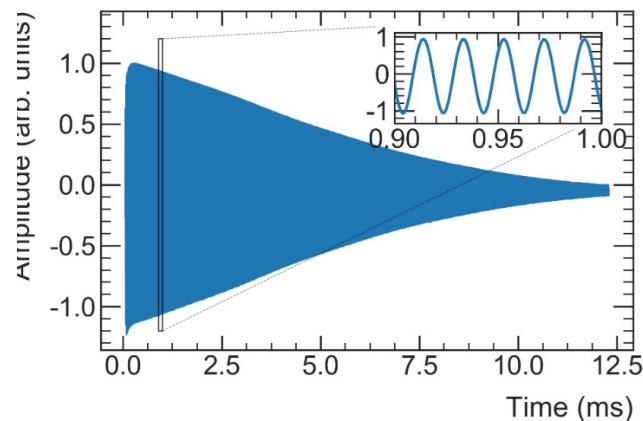
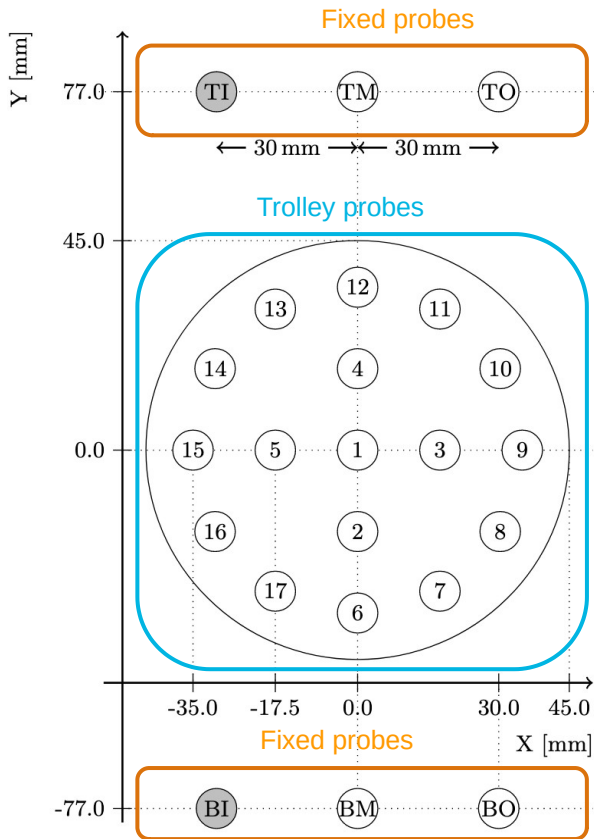
Temperature variations \rightarrow ppm changes in magnet geometry, and drift in the field

378 'fixed' NMR probes around the ring measure the drift continuously

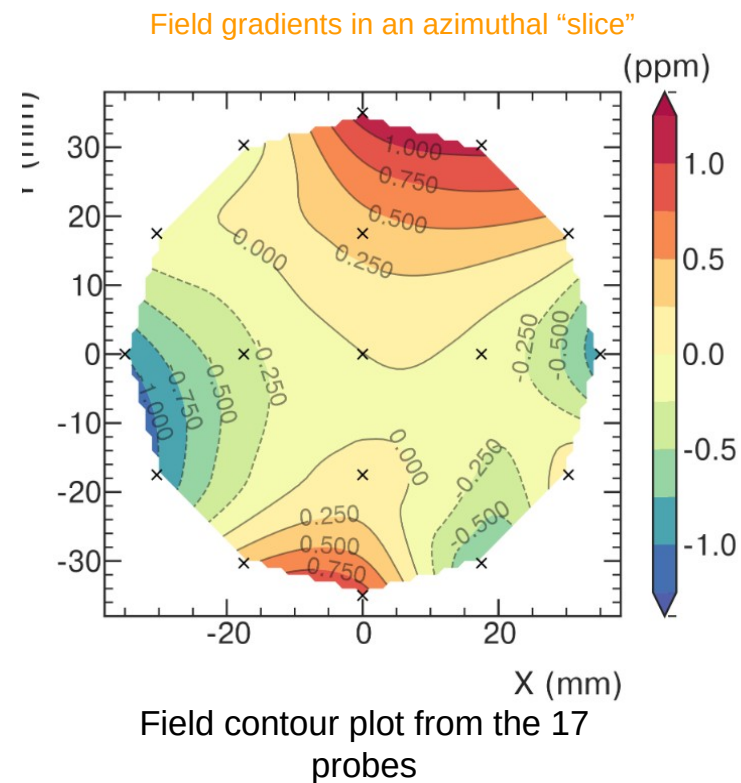
\rightarrow feedback to the magnet power supply to keep the dipole (vertical) term constant

