

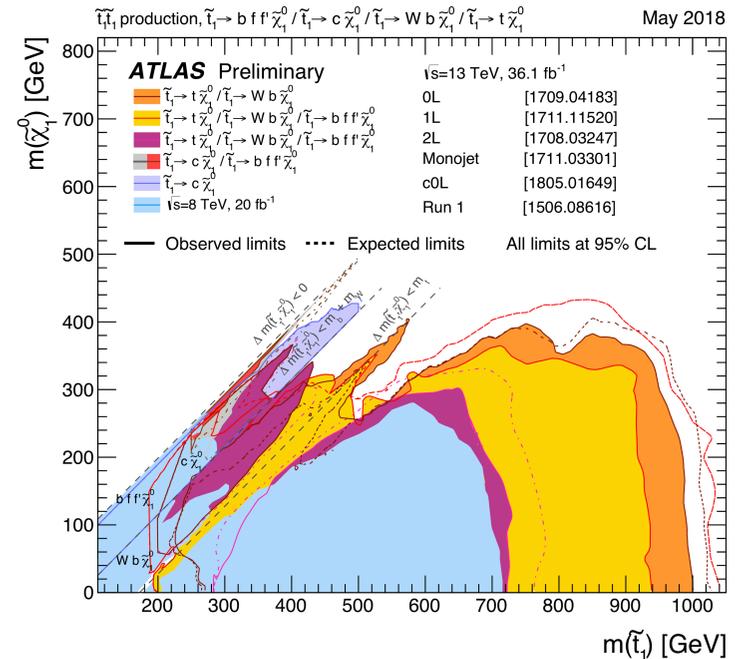
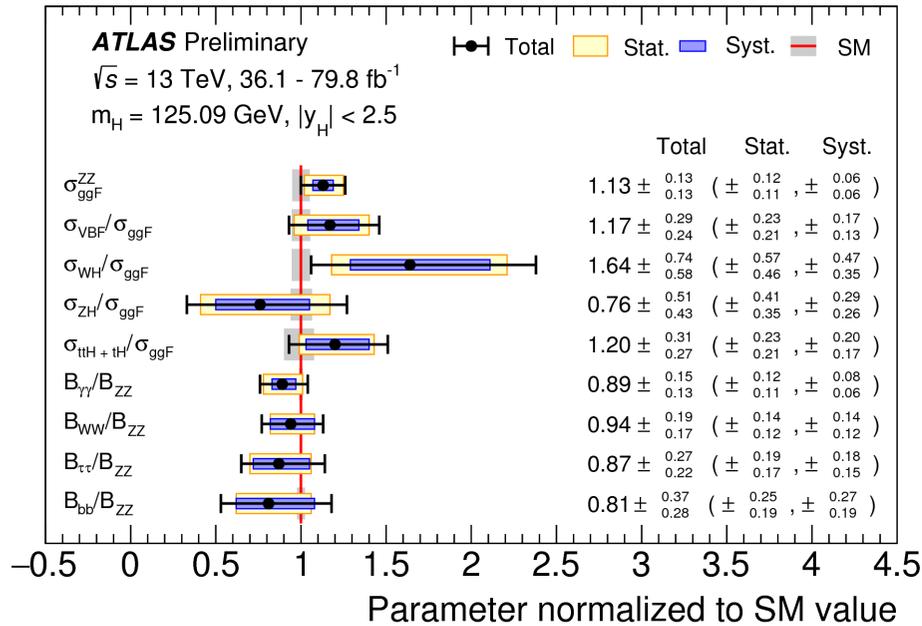
# Dual Readout Calorimetry

## For future $e^+e^-$ colliders

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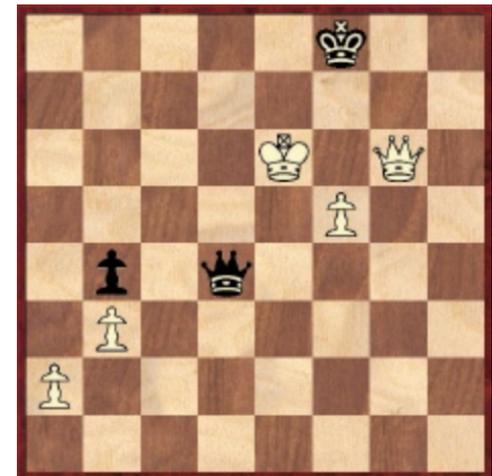
Iacopo Vivarelli  
University of Sussex

# Particle physics in stalemate?



- Did not checkmate fundamental laws of nature
- ... but maybe there is no next move that can be done?

Black to move



# Not really...

---

- Next move of experimentalists is obvious

Put the Higgs under the microscope (a  $e^+e^-$  Higgs factory)



# e<sup>+</sup>e<sup>-</sup> Higgs factories on the table

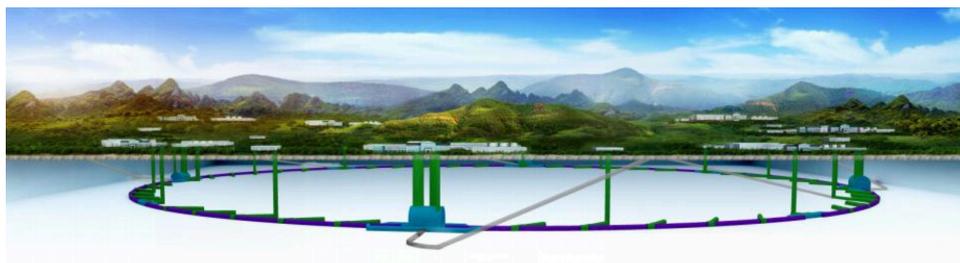
Circular

Linear



China

$$\sqrt{s} = 90 - 240 \text{ GeV}$$



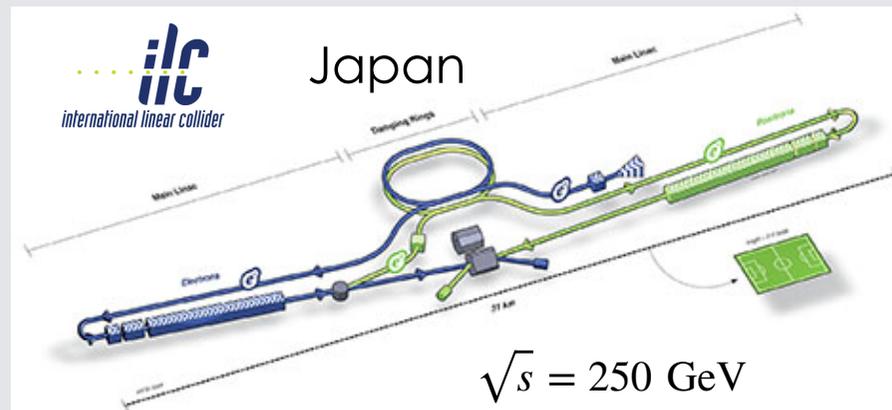
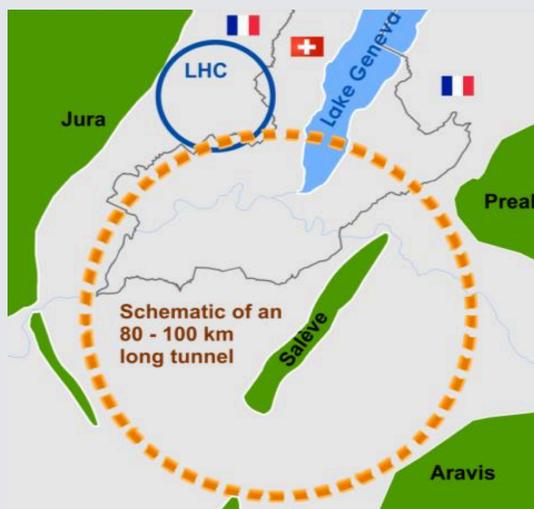
CERN

$$\sqrt{s} = 380 - 3000 \text{ GeV}$$



CERN

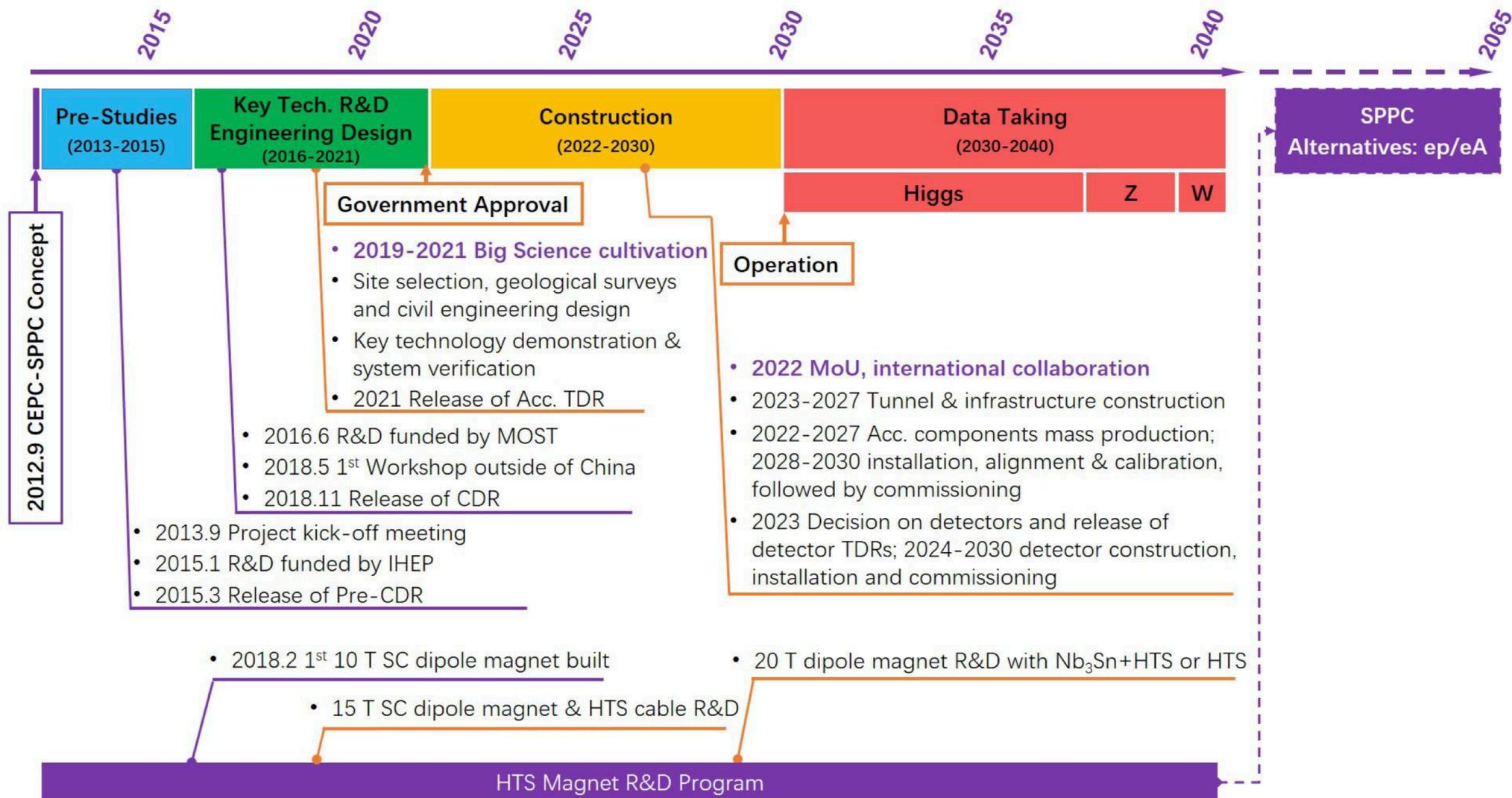
$$\sqrt{s} = 90 - 375 \text{ GeV}$$



$$\sqrt{s} = 250 \text{ GeV}$$

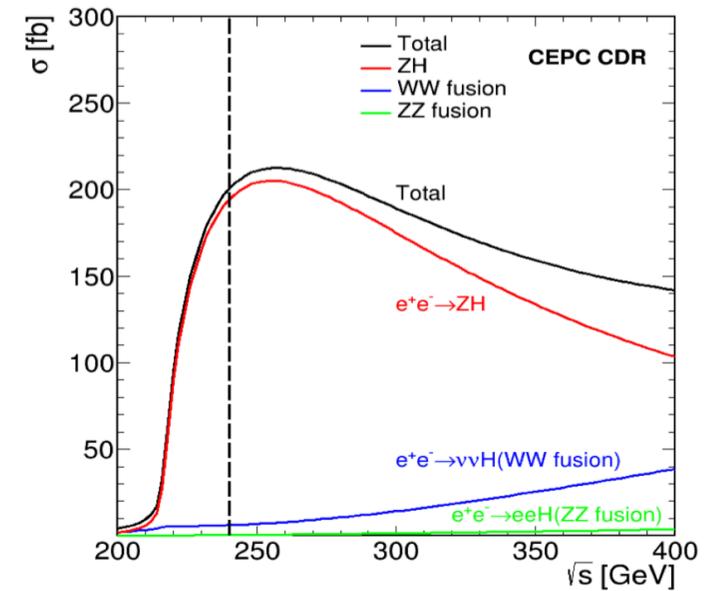
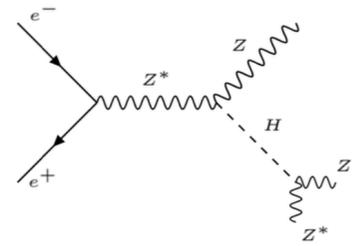
# Far future?

## CEPC Project Timeline

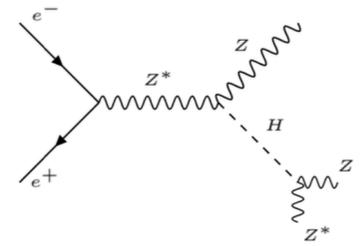


# $e^+e^-$ at ZH threshold

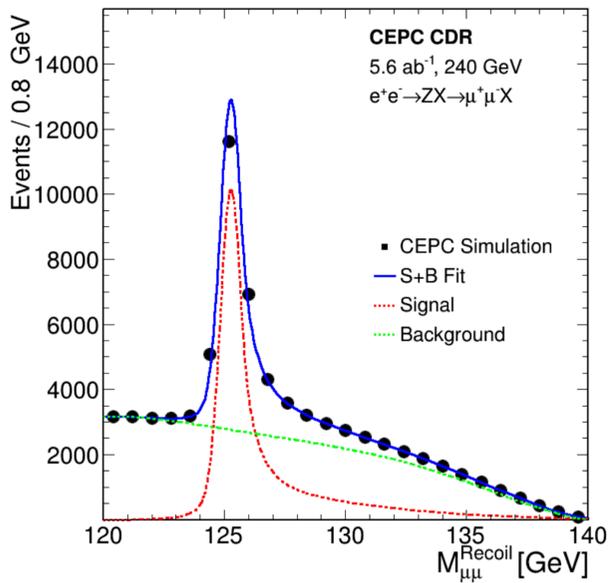
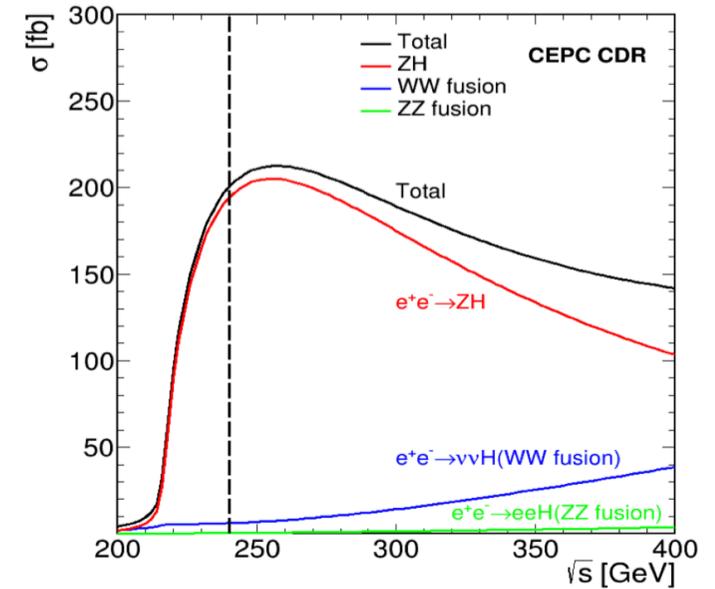
- For CEPC (but similar number and performance for FCC):
  - **5.6 ab<sup>-1</sup>** translate in **a million ZH events**



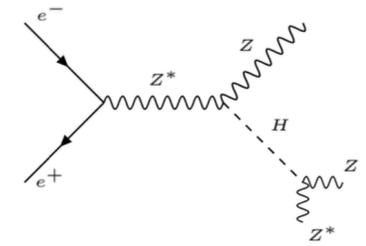
# $e^+e^-$ at ZH threshold



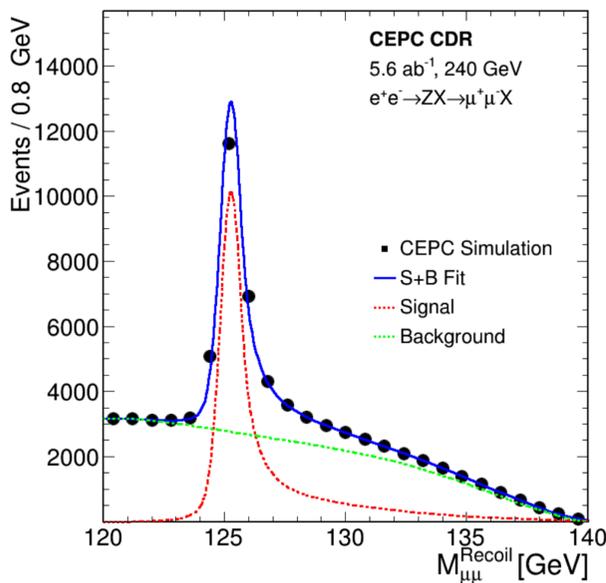
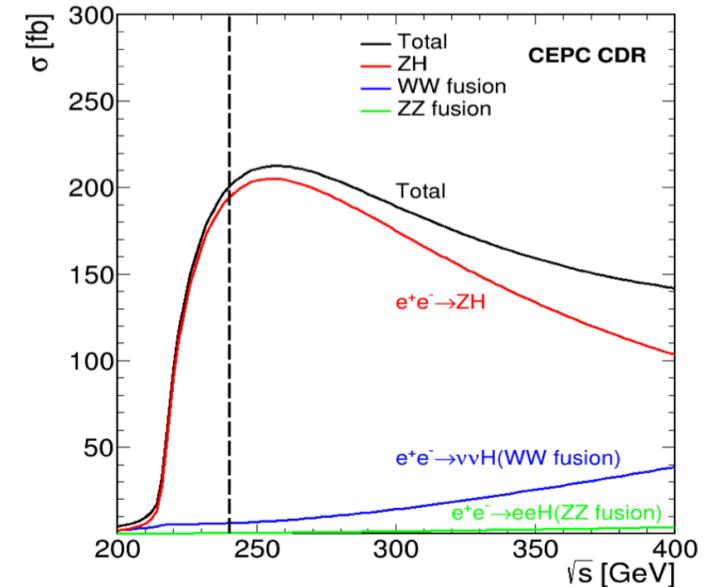
- For CEPC (but similar number and performance for FCC):
  - **$5.6 \text{ ab}^{-1}$**  translate in **a million ZH events**
  - Higgs boson tagging gives unique access to  $e^+e^- \rightarrow ZH$  production cross section