An in-vacuum trolley drives around the ring every ~3 days

- 17 NMR probes x 9000 measurements, mapping out the field components



Field measured by extracting frequency from a Free Induction Decay (FID) spectrum

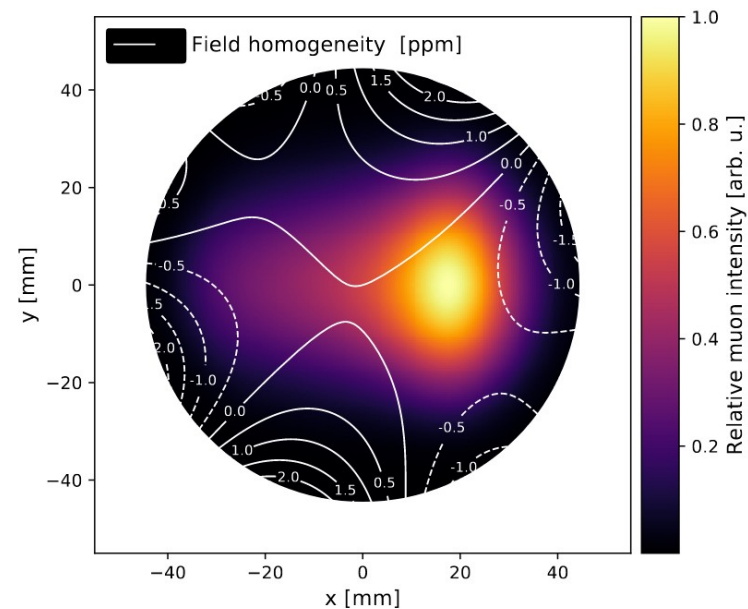
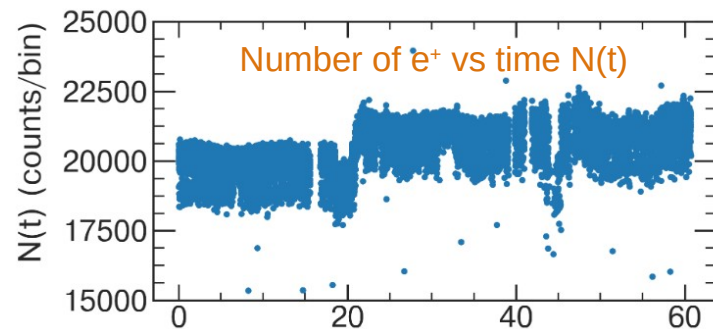


Frequency maps from trolley and fixed probe data need to be weighted by the muon distribution

$$\tilde{\omega}_p^{\parallel} = \left\langle \frac{\int \omega_p^{\parallel}(x, y, \phi) M(x, y, \phi) dx dy}{\int M(x, y, \phi) dx dy} \right\rangle$$

2D beam distribution obtained from the straw trackers

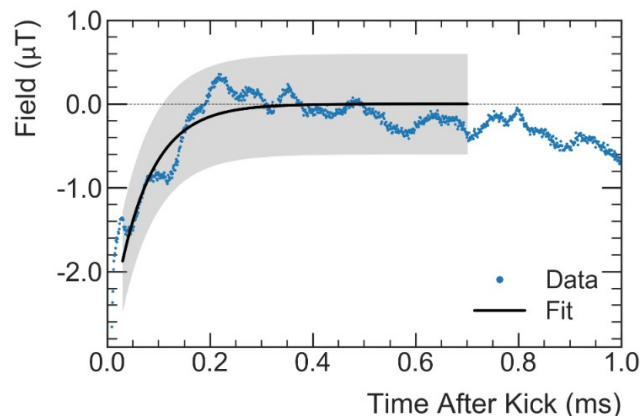
- includes beam dynamics information
- and detector acceptances



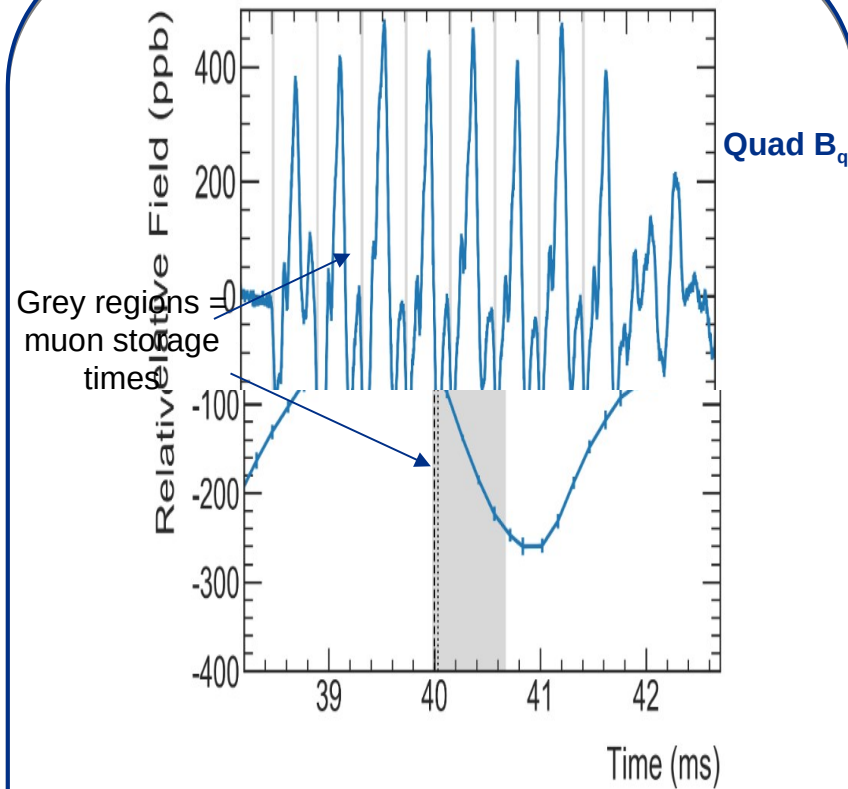
Muons experience “fast transient” fields from the pulsed kickers & quadrupoles

→ invisible to fixed probes due to shielding

Measured during dedicated campaigns



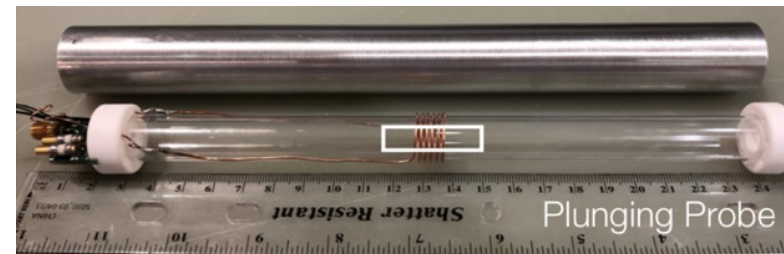
- Kicker pulse of **22 mT for 150 ns** just after muon injection.
- Field change caused by residual field after kicker pulse. Muons present from **30 μs to 700 μs** after the kick (fit region)
- Kicker correction: **-27 (37) ppb**



- Measured with a dedicated in-vacuum NMR probe located between quad plates during pulsing
- Quad correction: **-17 (92) ppb**

Specially designed “plunging probe”:

- characterised in very stable, homogenous solenoid
 - accuracy of 15 ppb
- used to calibrate trolley probes in-situ