# $e^+e^-$ at ZH threshold



- For CEPC (but similar number and performance for FCC):
  - **$5.6 \text{ ab}^{-1}$**  translate in **a million ZH events**
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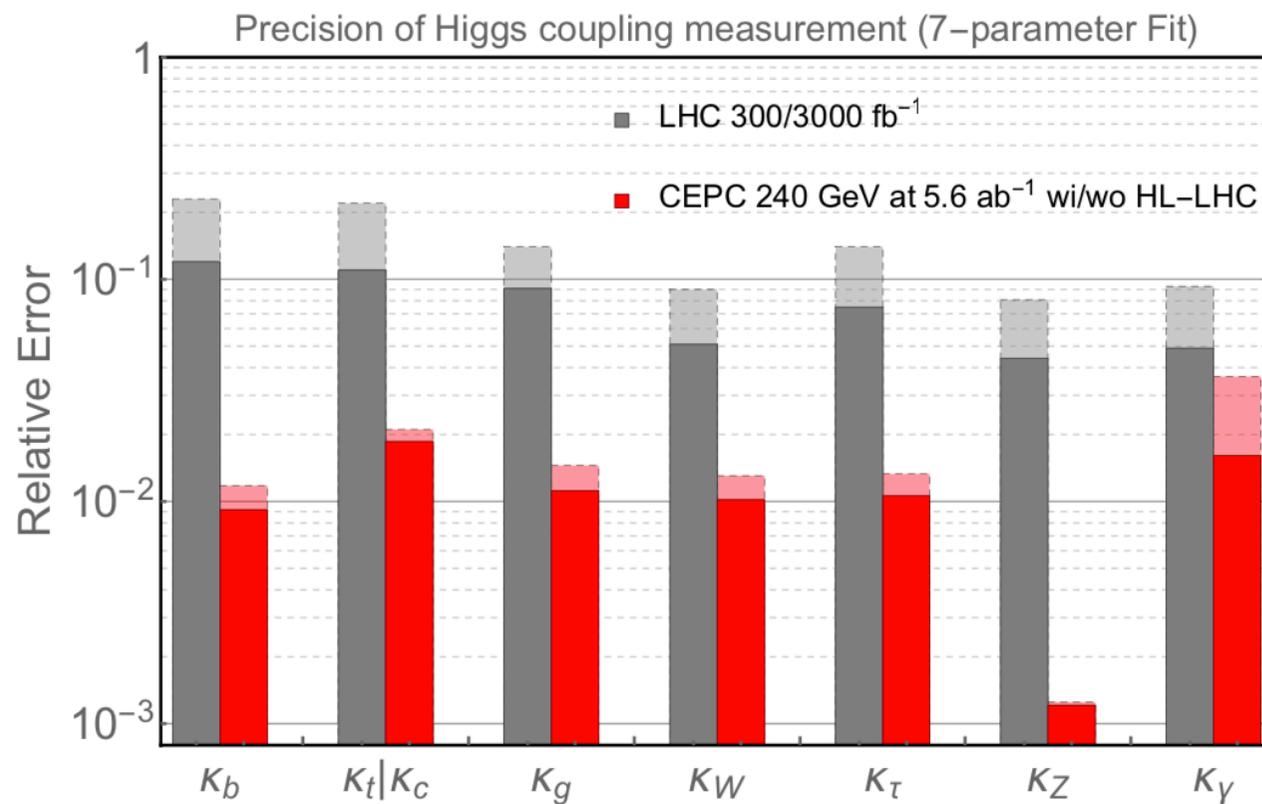


- **3% model-independent** Higgs boson width measurement from

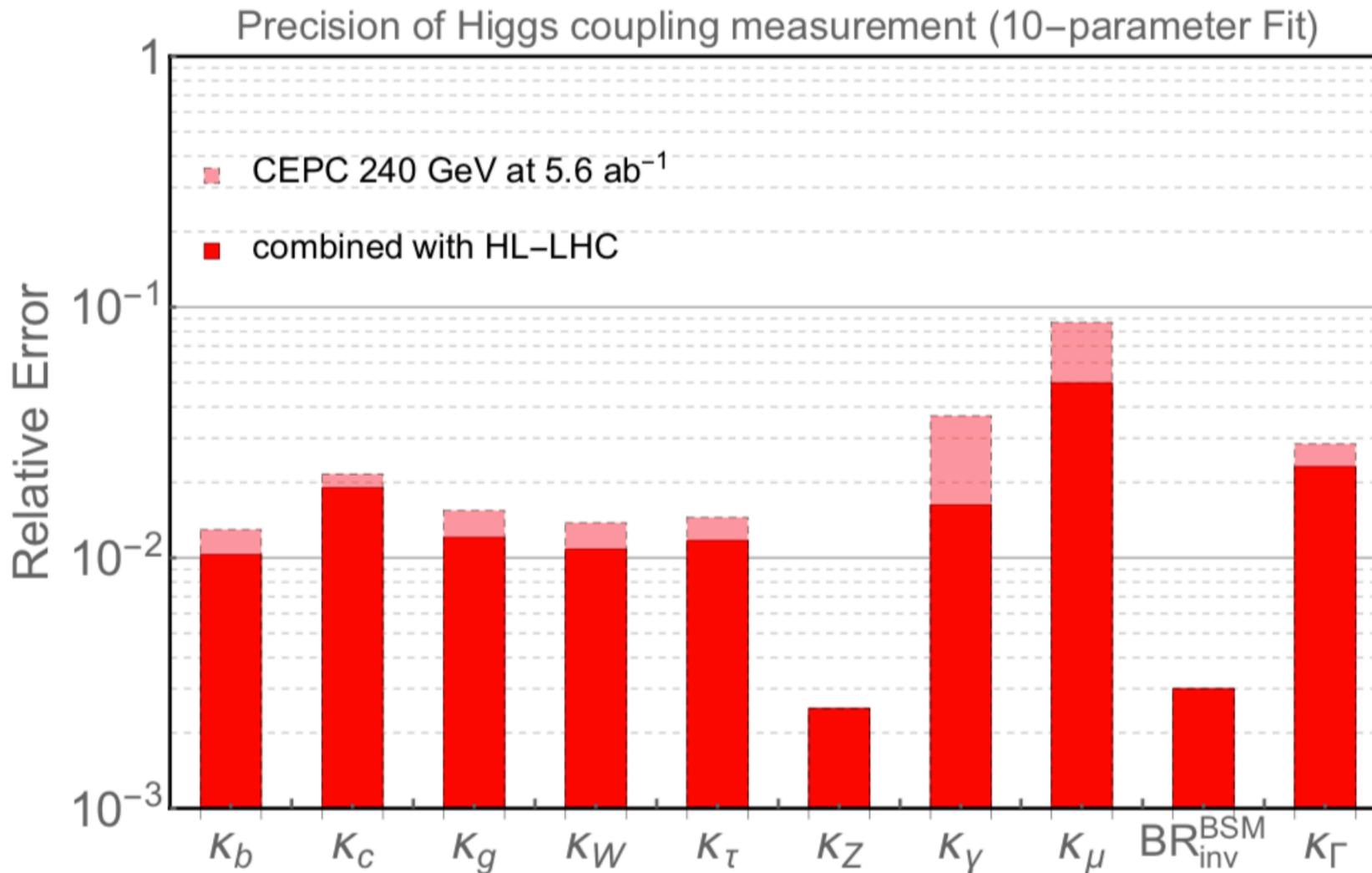
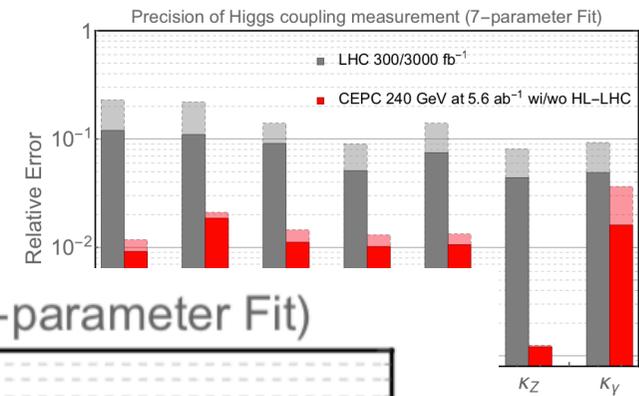
$$\Gamma_H = \frac{\Gamma(H \rightarrow ZZ^*)}{\text{BR}(H \rightarrow ZZ^*)} \propto \frac{\sigma(ZH)}{\text{BR}(H \rightarrow ZZ^*)}$$

$$\Gamma_H = \frac{\Gamma(H \rightarrow b\bar{b})}{\text{BR}(H \rightarrow b\bar{b})} = \frac{\Gamma(H \rightarrow WW^*)}{\text{BR}(H \rightarrow WW^*)} \propto \frac{\sigma(\nu\bar{\nu}H)}{\text{BR}(H \rightarrow WW^*)}$$

# Higgs boson sensitivity



# Higgs boson sensitivity



# Physics requirements for calorimetry

- **Precision physics** at  $e^+e^-$  collider calls for **high-resolution hadronic calorimetry**

Physics Process	Measured Quantity	Critical Detector	Required Performance
$ZH \rightarrow \ell^+ \ell^- X$	Higgs mass, cross section	Tracker	$\Delta(1/p_T) \sim 2 \times 10^{-5}$
$H \rightarrow \mu^+ \mu^-$	$\text{BR}(H \rightarrow \mu^+ \mu^-)$		$\oplus 1 \times 10^{-3} / (p_T \sin \theta)$
$H \rightarrow b\bar{b}, c\bar{c}, gg$	$\text{BR}(H \rightarrow b\bar{b}, c\bar{c}, gg)$	Vertex	$\sigma_{r\phi} \sim 5 \oplus 10 / (p \sin^{3/2} \theta) \mu\text{m}$
$H \rightarrow q\bar{q}, VV$	$\text{BR}(H \rightarrow q\bar{q}, VV)$	ECAL, HCAL	$\sigma_E^{\text{jet}} / E \sim 3 - 4\%$
$H \rightarrow \gamma\gamma$	$\text{BR}(H \rightarrow \gamma\gamma)$	ECAL	$\sigma_E \sim 16\% / \sqrt{E} \oplus 1\% (\text{GeV})$

# Physics requirements for calorimetry

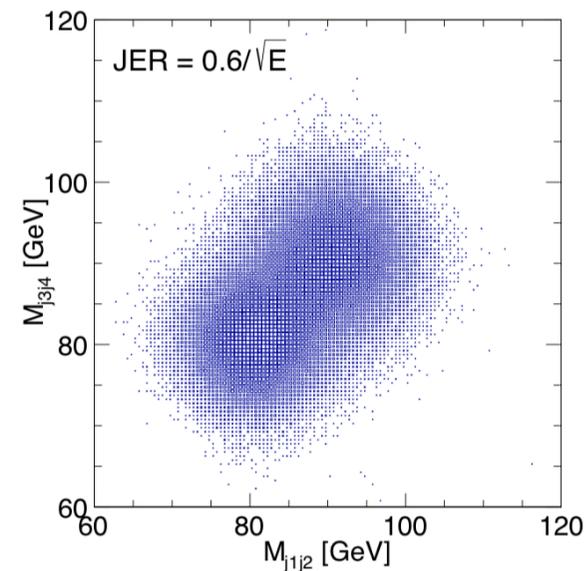
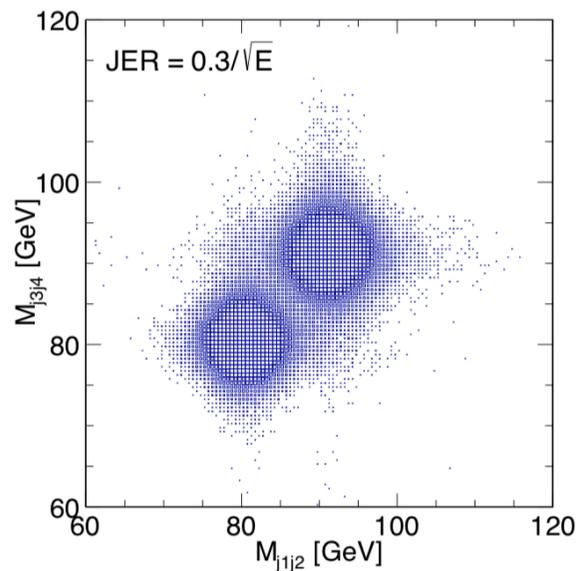
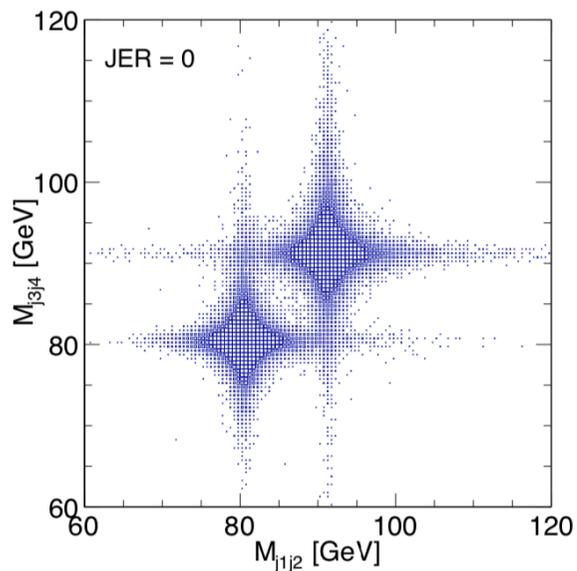
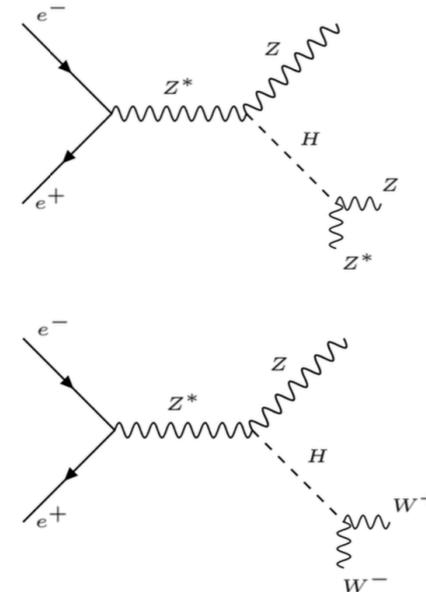
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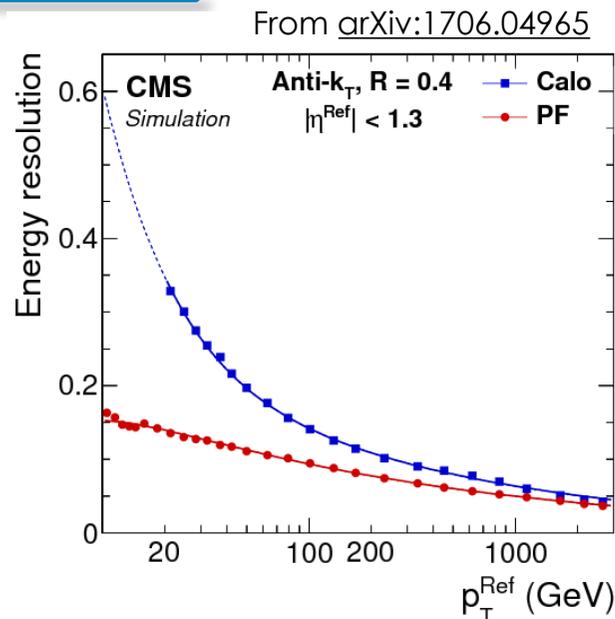
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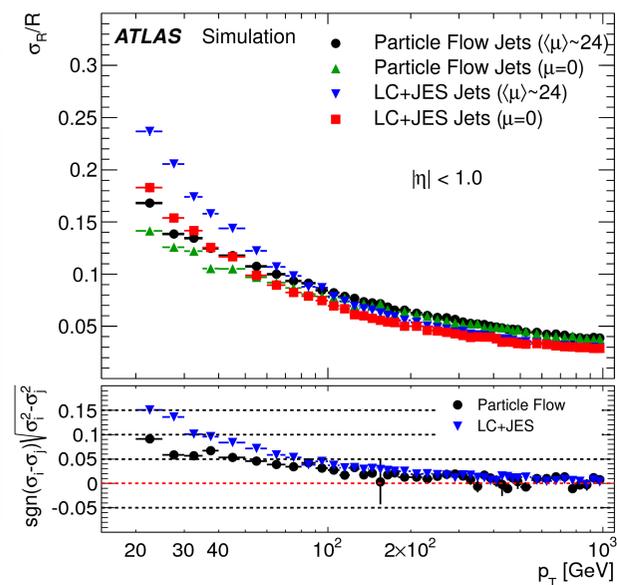


# Particle flow approach

- Basic idea: tracking wins at low particle energy
  - Use **tracks** to measure **charged particles** in the shower



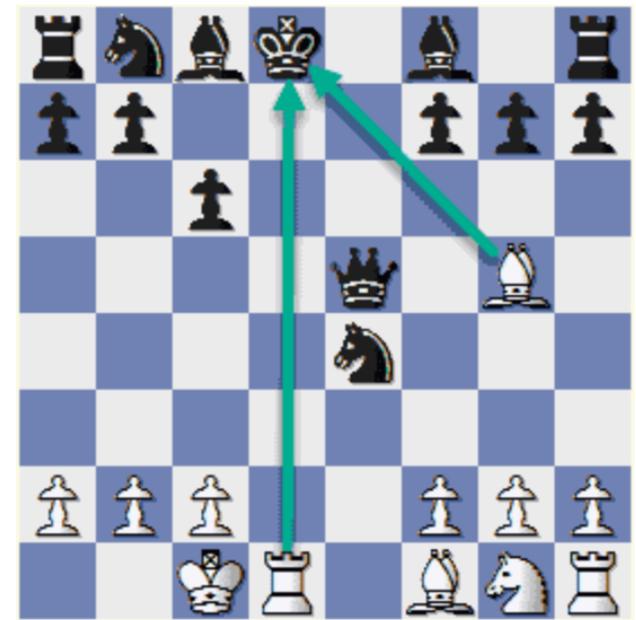
From *Eur. Phys. J C*77 (2017) 466



From [https://warwick.ac.uk/fac/sci/physics/staff/academic/boyd/warwick\\_week/detector\\_physics/warwick\\_lecture\\_calorimetry.pdf](https://warwick.ac.uk/fac/sci/physics/staff/academic/boyd/warwick_week/detector_physics/warwick_lecture_calorimetry.pdf)

# A different approach

- Attacking the problem from **two sides** always gives opportunities to learn
- Dual readout calorimetry:
  - Excellent **native electromagnetic and hadronic** calorimeter resolution
  - Excellent lateral granularity

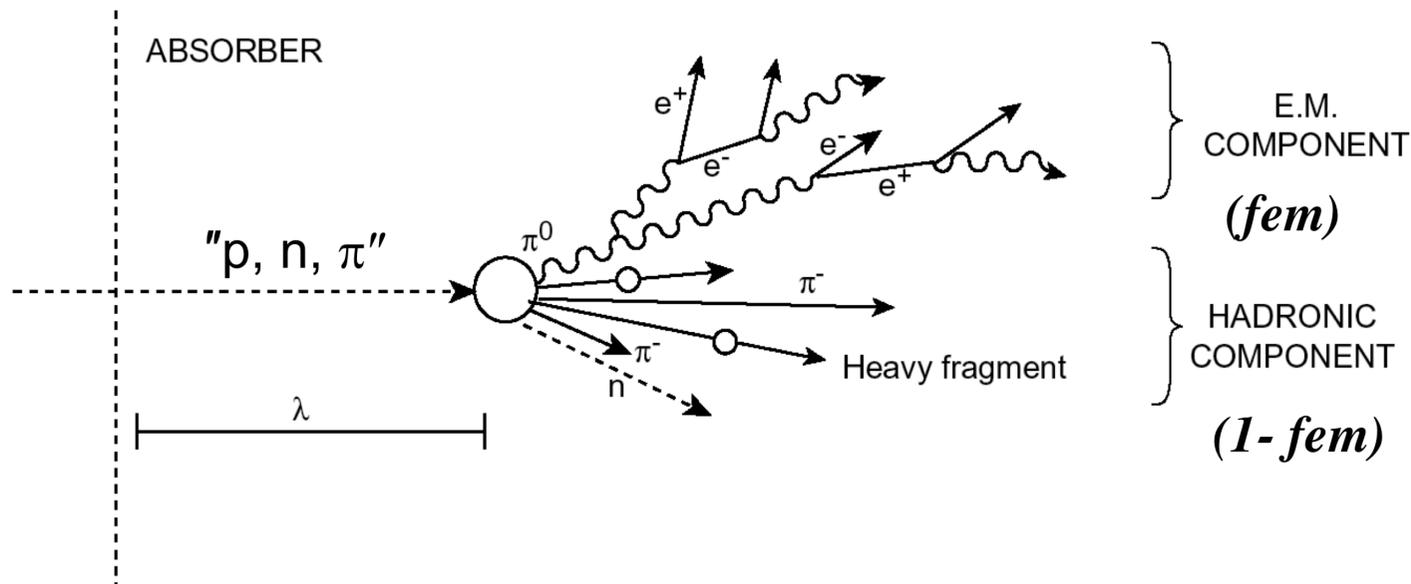


Mate in 2

# Hadronic Calorimetry - primer

# Hadronic Calorimetry - primer

- Hadronic shower:
  - Driven by a relatively **small number** of strong interactions with **nuclei**
  - **Strong** intrinsic interaction intensity, but **small targets** (scale to bear in mind  $1 \text{ fm} = 10^{-15} \text{ m}$ )
  - $\pi^0, \eta^0$  production leads to **EM component within hadronic shower**



# Hadronic showers

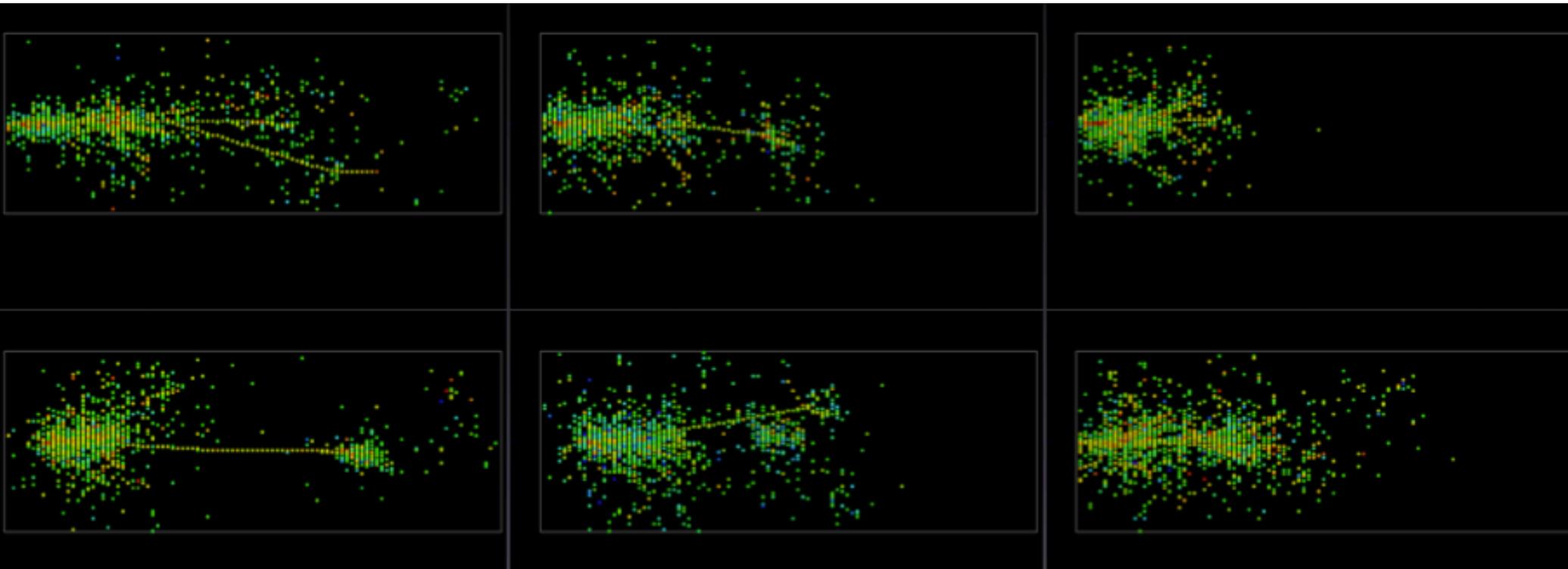
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- Lots of different physics processes at work in a hadronic shower:
  - EM component: **large fluctuations** in number and energy of  $\pi^0, \eta^0$
  - HAD component:
    - **Ionisation** from charged hadrons
    - **Nuclear remnants** (fission, knock-off)
    - Delayed photons
    - **Invisible component** (nuclei breakup)

# Hadronic showers

---

- Large **event-by-event** fluctuations in shape (and energy) deposit
- **Charged hadrons propagate the shower** on large scale ( $\lambda_1$ ), **local** EM showers from  $\pi^0, \eta^0$ .



# Why is hadronic calorimetry challenging?

Typically\* the calorimeter response to the electromagnetic (e) and hadronic (h) components is different (because of invisible energy)

\* This is actually true for non-compensating calorimeter (the vast majority of those in use nowadays)

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The fraction of energy in EM component and HAD component is energy dependent with large fluctuations

# Why is hadronic calorimetry challenging?

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\*This is actually true for non-compensating calorimeter (the vast majority of those in use nowadays)

The fraction of energy in EM component and HAD component is energy dependent with large fluctuations

The calorimeter response to hadrons is energy dependent and fluctuates a lot

# Calorimeter response to hadrons

- Typically calorimeters are calibrated to the **EM scale**
  - For example: you shoot 20 GeV electrons and want to read 20 GeV
  - Say  $r$  is the energy deposit in active material in the calorimeter. Then you choose  $k$  such that

$$E_e = ker$$

- Then the response to a hadron is

$$E_h = k(ef_{em} + h(1 - f_{em}))r$$

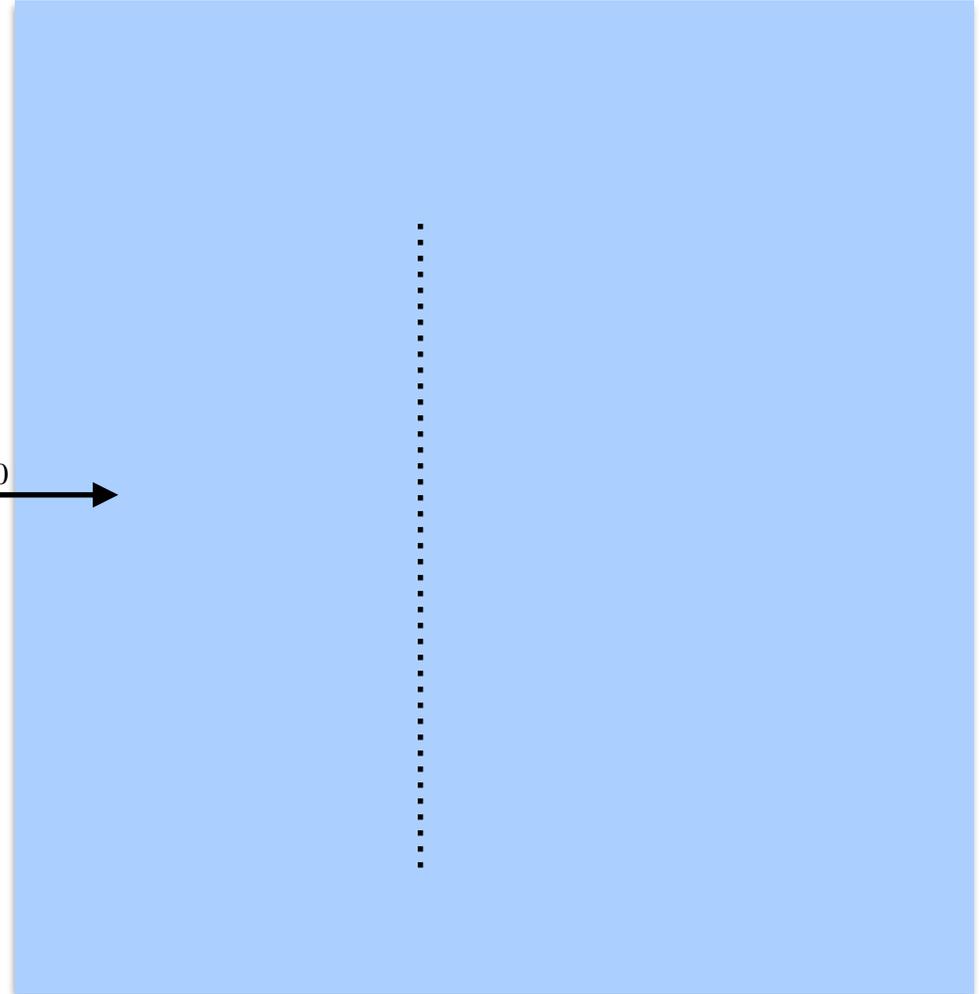
- and with respect to that of an electron of the same energy

$$E_h = E_e(f_{em} + \frac{h}{e}(1 - f_{em}))$$

# f<sub>em</sub> energy dependence

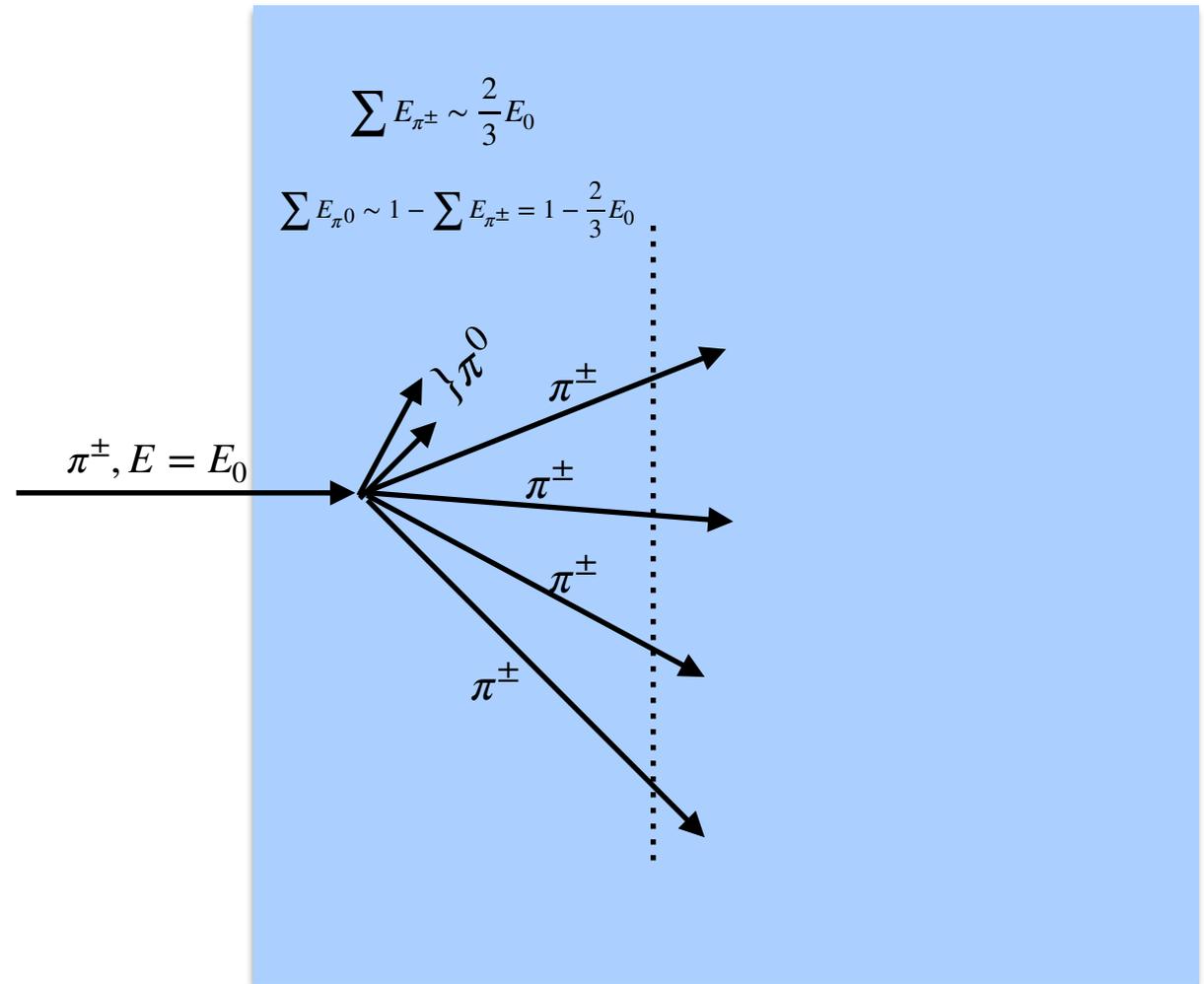
- Simple model: **only pions** are produced at each interaction, **respecting isospin symmetry**
- Then the math is:

$$\pi^\pm, E = E_0 \longrightarrow$$



# f<sub>em</sub> energy dependence

- Simple model: **only pions** are produced at each interaction, **respecting isospin symmetry**
- Then the math is:





# $f_{em}$ energy dependence

- Simple model: **only pions** are produced at each interaction, **respecting isospin symmetry**

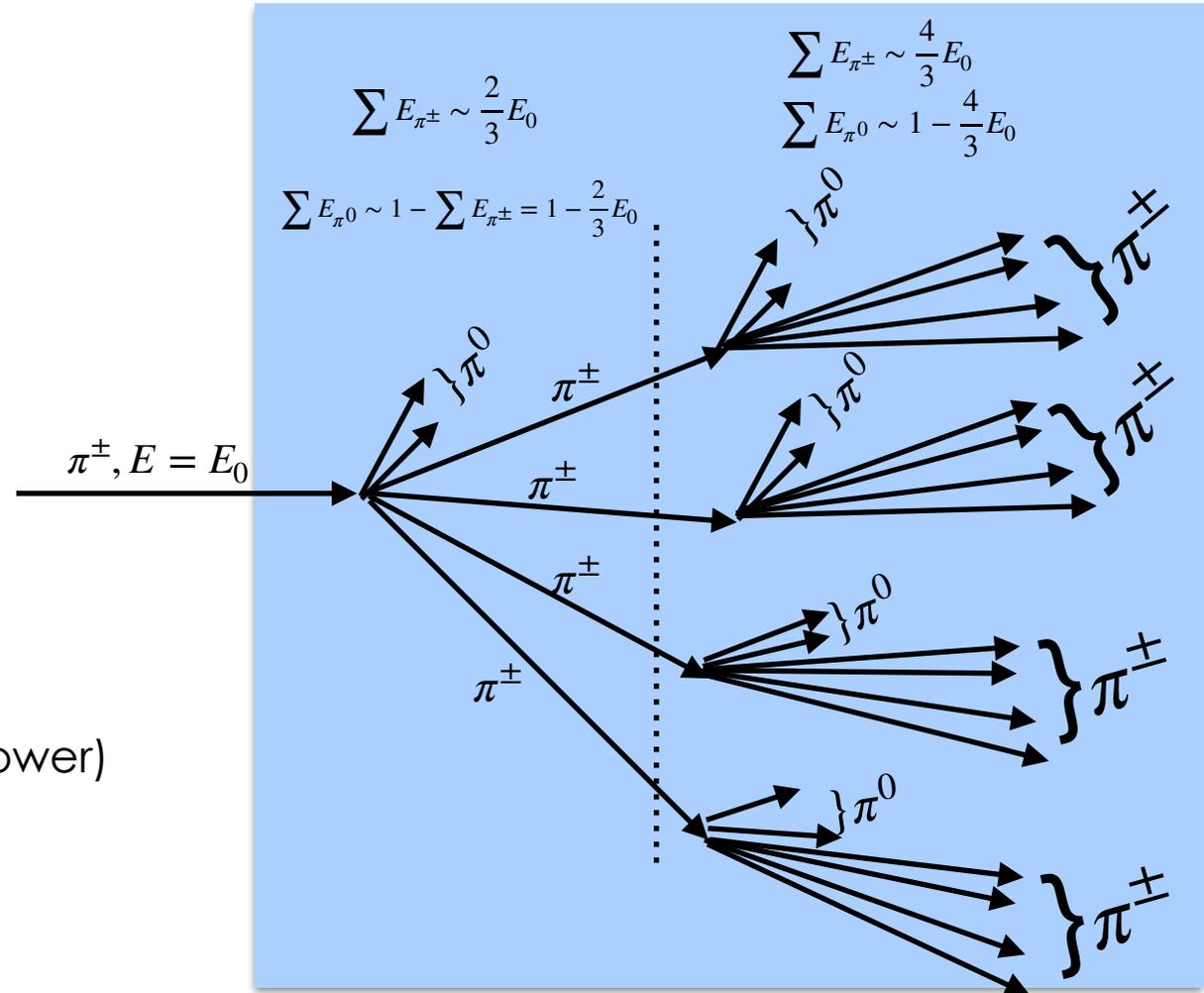
- Then the math is:

$$f_{em} = \frac{E_{em}}{E} = 1 - \left(\frac{2}{3}\right)^n$$

$$\left(\frac{2}{3}\right)^n E = E_{th} \quad (\text{Condition to stop shower})$$

$$n = p \ln \frac{E}{E_{th}} \quad (\text{for some number } p)$$

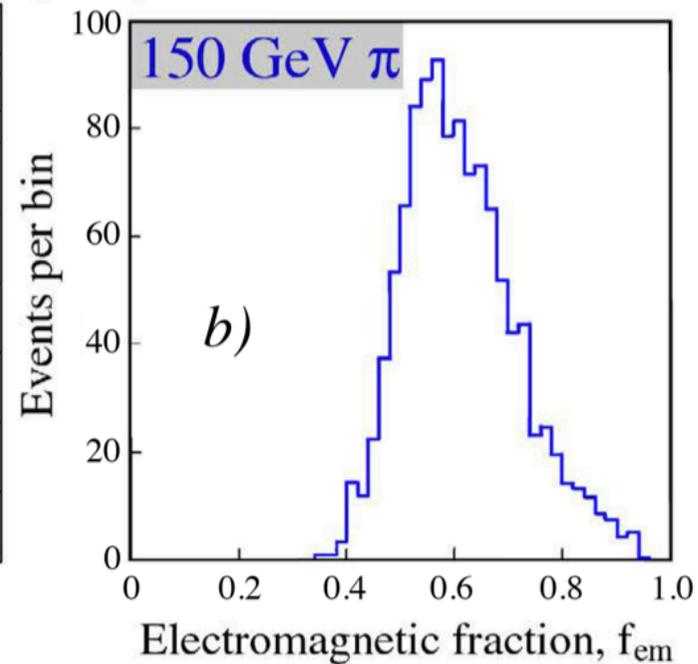
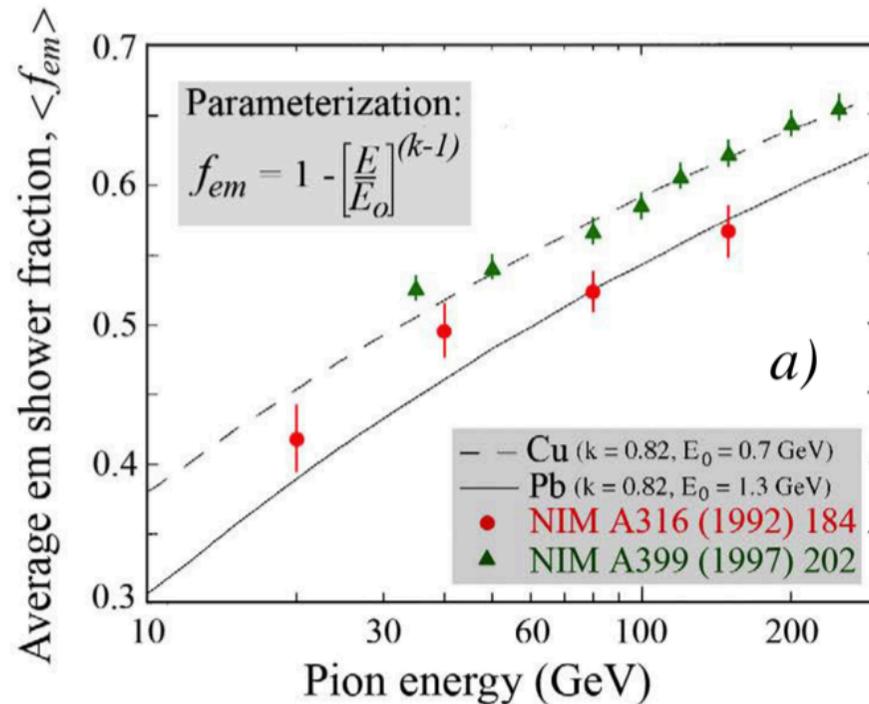
$$f_{em} = 1 - \left(\left(\frac{2}{3}\right)^{\ln \frac{E}{E_{th}}}\right)^p = 1 - \left(\frac{E}{E_{th}}\right)^{k-1} \quad (\text{for some number } k)$$