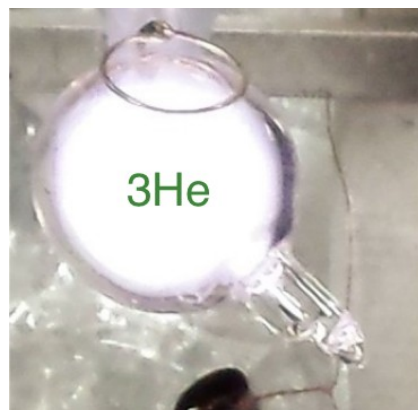


Calibrated against spherical, water-based probe from BNL to 6ppb

A novel ^3He NMR probe also developed (different systematics)

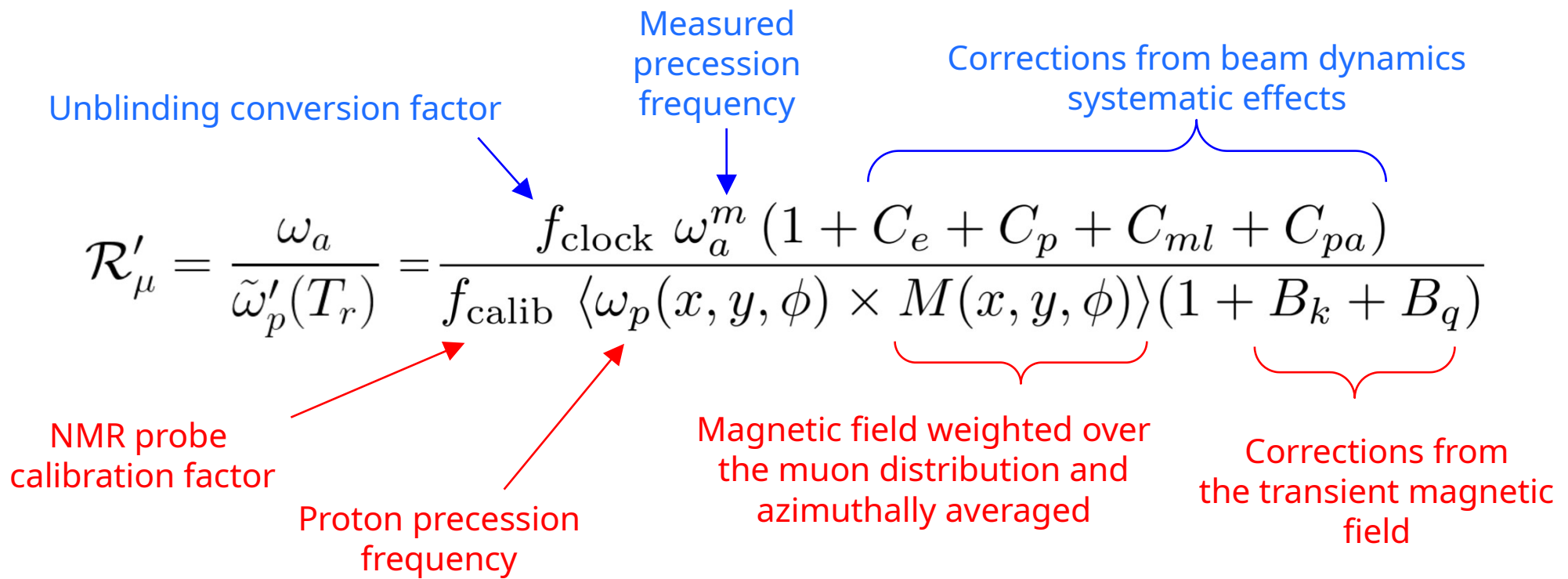
→ provides cross-check to 38 ppb (PRL 124, 223001 (2020))

Ongoing effort to cross-calibrate E989 probe with J-PARC Muon g-2/EDM probe



Quantity	Symbol	Value (ppb)	Uncertainty (ppb)
Diamagnetic Shielding T dep	δ^T	-99.1 to -86.0	5
Bulk Magnetic Susceptibility	δ^b	-1505.9 to -1505.6	6
Material Perturbation	δ^s	15.2	12
Water Sample and Sample Holder	δ^w	0	2
Radiation Damping	δ^{RD}	0	3
Proton Dipolar Fields	δ^d	0	2
TOTAL		-1589.8 to -1576.4	15

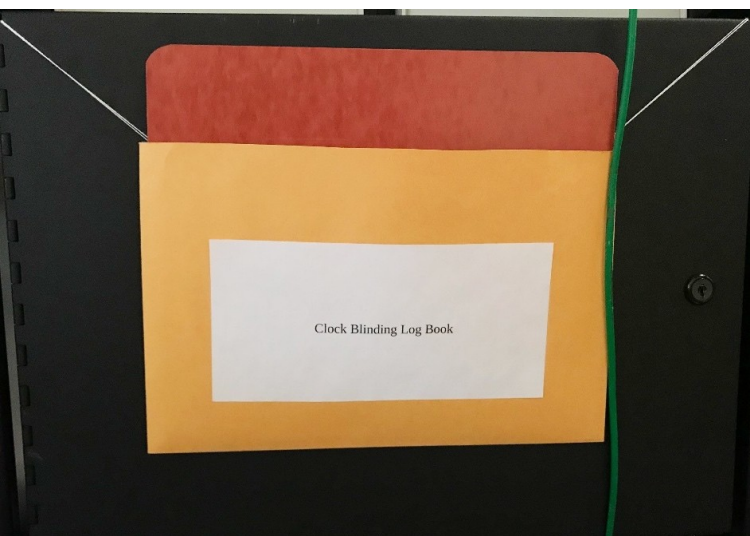
$$a_\mu = \frac{\omega_a}{\tilde{\omega}_p(\text{Tr})} \frac{\mu_p(\text{Tr})}{\mu_e(\text{H})} \frac{\mu_e(\text{H})}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2}$$



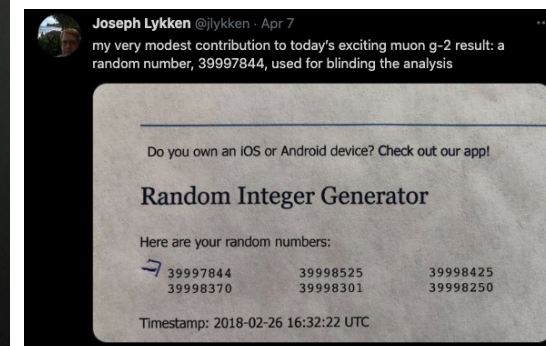
Hardware-blinded clock with frequency $(40 + x)$ MHz
- offset was ~ 25 ppm (approx 10x the BNL-SM difference)

Set by two people outside the collaboration, stored in locked cabinet!

Locked cabinet



Greg Bock and Joe Lykken (2018)



*"The code was picked by the Fermilab theorist,
and he is the only person to know it.*

*This theorist now refuses to give away the code.
It is not clear why.*

*One time he said he had forgotten the envelope with the code
on a train, another time he said the dog had eaten it."*