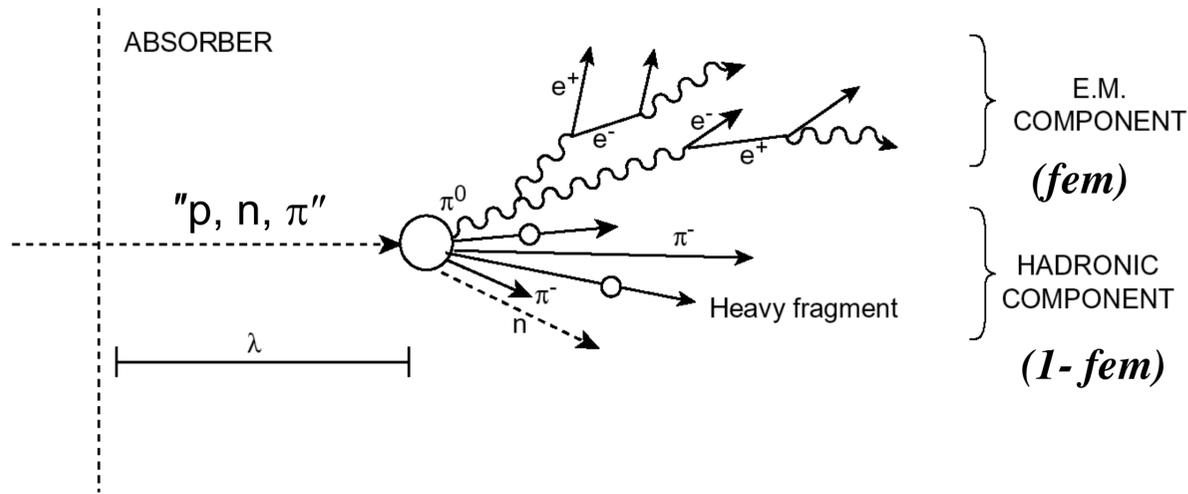
# $f_{em}$ energy dependence

$$f_{em} = 1 - \left( \frac{E}{E_0} \right)^{k-1}$$

- $f_{em}$  depends on the **incoming particle energy**
- Fluctuations in  $f_{em}$  are **large and non-poissonian**

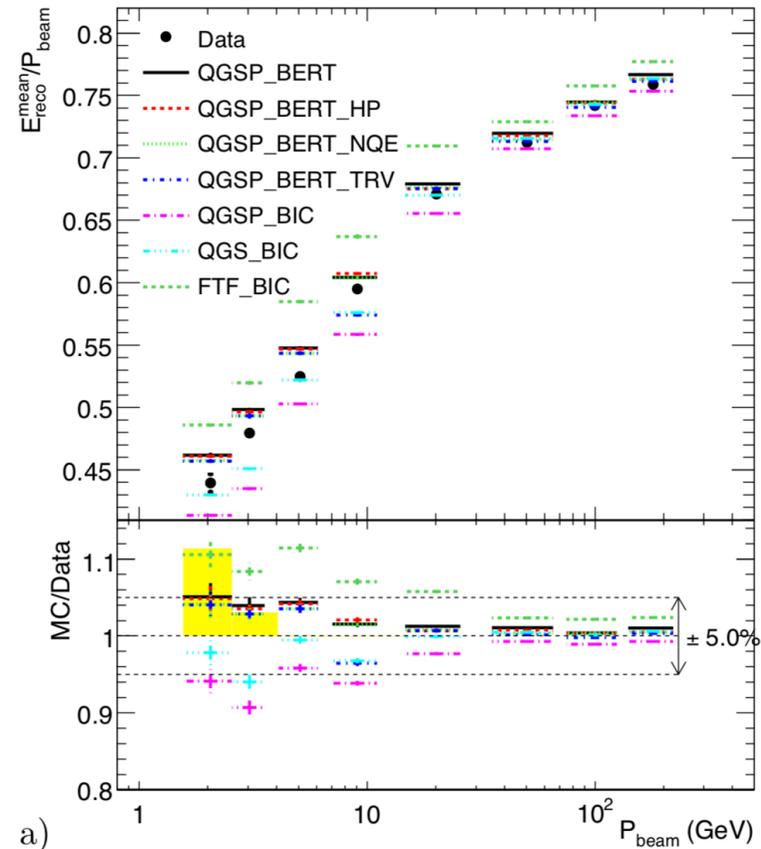


# The curse of hadronic calorimetry



- **Non-compensating calorimeters:** response to em part different from that to non-em part.  $h/e < 1$
- $\langle f_{em} \rangle$  energy dependent  $\Rightarrow$  **Non-linear calorimeter response** to hadrons

ATL-CAL-PUB-2010-001

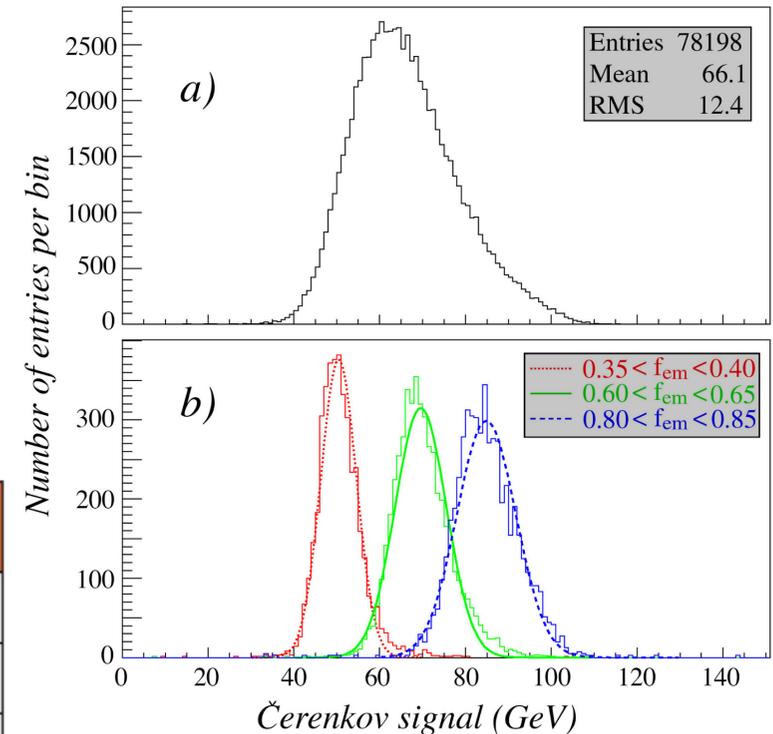


$$E_{meas} = E \left( f_{em} + \frac{h}{e} (1 - f_{em}) \right)$$

# The curse of hadron calorimetry (2)

- Event-by-event  $f_{em}$  fluctuations dominate the hadronic calorimeter resolution

Experiment	Detector	Absorber material	$e/h$	Energy resolution (E in GeV)
UA1 C-Modul	Scintillator	Fe	$\approx 1.4$	$80\%/\sqrt{E}$
ZEUS	Scintillator	Pb	$\approx 1.0$	$34\%/\sqrt{E}$
WA78	Scintillator	U	0.8	$52\%/\sqrt{E} \oplus 2.6\%^*$
D0	liquid Ar	U	1.11	$48\%/\sqrt{E} \oplus 5\%^*$
H1	liquid Ar	Pb/Cu	$\leq 1.025^*$	$45\%/\sqrt{E} \oplus 1.6\%$
CMS	Scintillator	Brass (70% Cu / 30% Zn)	$\neq 1$	$100\%/\sqrt{E} \oplus 5\%$
ATLAS (Barrel)	Scintillator	Fe	$\neq 1$	$50\%/\sqrt{E} \oplus 3\%$
ATLAS (Endcap)	liquid Ar	Brass	$\neq 1$	$60\%/\sqrt{E} \oplus 3\%$



N. Akchurin *et al.*, *Nucl. Instr. and Meth.* A537 (2005) 537 .

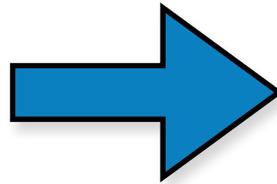
# Dual readout

# Dual readout - the principle

- Suppose I read out two signals, S and C, with different  $h/e$ . Then:

$$E_S = E \left( f_{\text{em}} + \left( \frac{h}{e} \right)_S (1 - f_{\text{em}}) \right)$$

$$E_C = E \left( f_{\text{em}} + \left( \frac{h}{e} \right)_C (1 - f_{\text{em}}) \right)$$



$$f_{\text{em}} = \frac{\left( \frac{h}{e} \right)_C - \left( \frac{h}{e} \right)_S \left( \frac{E_C}{E_S} \right)}{\left( \frac{E_C}{E_S} \right) \left( 1 - \left( \frac{h}{e} \right)_S \right) - \left( 1 - \left( \frac{h}{e} \right)_C \right)}$$

$$E = \frac{(E_S - \chi E_C)}{1 - \chi}$$

$$\chi = \frac{1 - \left( \frac{h}{e} \right)_S}{1 - \left( \frac{h}{e} \right)_C}$$

Depends only on the detector, it can be determined in test beam for example

# Dual readout - the signals

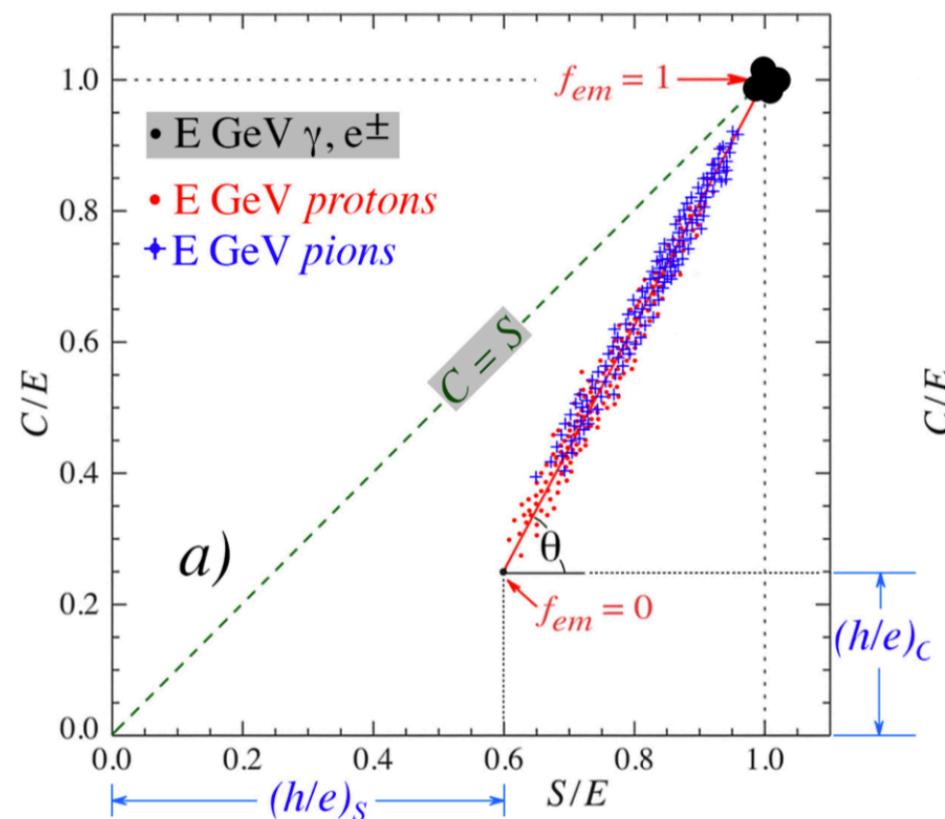
• In practice, this is realised with:

- S: **scintillating fiber** signal measuring  $dE/dx$  of particles. Sensitive to **all shower components** -  $h/e < 1$
- C: **undoped fibres** sensitive to **Cherenkov** signal from relativistic particles in the shower (essentially only EM component) -  $h/e \sim 0.2$

$$\frac{S}{E} = \left( f_{em} + \left( \frac{h}{e} \right)_S (1 - f_{em}) \right)$$

$$\frac{C}{E} = \left( f_{em} + \left( \frac{h}{e} \right)_C (1 - f_{em}) \right)$$

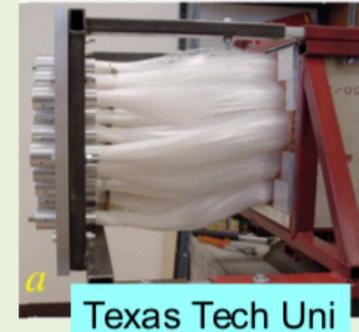
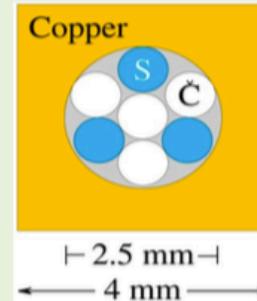
$$\chi = \frac{1 - \left( \frac{h}{e} \right)_S}{1 - \left( \frac{h}{e} \right)_C} = \cot \theta$$



# Dual readout calorimeters (PMT readouts)

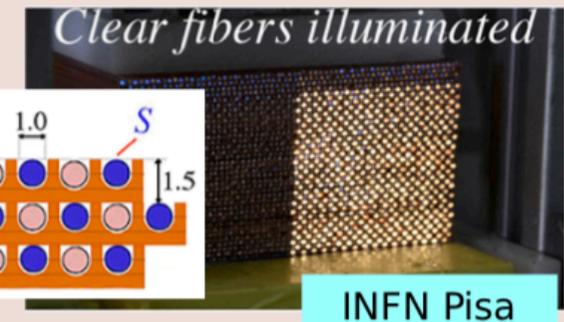
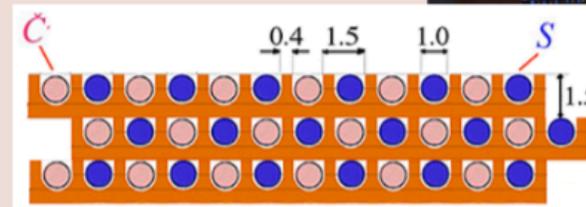
2003  
DREAM

Cu: 19 towers, 2 PMT each  
2m long, 16.2 cm wide  
Sampling fraction: 2%



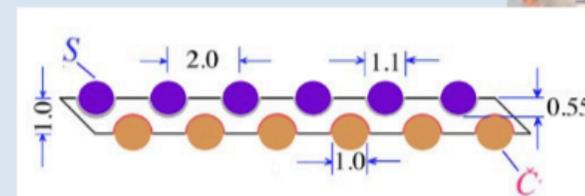
2012  
RD52

Cu, 2 modules  
Each module:  $9.2 \times 9.2 \times 250 \text{ cm}^3$   
Fibers: 1024 S + 1024 C, 8 PMT  
Sampling fraction:  $\sim 4.6\%$   
Depth:  $\sim 10 \lambda_{\text{int}}$

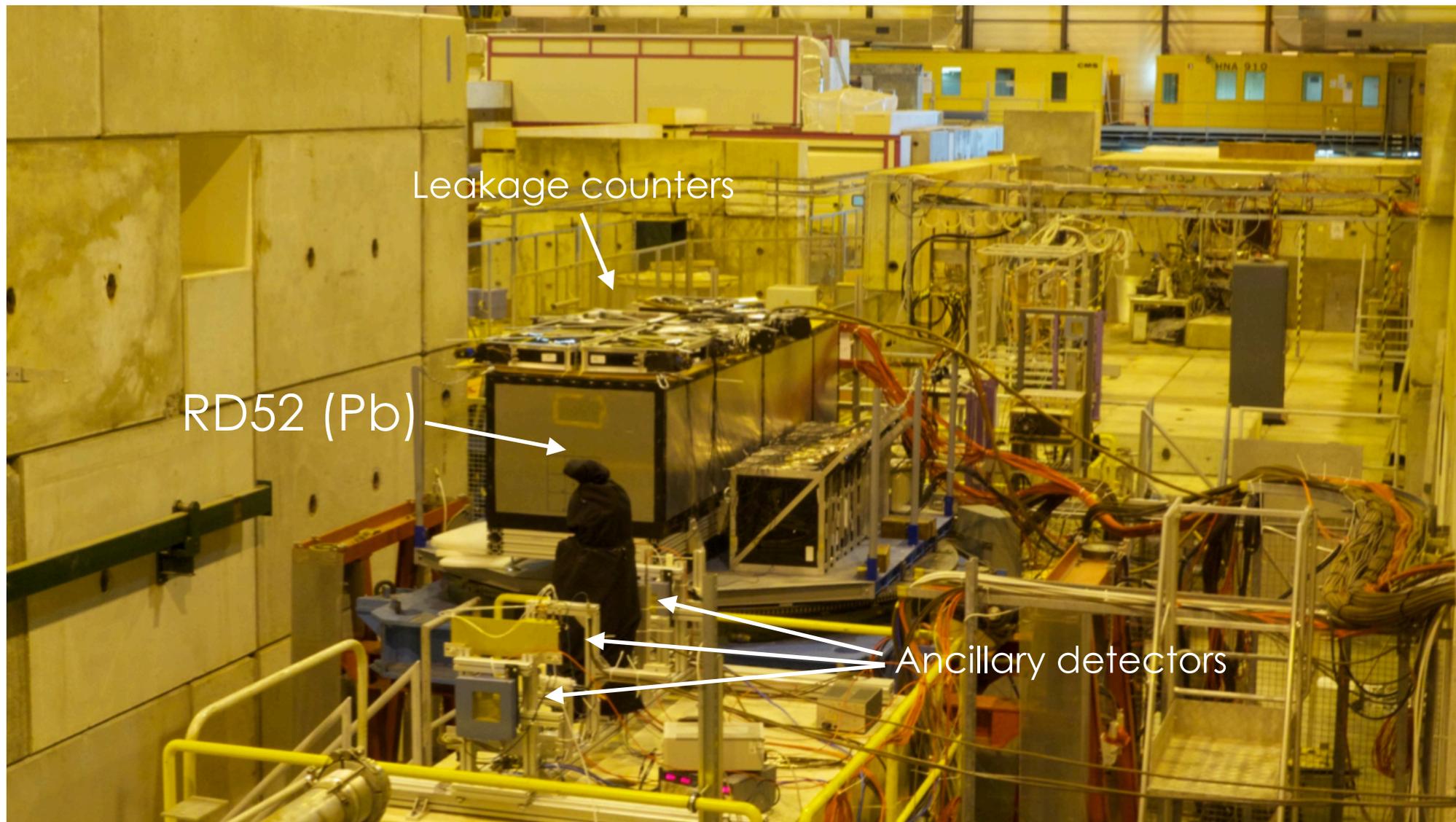


2012  
RD52

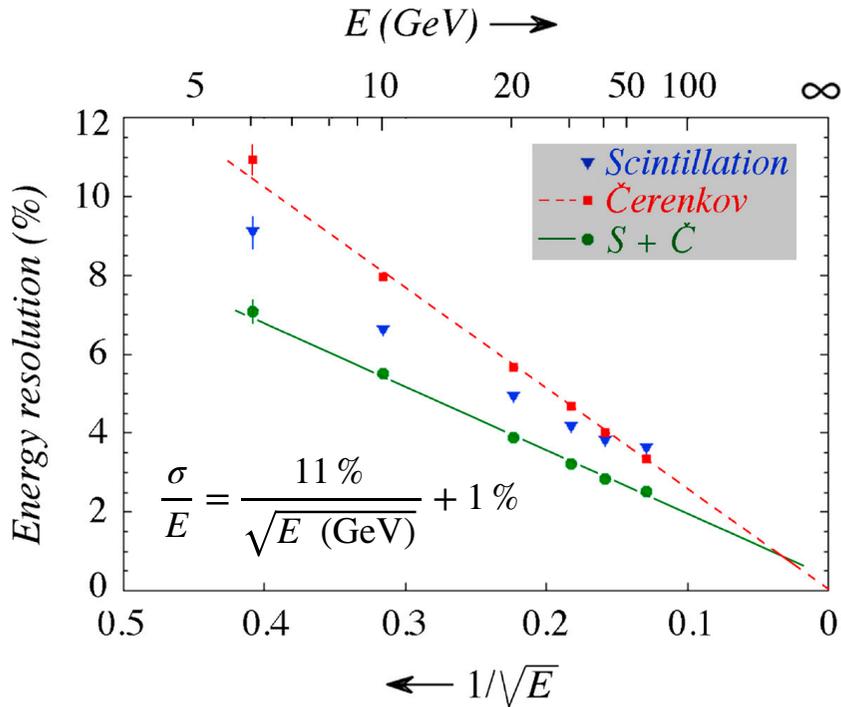
Pb, 9 modules  
Each module:  $9.2 \times 9.2 \times 250 \text{ cm}^3$   
Fibers: 1024 S + 1024 C, 8 PMT  
Sampling fraction:  $\sim 5.3\%$   
Depth:  $\sim 10 \lambda_{\text{int}}$



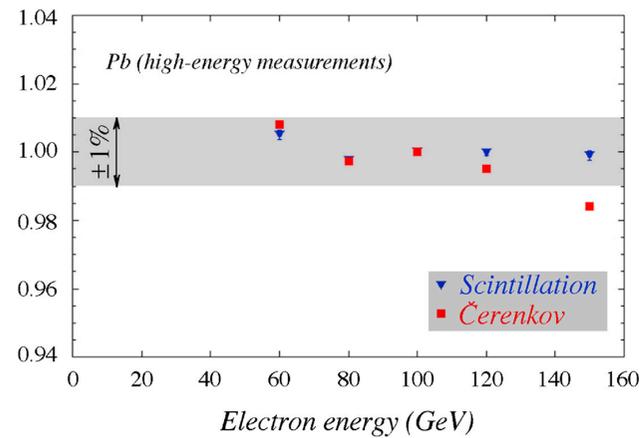
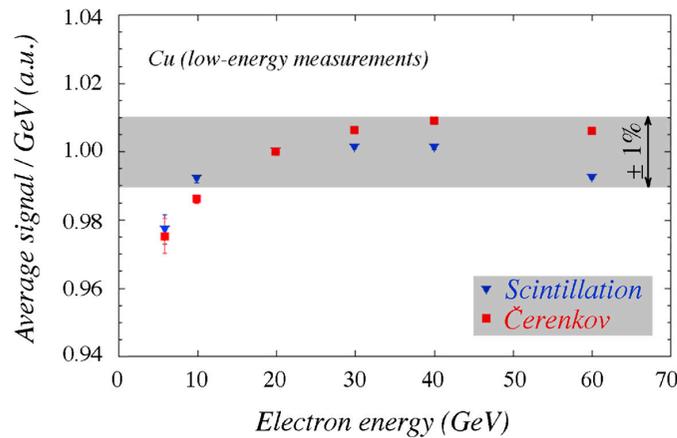
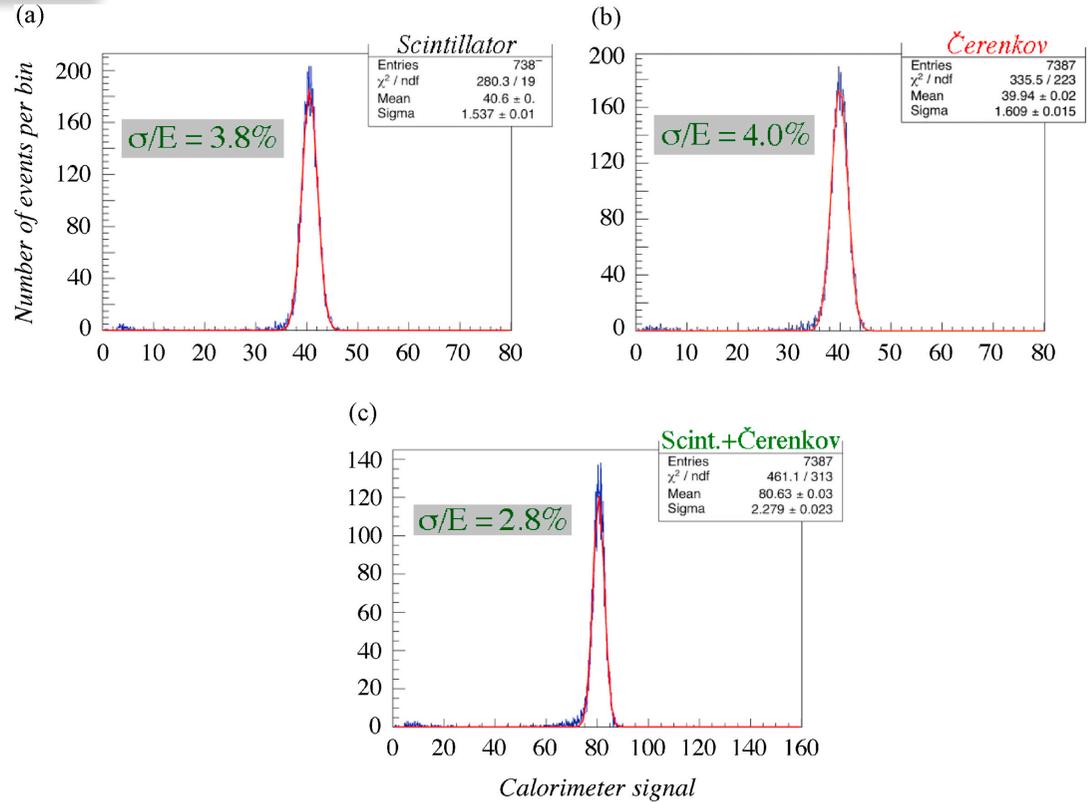
# Dual readout calorimeter at work



# Electron response



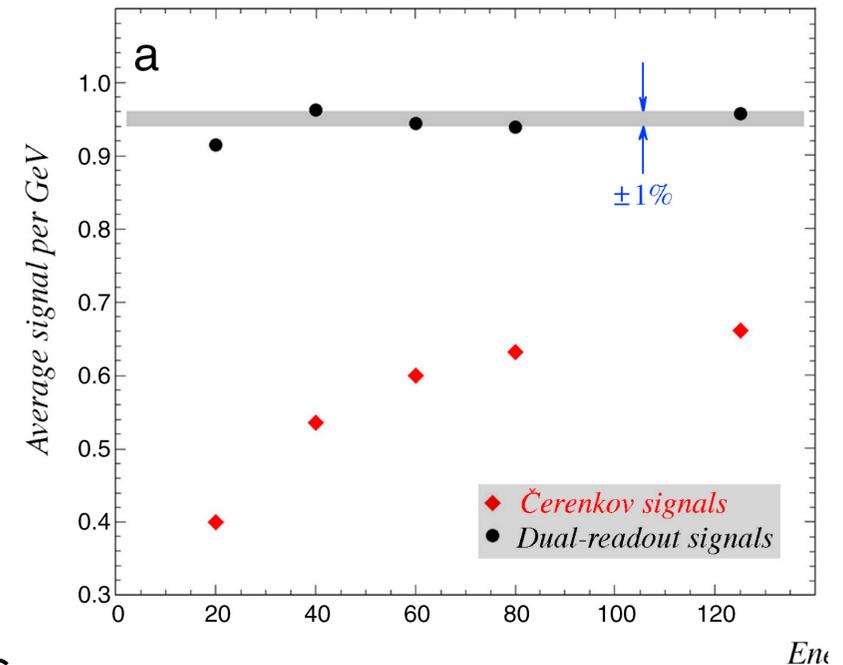
## 40 GeV electrons



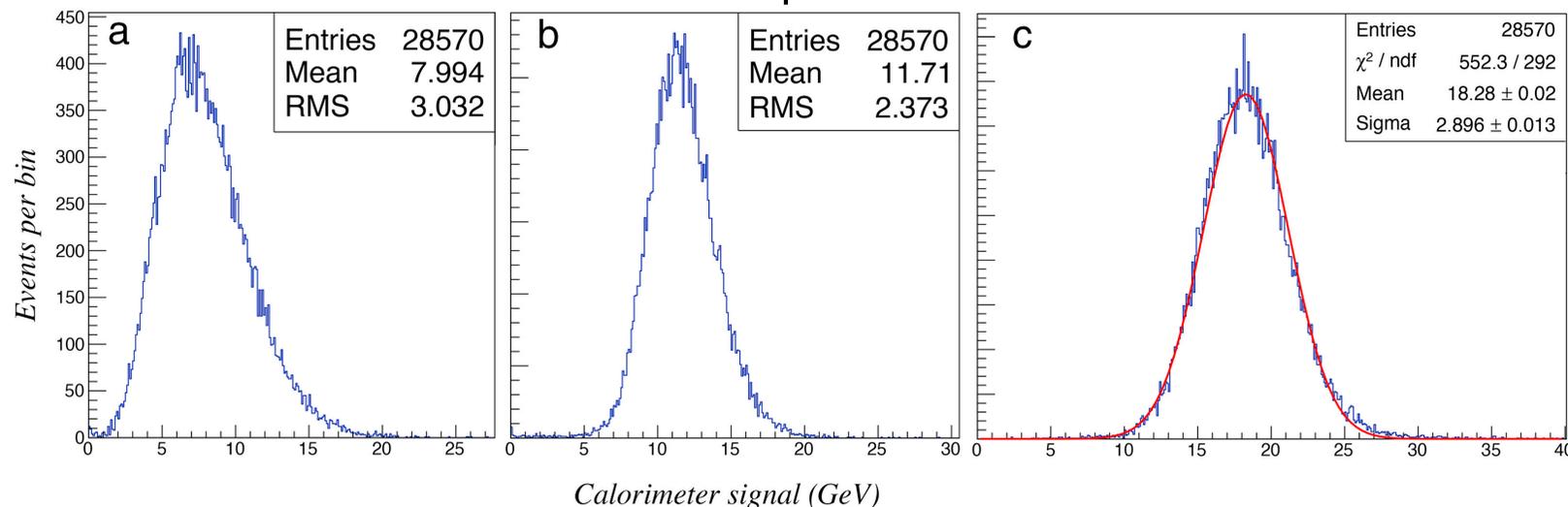
# Single hadron response - linearity

NIM A 866 (2017) 76

- Dual readout signal largely recovers linearity while vastly improving resolution

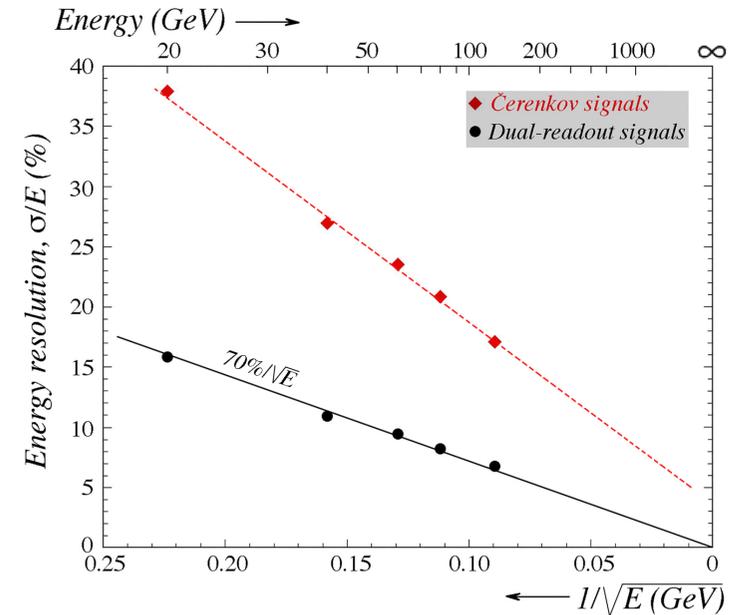
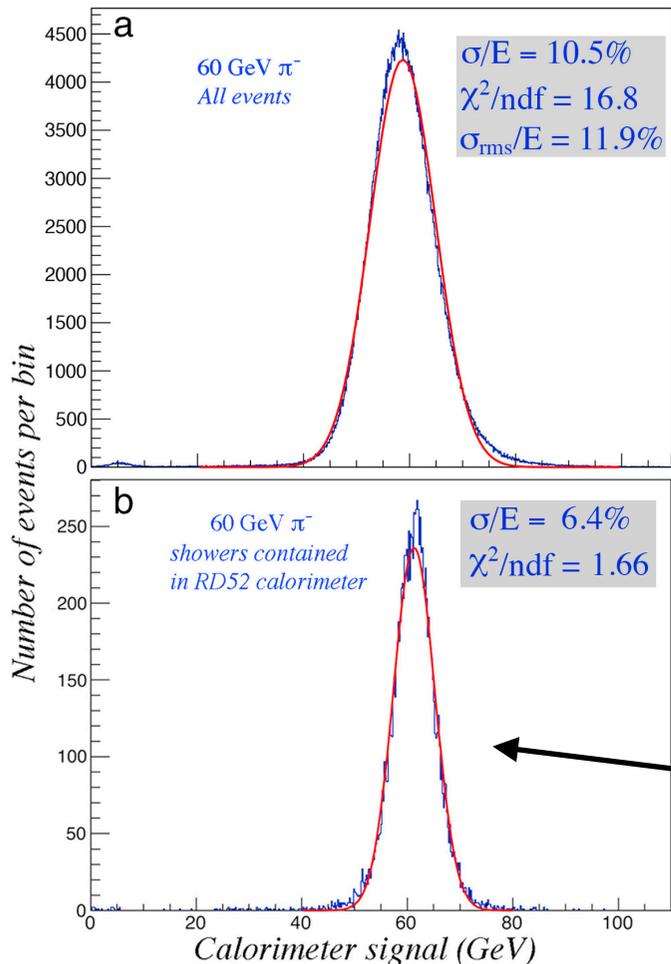


20 GeV pions



# Single hadron response - resolution

- Curse of calorimeter R&D: a fully containing calorimeter is expensive



No signal in leakage counters

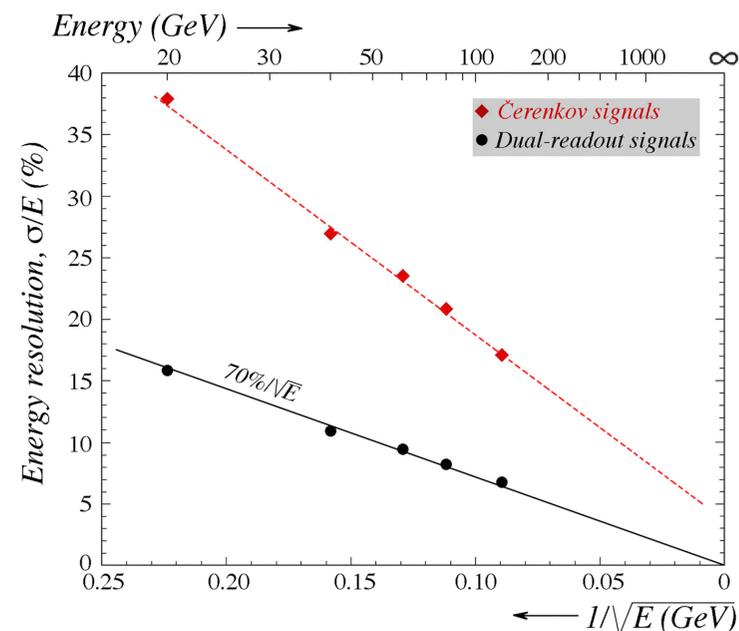


# Performance of Dual Readout

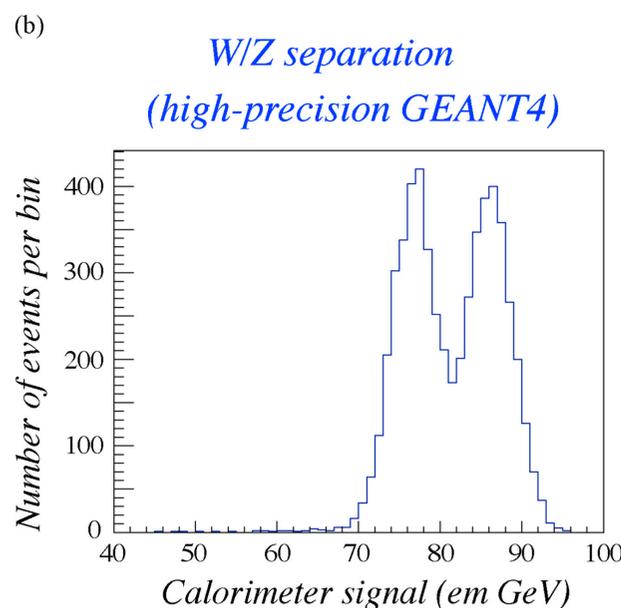
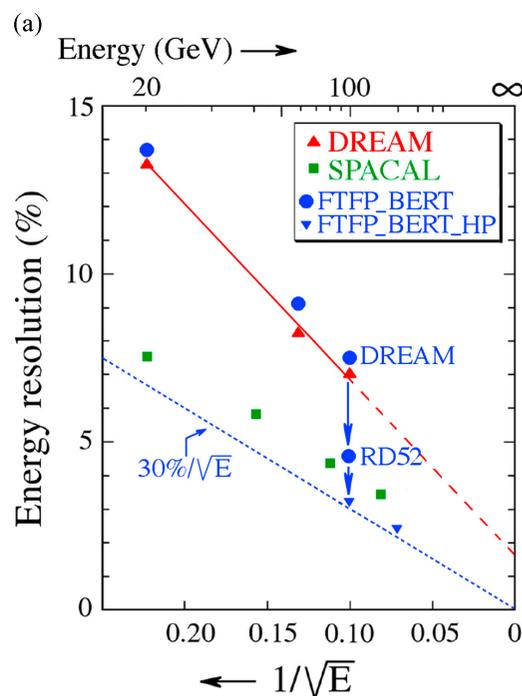
- **Hadronic resolution** comparable to **compensating calorimeters**.