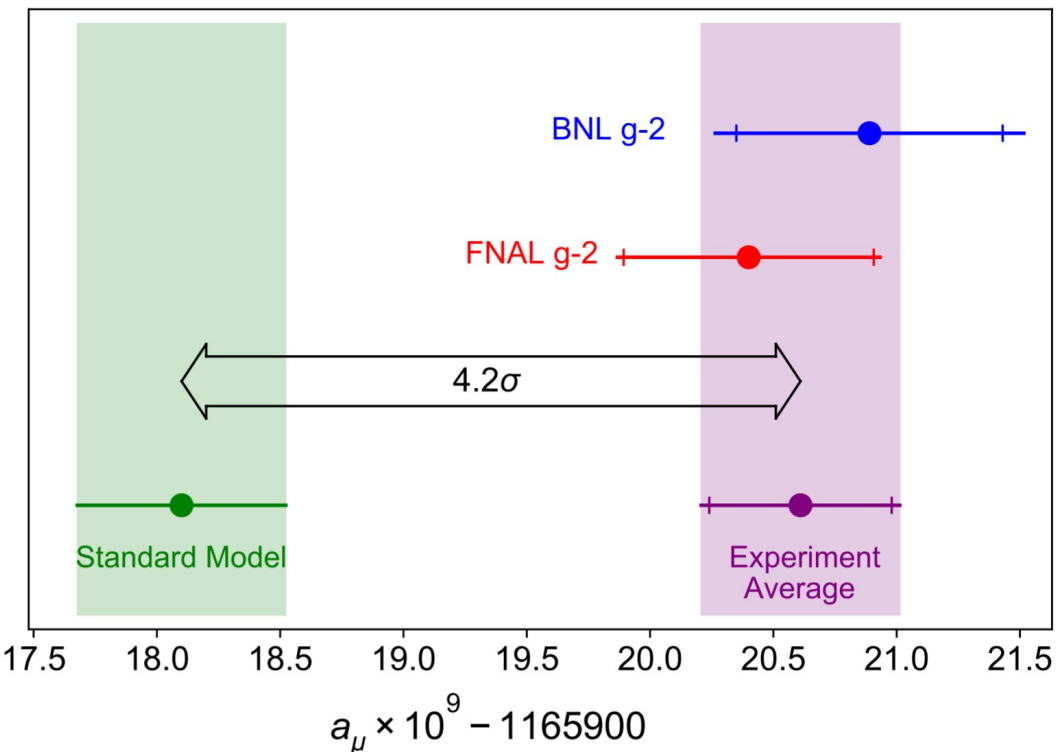
Unblinding meeting 25th Feb 2021

After all corrections uncertainties were finalized, the collaboration unanimously vote to unblind.



No changes to the result were made since then





Quantity	Correction Terms (ppb)	Uncertainty (ppb)
ω_a^m (statistical)	–	434
ω_a^m (systematic)	–	56
C_e	489	53
C_p	180	13
C_{ml}	-11	5
C_{pa}	-158	75
$f_{\text{calib}} \langle \omega_p(x, y, \phi) \times M(x, y, \phi) \rangle$	–	56
B_k	-27	37
B_q	-17	92
$\mu'_p(34.7^\circ)/\mu_e$	–	10
m_μ/m_e	–	22
$g_e/2$	–	0
Total systematic	–	157
Total fundamental factors	–	25
Totals	544	462