- Resolution at TB (dominated by leakage). G4 estimate **with full containment**

$$\frac{\sigma}{E} = \frac{34\%}{\sqrt{E}}$$

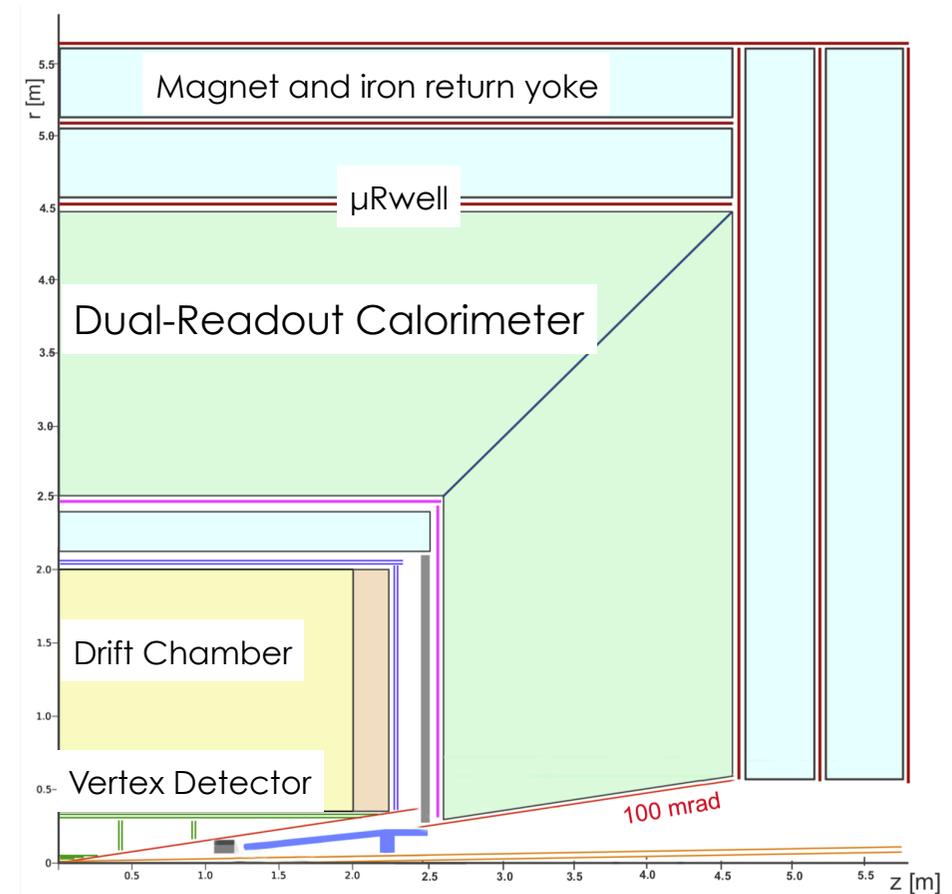
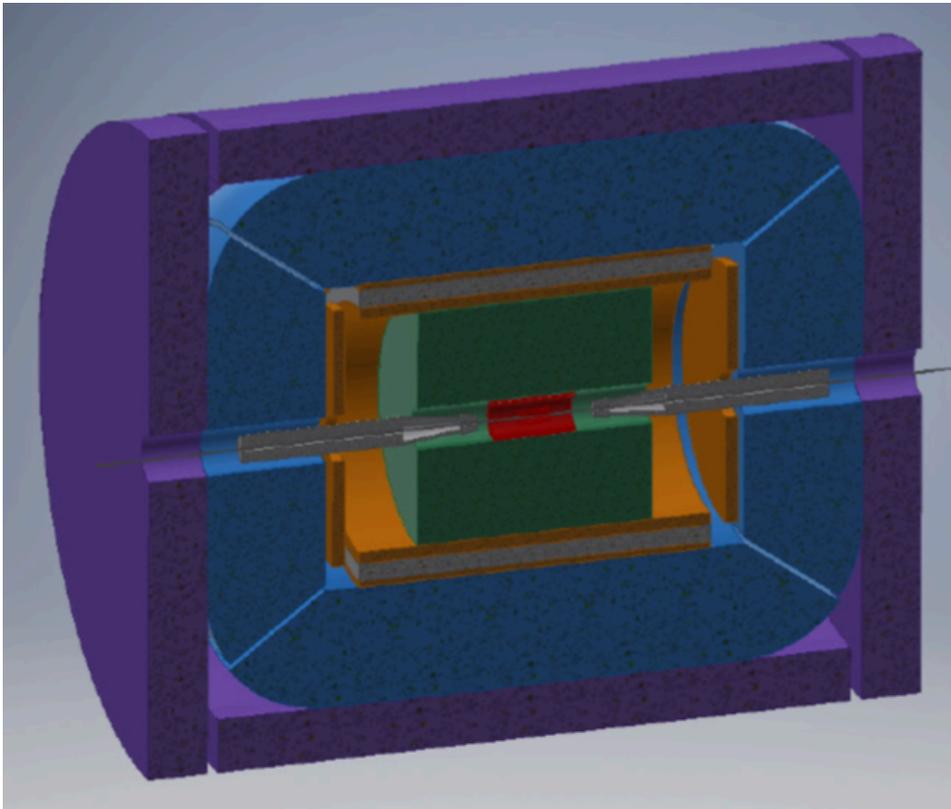


See <https://doi.org/10.1016/j.pnpnp.2018.07.003>



# Recent results

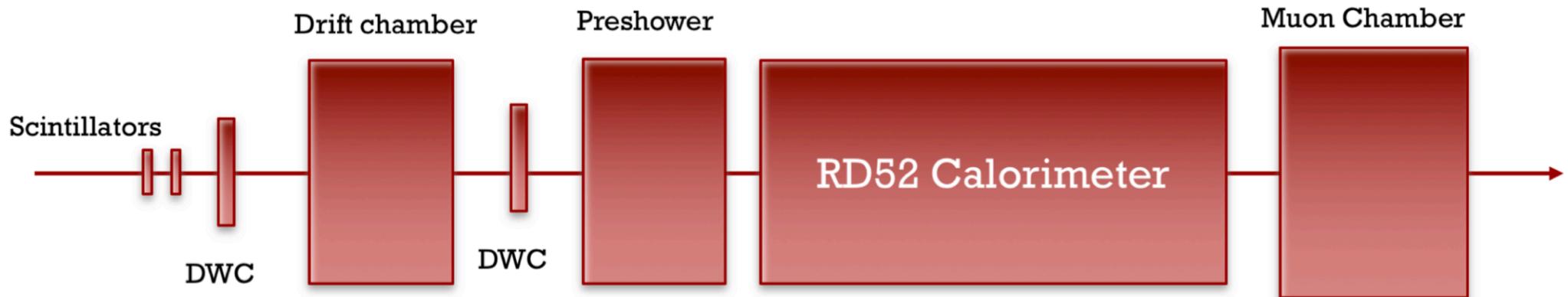
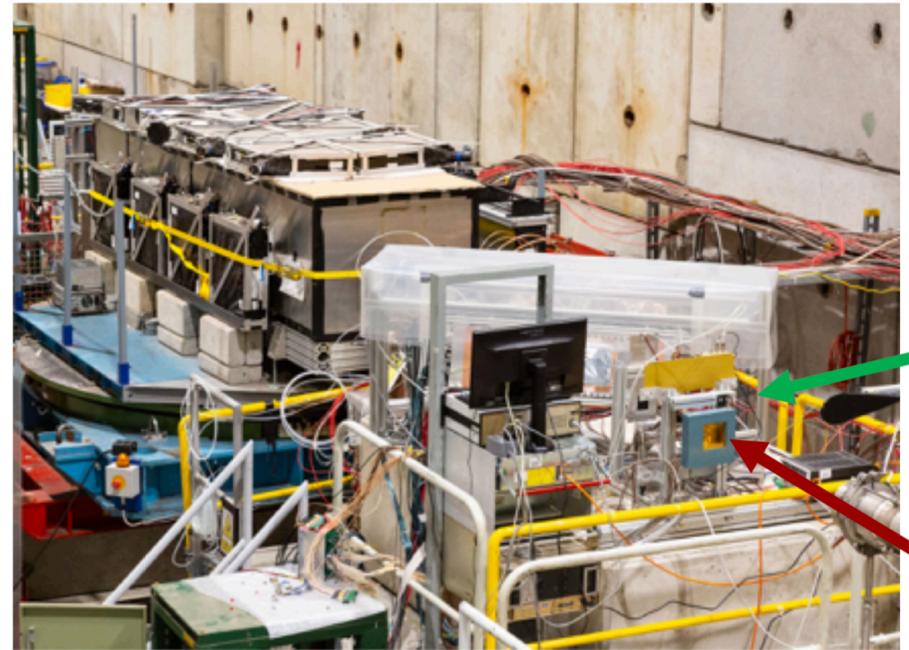
# A practical implementation: IDEA



See [here](#) for additional information

# IDEA slice on beam (2018)

- A full **combined test** of IDEA:
  - Drift chamber prototype
  - GEM as preshower +  $\mu$ RWell for  $\mu$  detection
  - Several calorimeter options tested on beam



# From PMTs to SiPM

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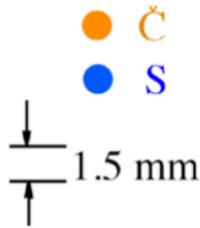
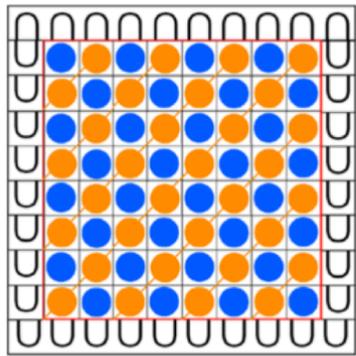
## SiPM pros:

- Compact readout
- Resilience to magnetic field
- Large light yields
- High readout granularity (particle flow “friendly”)
- Photon counting (calibration)
- High timing resolution

## SiPM cons:

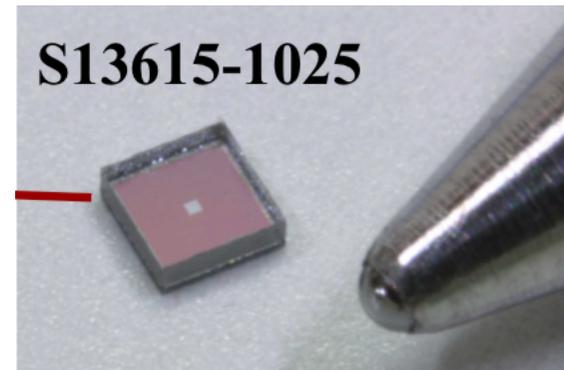
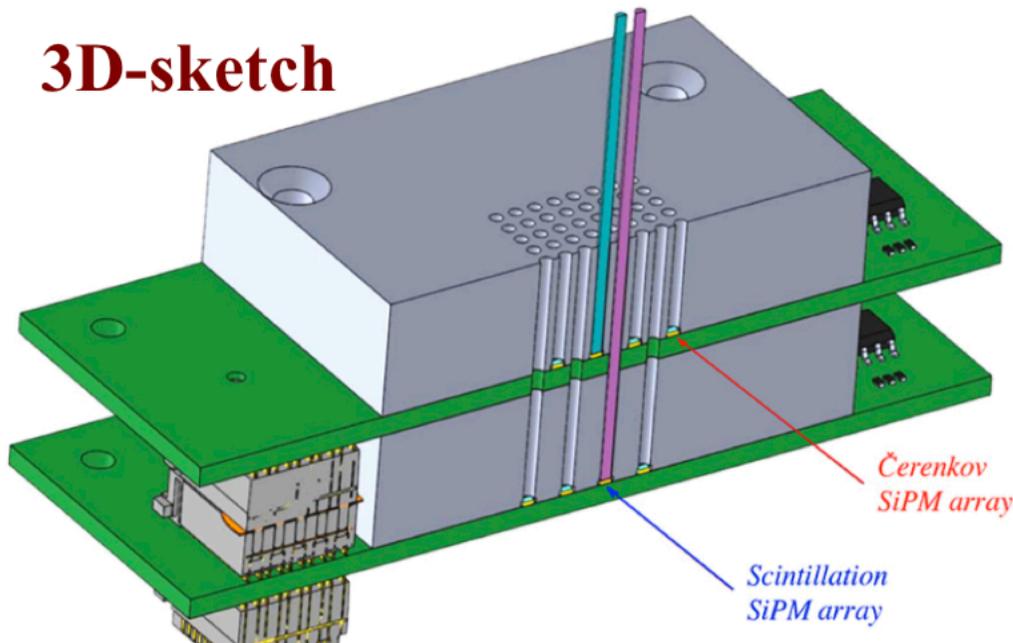
- Signal saturation/dynamic range
- Cross talk between C and S signals
- Instrumental effects

# SiPM dual readout

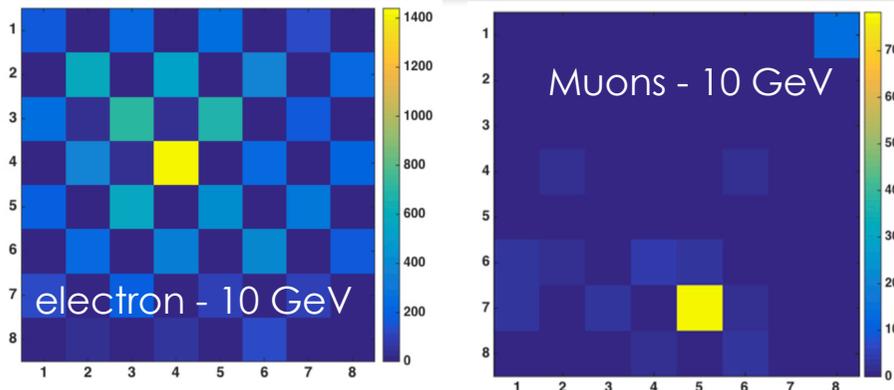


- Single fibre readout with **HAMAMATSU SiPM**
- Readout for Cherenkov and Scintillation light separated to minimise cross talk (the latter expected to be  $\sim 50$  times larger if not attenuated)

## 3D-sketch



# SiPM dual readout

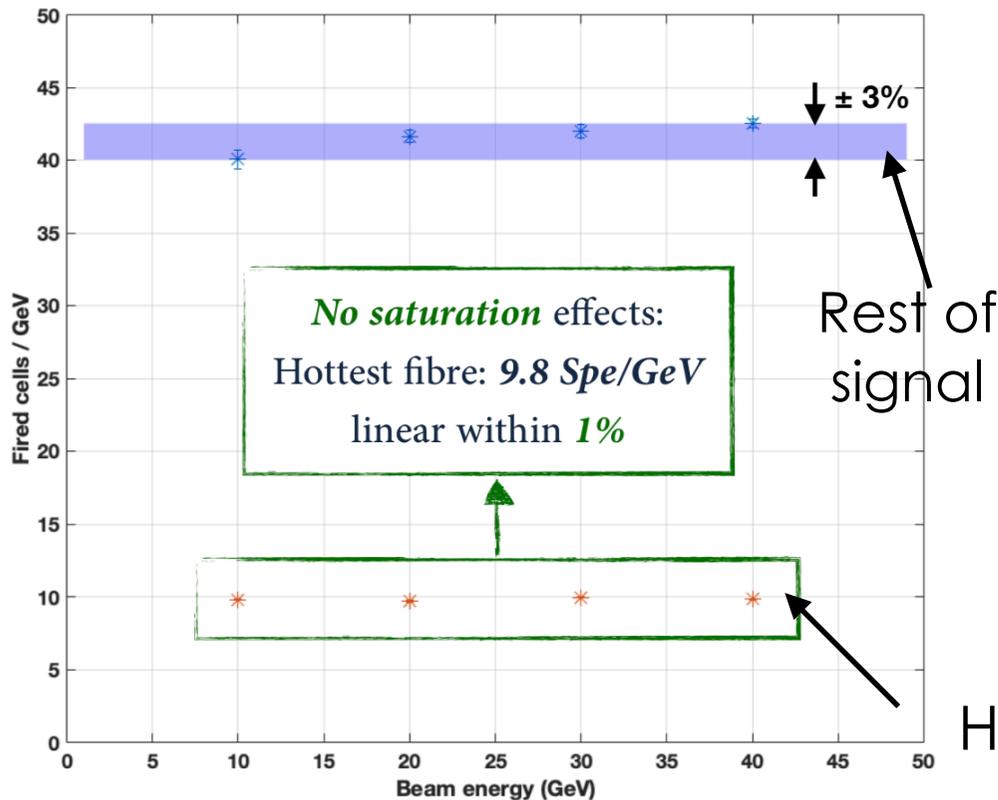


Operating with 5.5  $V_{OV}$  - PDE  $\sim$  22%

Cherenkov light yield (70 Spe/GeV)  
 $\sim$  a factor 2 larger than what  
 measured with PMT

(Filtered) scintillation light yield  
 under control ( $\sim$ 95 Spe/GeV).

EM stochastic term  $\sim$  10% is  
 achievable



Hottest fiber

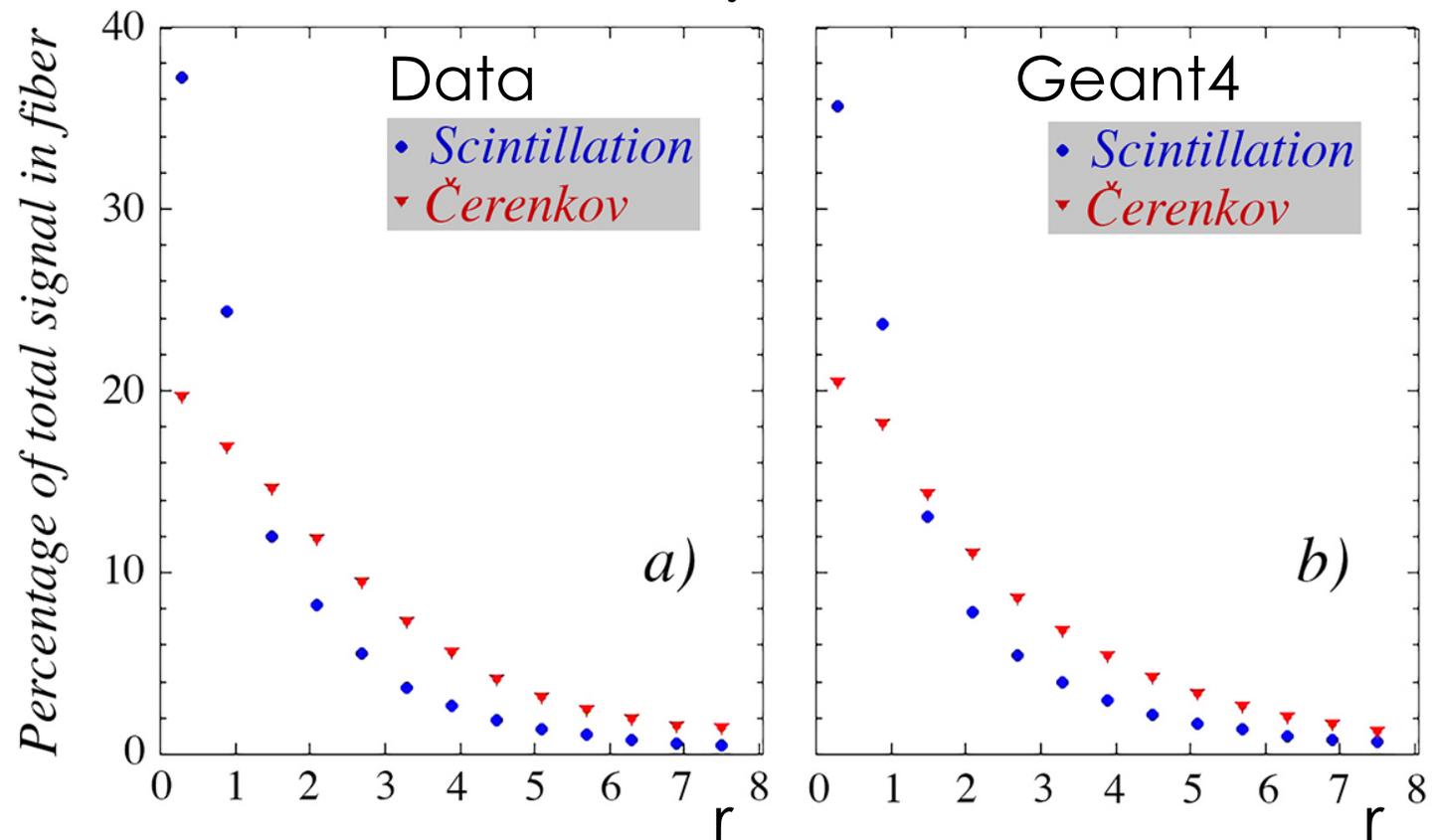
# SiPM dual readout (shower shape)

- Readout of single fibre gives **unprecedented lateral segmentation**
- Em lateral shower shape measured with **~ 1 mm precision.**

[Doi:10.1016/j.nima.2018.05.016](https://doi.org/10.1016/j.nima.2018.05.016)

$$\bar{x} = \frac{\sum_i x_i E_i}{\sum_i E_i}; \bar{y} = \frac{\sum_i y_i E_i}{\sum_i E_i}$$

$$r = \sqrt{(x_i - \bar{x})^2 + (y_i - \bar{y})^2}$$

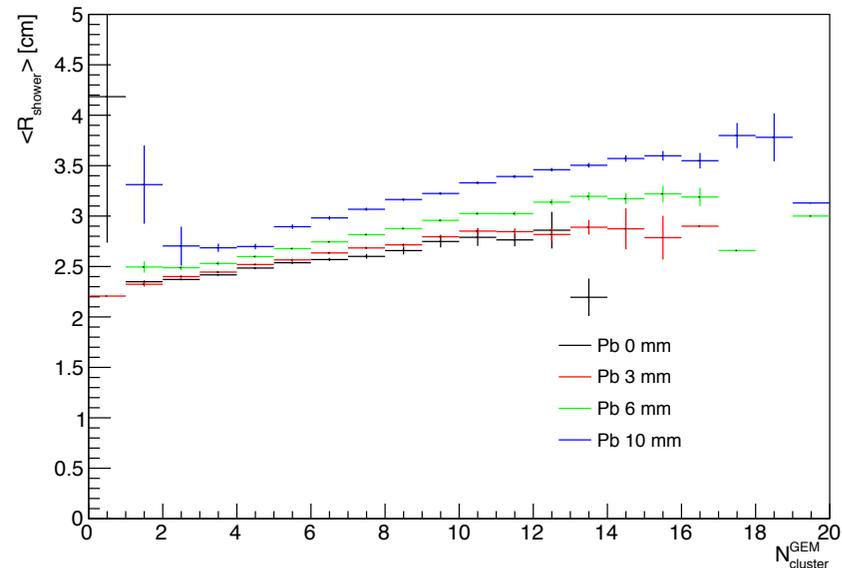
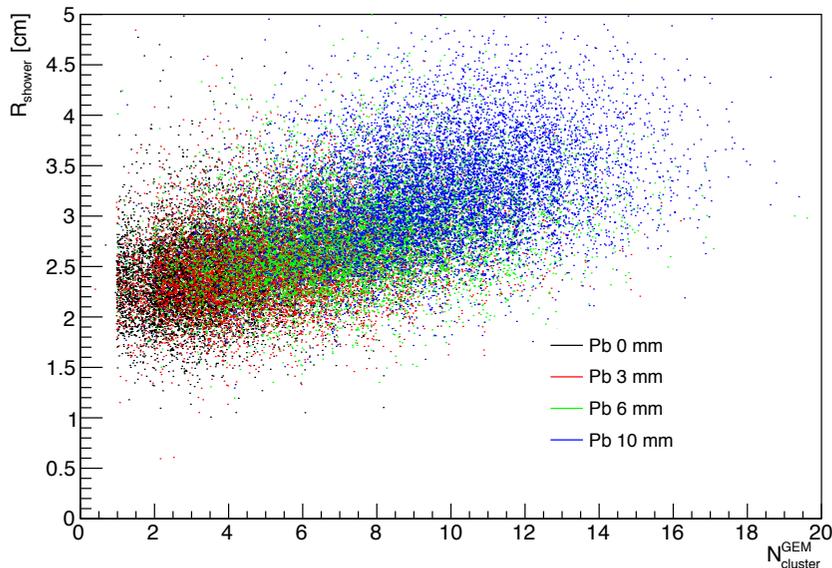


# Combined measurements (RD52)

- RD52 performance studied in detail elsewhere
  - Focus on DAQ combination and combined runs with GEM-based preshower

$$R_{\text{shower}} = \frac{\sum_{\text{ch}} E_{\text{ch}} \cdot \sqrt{x_{\text{ch}}^2 + y_{\text{ch}}^2}}{\sum_{\text{ch}} E_{\text{ch}}}$$

Shower width from 5 mm Pb + additional material correlates with number of clusters in GEM preshower



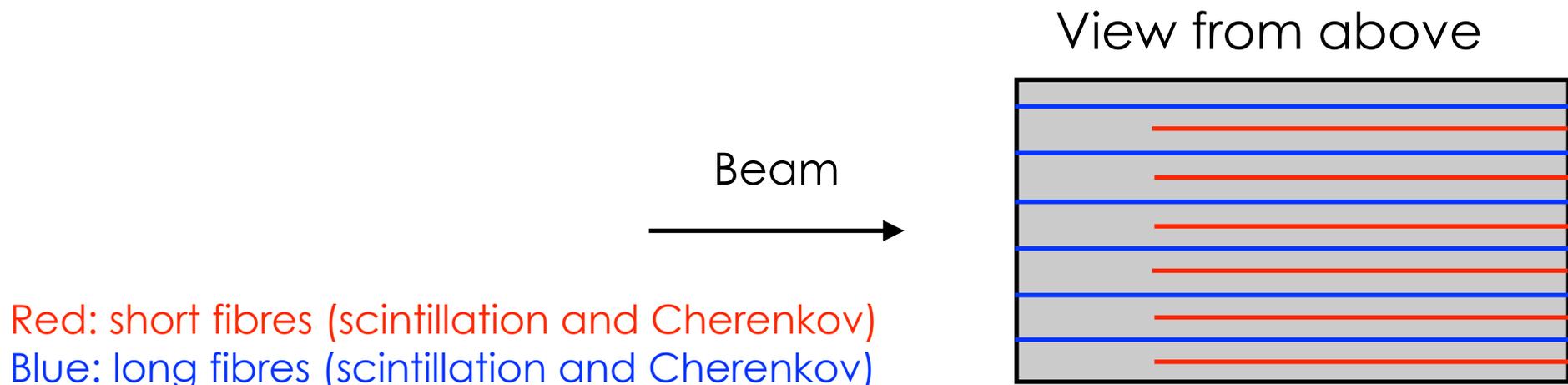
# Longitudinal segmentation (standalone test)

- Particle identification (e.g. hadronic tau decay) may benefit from **longitudinal segmentation**.
- “Staggered” option tested on beam

“HAD” section: E(short fibres)

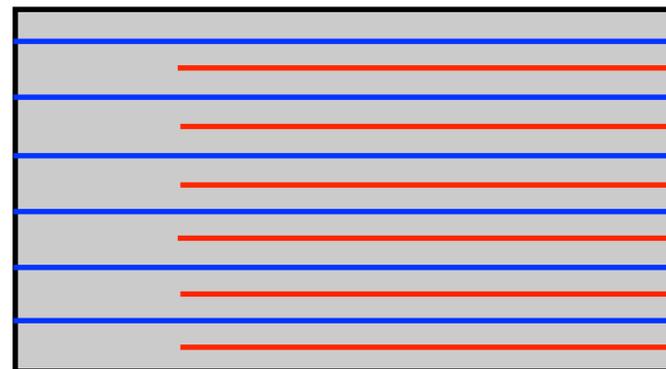
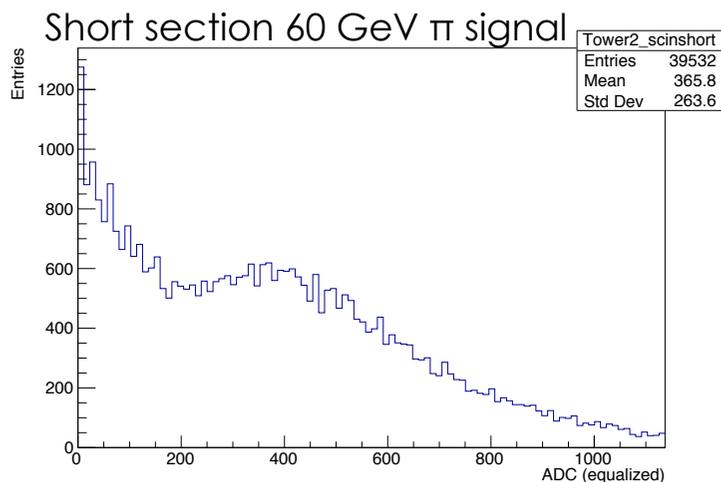
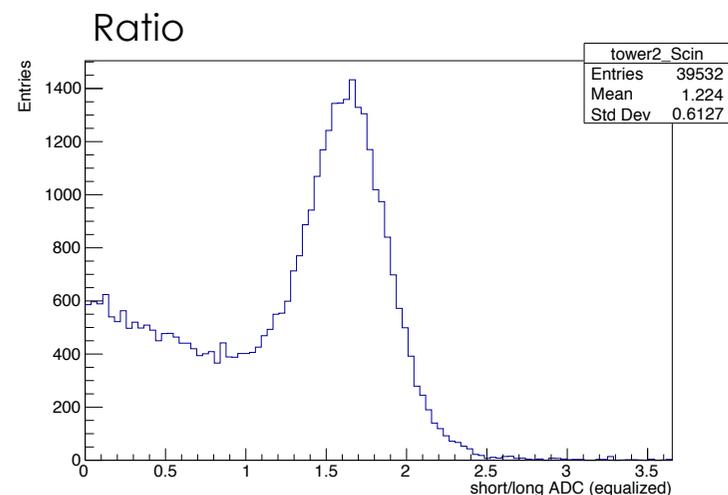
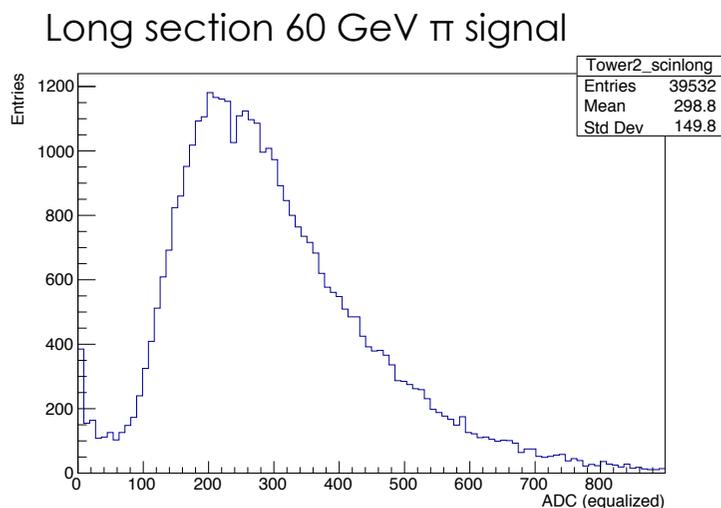
“EM” section: E (long fibres) - E (short fibres)

- Challenge: calibration of the short section.



# Longitudinal segmentation

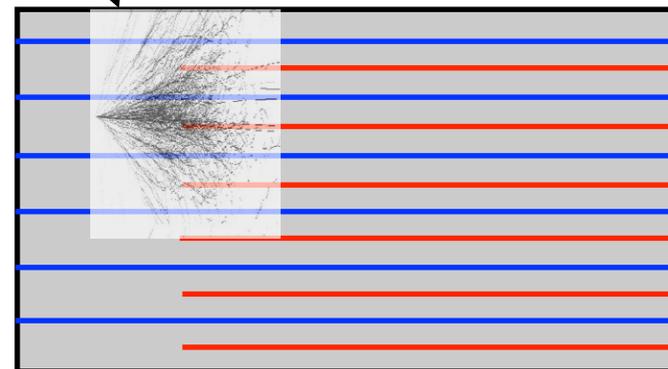
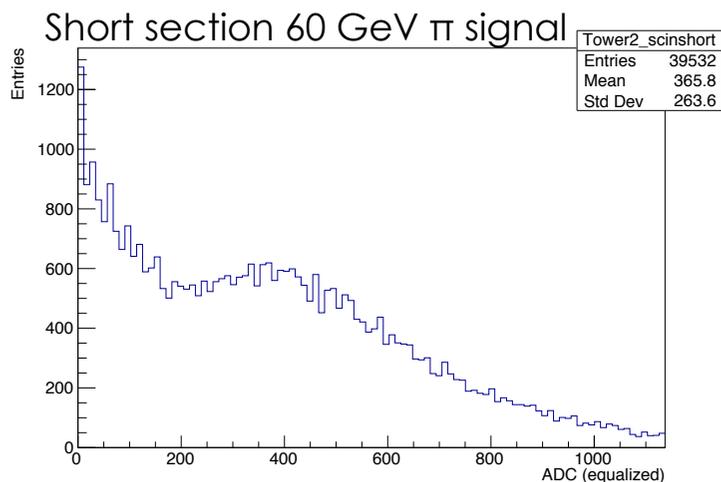
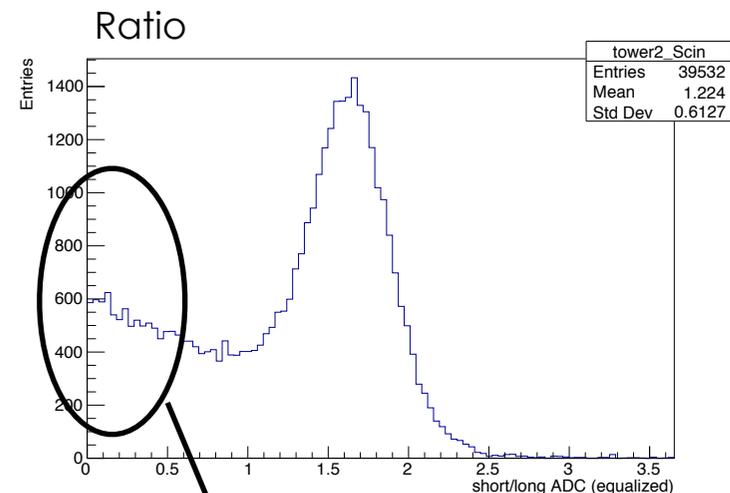
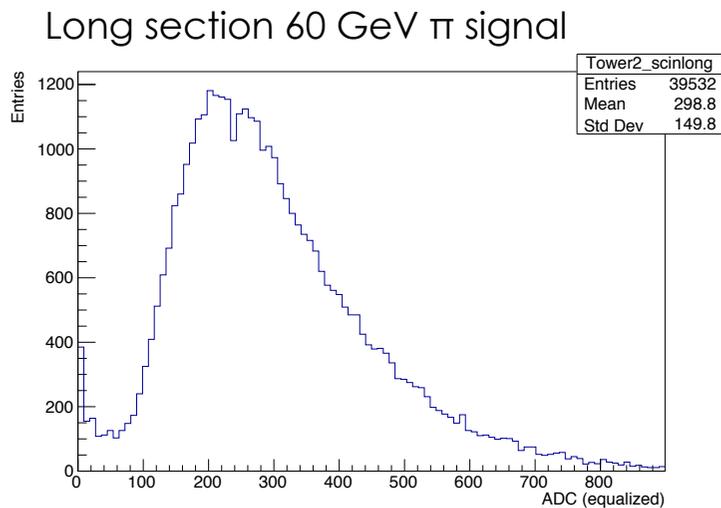
- Problem: how do we calibrate short section?
- Idea: **propagate calibration** from long section **using hadrons**.



Scintillation fibers

# Longitudinal segmentation

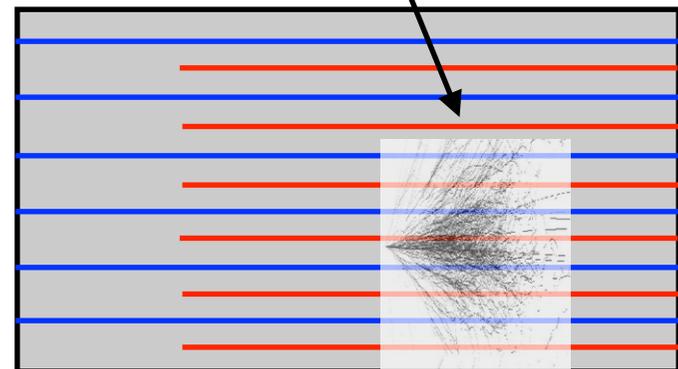
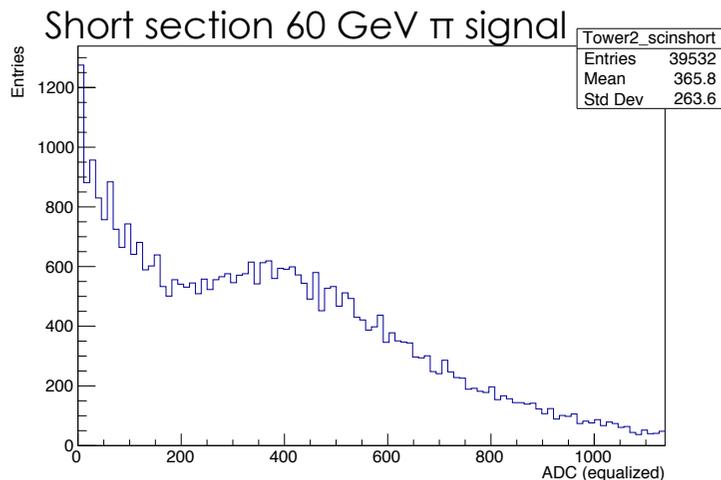
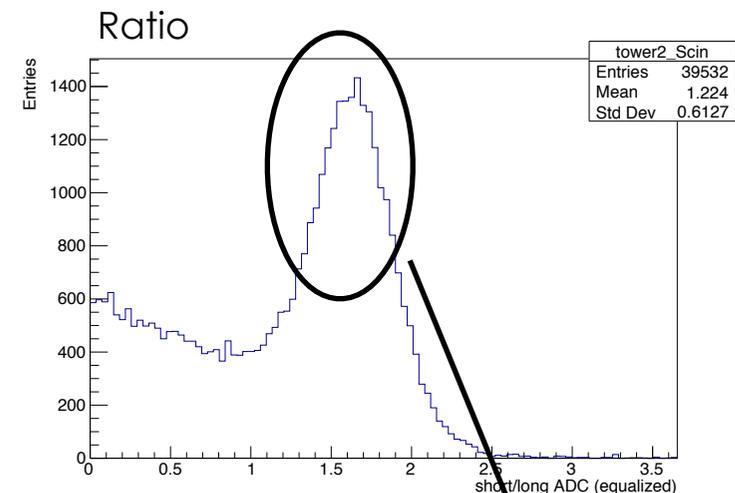
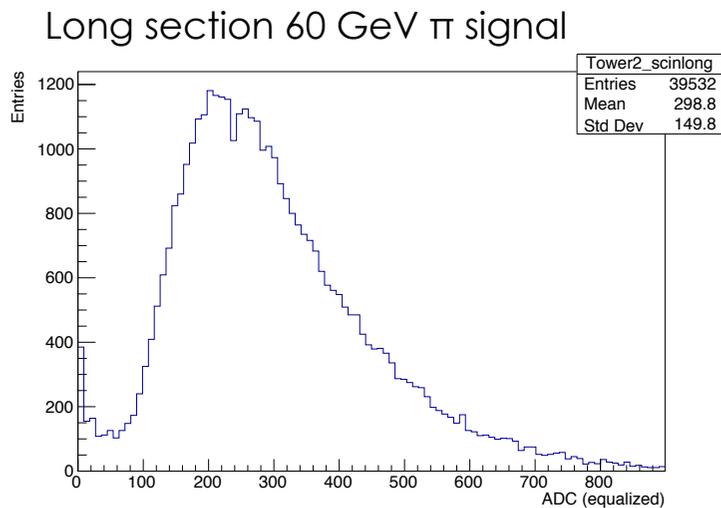
- Problem: how do we calibrate short section?
- Idea: **propagate calibration** from long section **using hadrons**.



Scintillation fibers

# Longitudinal segmentation

- Problem: how do we calibrate short section?
- Idea: **propagate calibration** from long section **using hadrons**.