$$a_\mu(\text{FNAL}) = 116592040(54) \times 10^{-11}$$

$$a_\mu(\text{Exp}) = 116592061(41) \times 10^{-11}$$

$$a_\mu(\text{SM}) = 116591810(43) \times 10^{-11}$$

Discrepancy is ~2.5 ppm

- x2 the size of the electroweak contribution, x1/30 the size of QCD

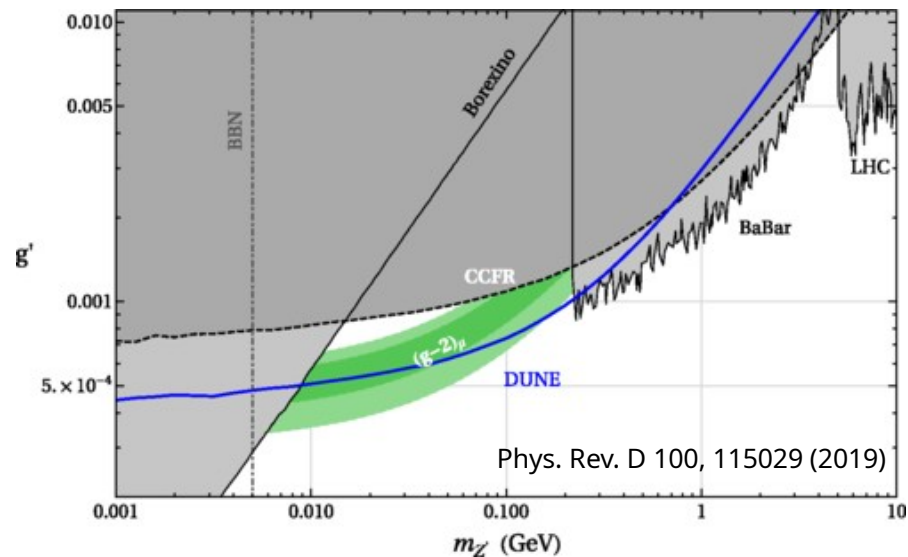
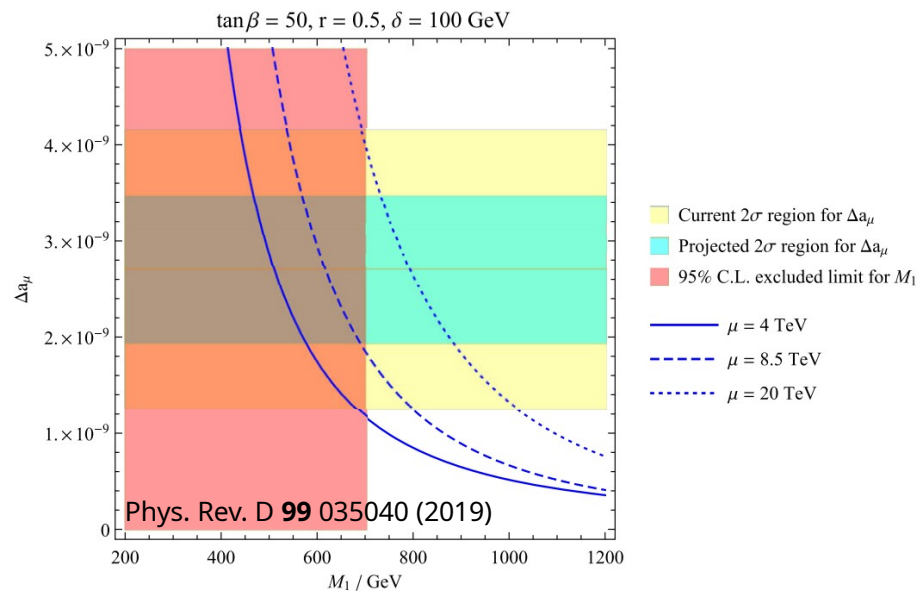
Many models can explain the anomaly and evade other constraints (dark matter, LHC limits, ...)

- TeV leptiquarks, Z', ALPs, SUSY, 2-Higgs doublets, ...

Chirality-flipping, flavour conserving

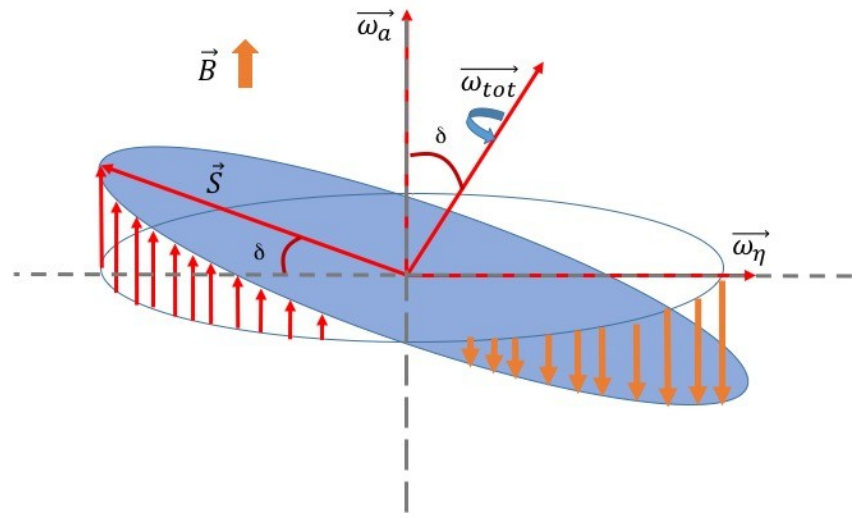
$$\Delta a_\mu^{\text{BSM}} = C_{\text{BSM}} \left(\frac{m_\mu}{M_{\text{BSM}}} \right)^2$$

arXiv:2104.03691



Search for a muon electric dipole moment:

- ~zero in the Standard Model
- possible source of CP-violation



Tilts precession plane towards center of ring

- vertical oscillation, 90° out of phase with a_μ
- tracking detectors key to sensitivity

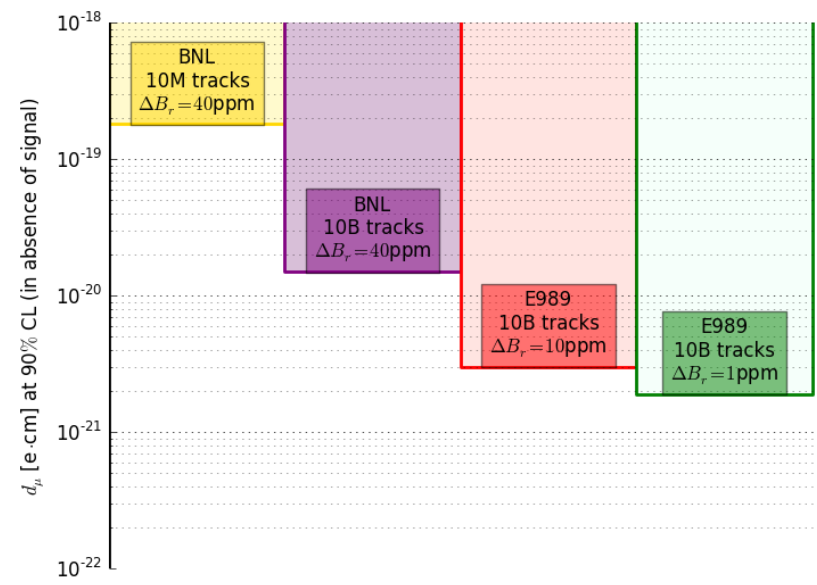
$$\vec{d} = \eta \left(\frac{Qe}{2mc} \right) \vec{s}$$

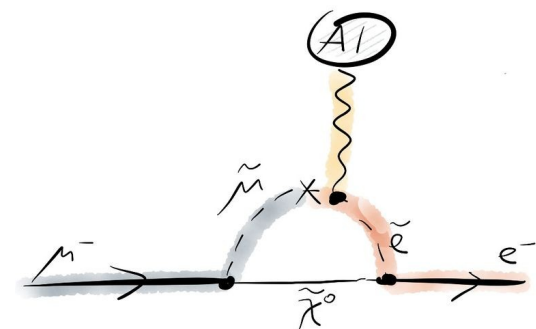
World's best limit:

$$|d_\mu| > 1.9 \times 10^{-19} \text{ e.cm (BNL)}$$

Target:

$$|d_\mu| \approx 1.9 \times 10^{-21} \text{ e.cm}$$



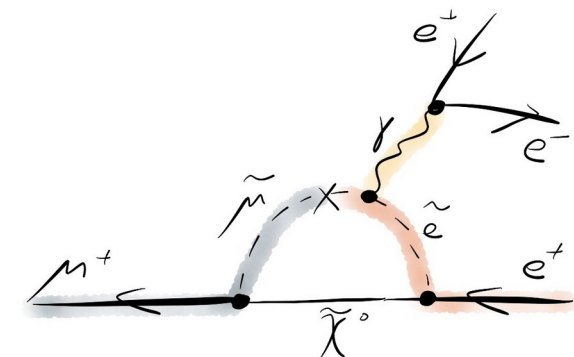


Charged lepton flavour violation:

- again, \sim zero in the Standard Model
- predicted by many BSM models (particularly leptogenesis)

Mu2e:

- follows g-2 at Fermilab
- decay of captured muons



Mu3e:

- data taking in 2023 @ PSI
- full reconstruct $\mu \rightarrow e e e$ final state

Aim to push limits on Br to 10^{-17} (factor $\times 10^4$)

- mass scales up to 10,000 TeV

The FNAL Muon g-2 experiment measured a_μ to 0.46 ppm

- result consistent with BNL
- combined result differs from SM by 4.2σ
- statistics limited, based on 6% of target stats

Upgrades since Run-1:

- Replaced damaged resistors, reducing C_{pa}
- Higher kicker voltage to center beam radially
- Thermal magnet insulation & hall cooling improve field stability

EDM search & CLFV: complementary information on any BSM scenario