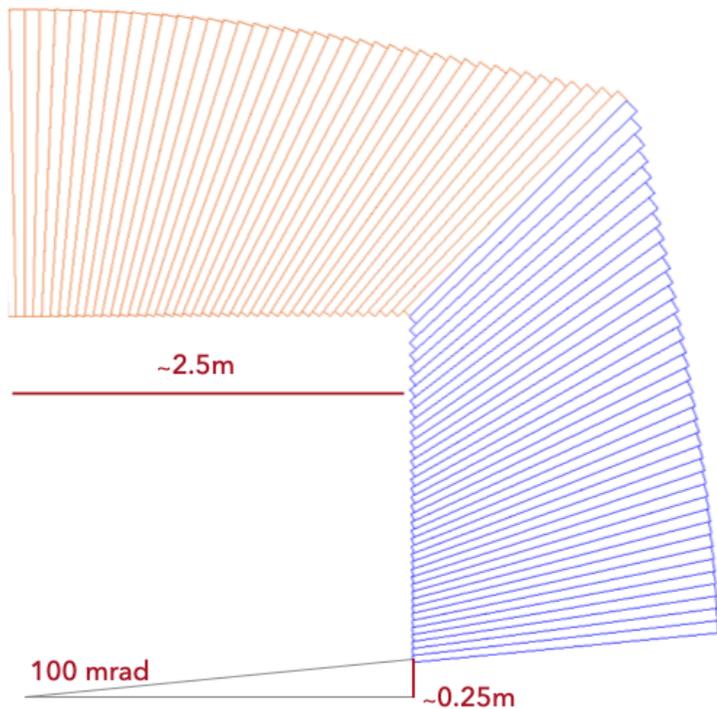


Scintillation fibers

Open questions

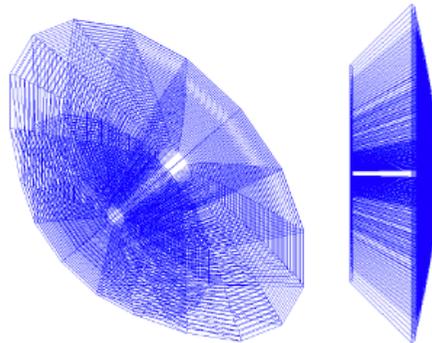
# Simulation

- Final choices need **reasonably well understood G4** simulation
- Existing fast sim + **full G4 sim projective geometry**



**Barrel (right):** 40 towers  
Inner length 2.5m  
Tower height 2.0m

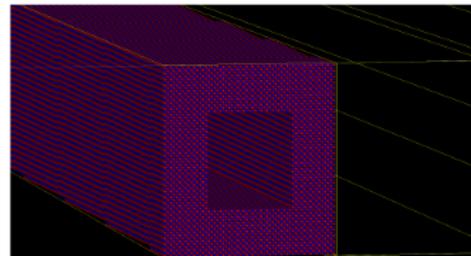
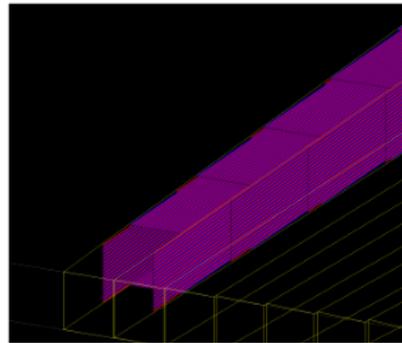
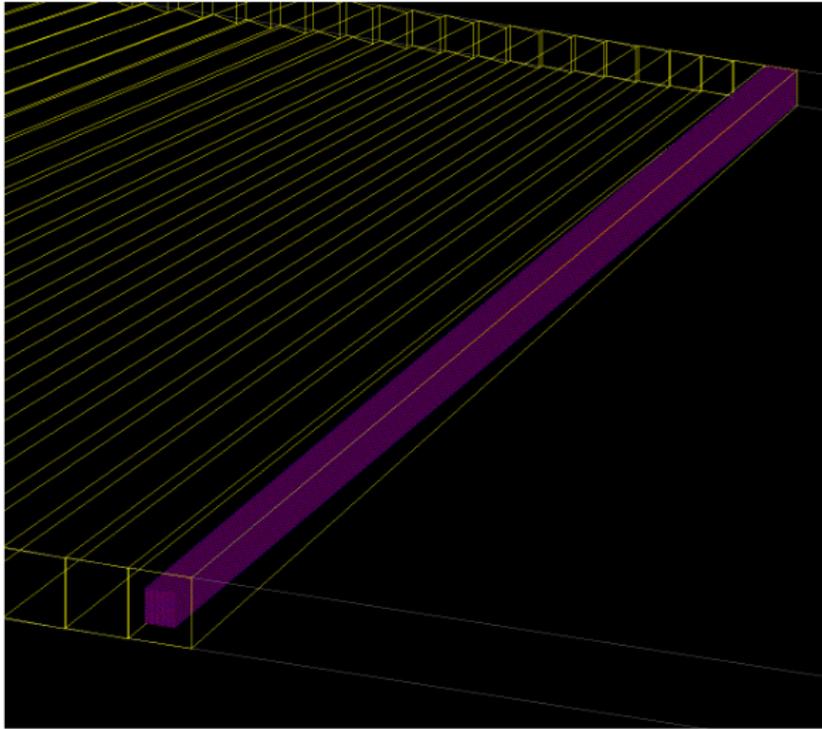
**Endcap (right):** 39 towers  
Inner length 2.25m  
Tower height 2.0m



- Open points:
  - Reliability of G4 to reproduce  $\chi$  value.
  - More structured software effort for full IDEA project

# Absorber and layout

- Copper/brass calorimeter would **avoid problems with e/mip**
  - But lead more cost-effective
  - Options to be finalised with G4 simulation



**70447104 fibers in the barrel**

- Readout with **longitudinally distributed fibres** is the baseline:
  - Excellent lateral segmentation and **electronics all outside the calorimeter**
  - But **large number of fibres/SiPM** readout channels needed
- Other options being investigated:
  - Tiles
  - More creative options

Courtesy of L. Pezzotti

# Longitudinal segmentation - use timing?

- Profit from **excellent SiPM time resolution** - derive **energy deposit longitudinal position** with time
- More studies with simulation needed

SiPM readout (2):

digitiser (ASIC) & feature extraction (FPGA)

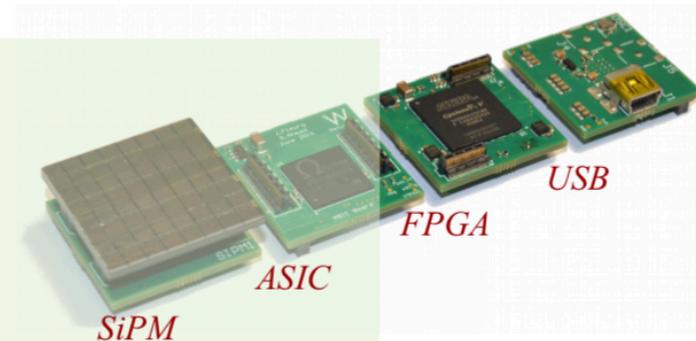
→ started investigating

→ time stamps w/  $O(100 \text{ ps res})$

→ get shower longitudinal development

$[\Delta x \sim 5 \text{ cm} \Rightarrow \Delta t \sim 100 \text{ ps}]$

→ started looking at neural network implementation



# Costs

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- Total foreseen cost at the moment ~ 150 M€.
  - Dominated by optical elements (fibre + SiPM)

	Unit cost	Quantity	Total (M€)
<b>Fibres</b>	0.25 €/m	190 M x 1.2 m	57.4
<b>SiPM</b>	0.25 € each	190 M	47.7
<b>Absorber</b>	2000 €/ton	3.4 kton	6.7
<b>Front-end</b>	-	-	17.9

- Actively looking at ways of reducing these costs

# Summary

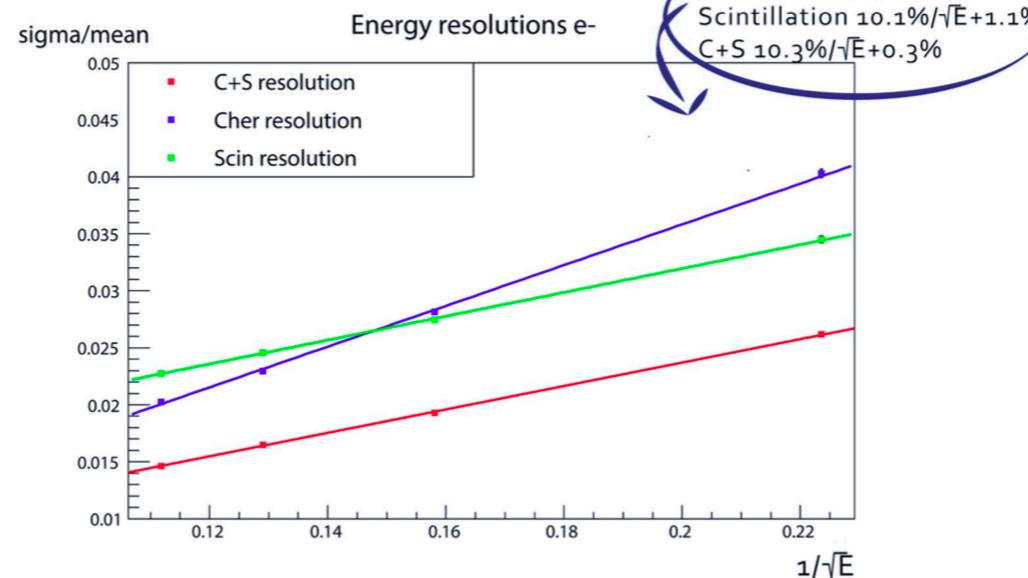
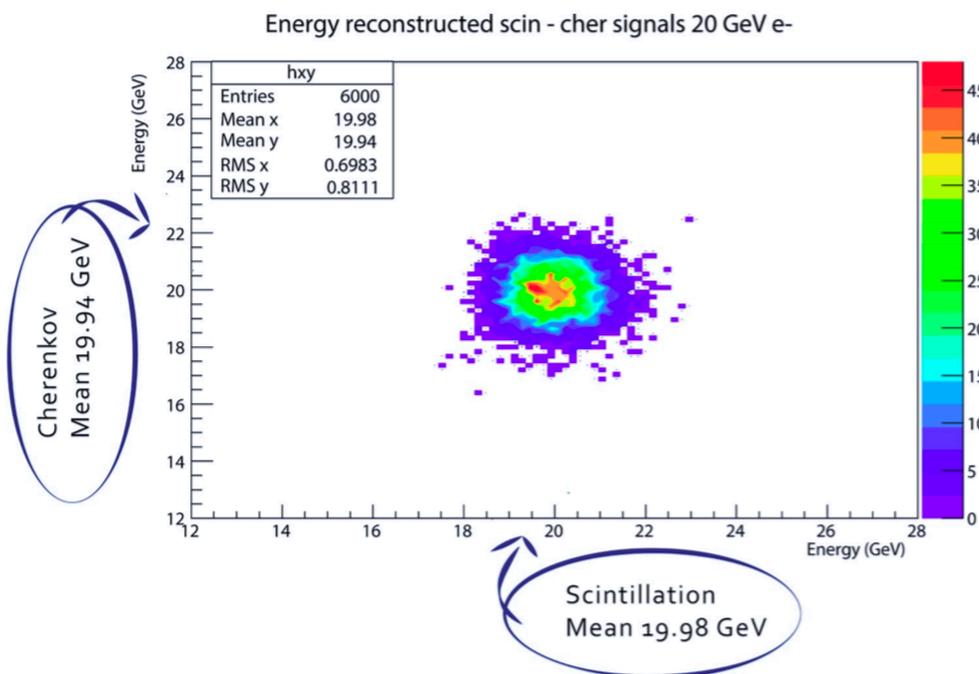
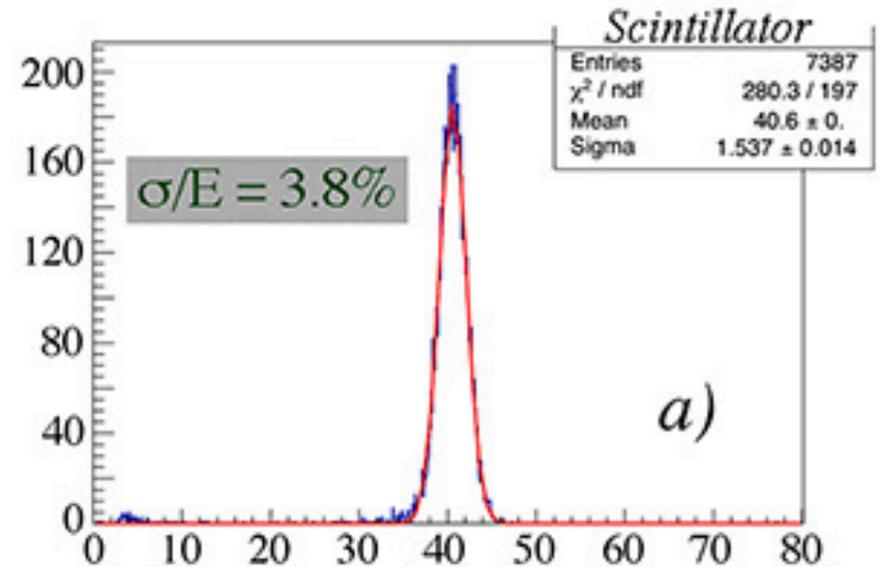
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- Dual readout:
  - **Complementary** principles w.r.t. particle flow.
  - Excellent EM and HAD **native resolution**.
  - Could be combined with pFlow approach if need be.
- Towards a prototype. Help welcome on each of the open questions

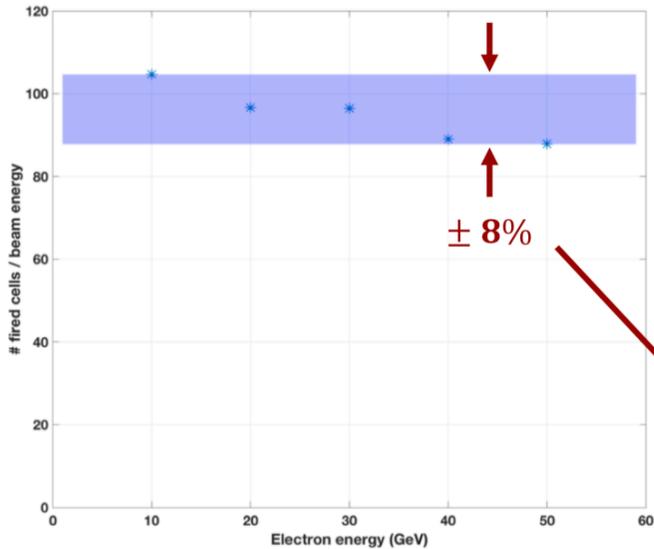
Backup

# EM performance

- Excellent e and  $\gamma$  calorimeter performance thanks to high sampling fraction.
- EM and HAD calorimetry in one device.



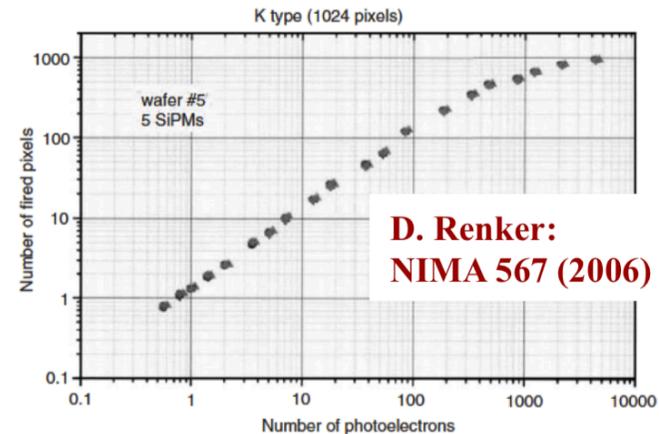
## Results from 2017 test beam



- Detector operated at 0.5V over breakdown (PDE  $\approx$  2%)
- Temperature stability correction:
  - $< 0.5^\circ\text{C}$  during a single run (negligible)
  - $< 2^\circ\text{C}$  during the full scan (considered)
- PDE correction for temperature variation

Even if the SiPMs are not saturated with this setting, they are working in a strongly non linear regime: a correction is required

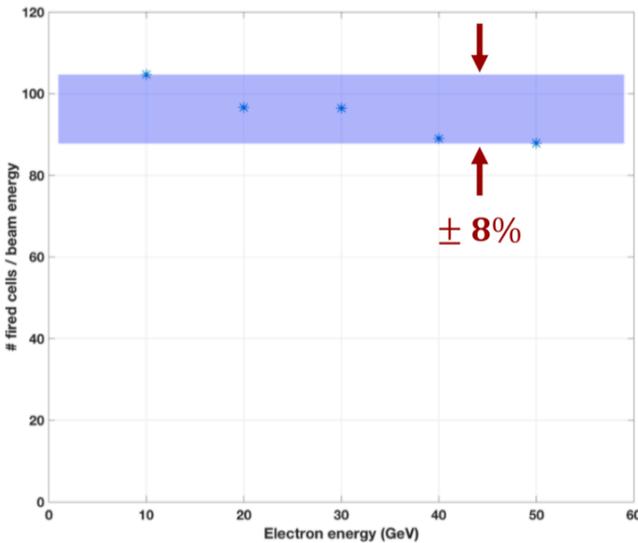
$$N_{fired} = N_{total} \times \left[ 1 - e^{-\frac{N_{photons} \times PDE}{N_{total}}} \right]$$



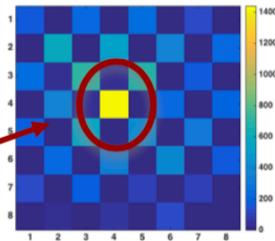
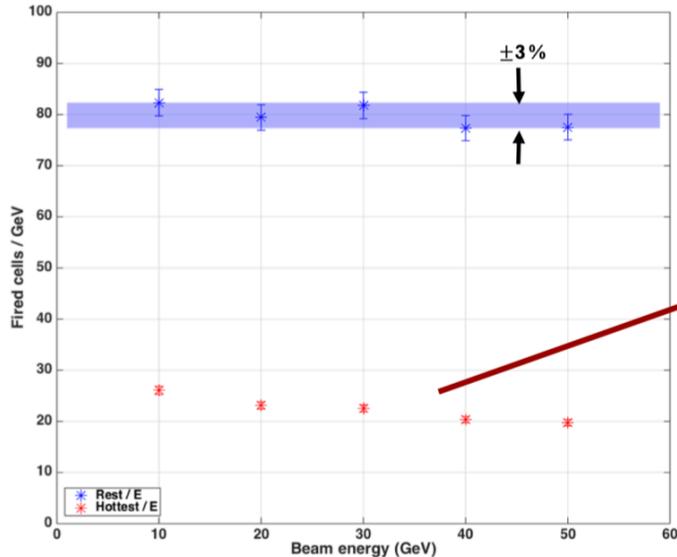
Valid as a first approximation: the light uniformly illuminate the SiPMs, all photons come at the same time and spurious effects are negligible

# Scintillation light

## Results from 2017 test beam



- Detector operated at 0.5V over breakdown (PDE  $\approx 2\%$ )
- Temperature stability correction:
  - $< 0.5^\circ\text{C}$  during a single run (negligible)
  - $< 2^\circ\text{C}$  during the full scan (considered)
- PDE correction for temperature variation



Once the correction is applied, the linearity is improved even if it is not fully recovered (i.e. signal from the seed)

**To reduce this effect we decided to attenuate the scintillating light using a yellow filter**

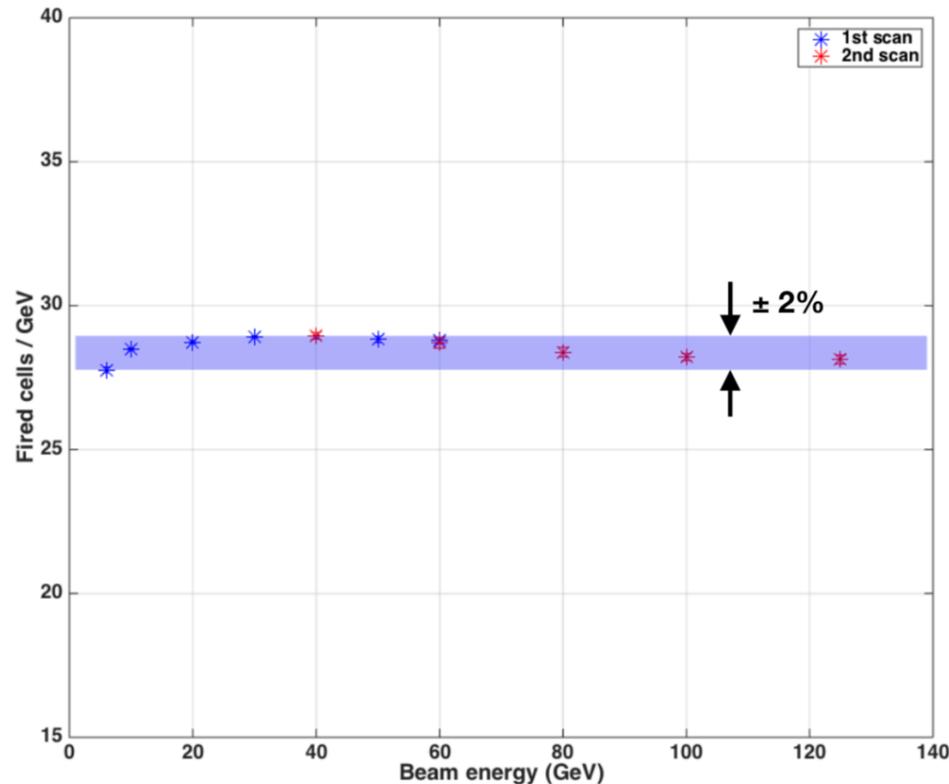
# Cherenkov light

## Cherenkov light yield:

$V_{\text{Bias}} = 5.5 V_{\text{ov}}$  (57.5 V) and  $PDE \sim 25\%$ .

$\sim 28.6$  Cpe/GeV, 2% linear from 6 to 125 GeV.

Correcting for 36% e.m. energy containment:  $\sim 69 \pm 5$  Cpe/GeV.



More than **2 times larger** than what measured with the previous\* PMT-based modules.

### Example:

Stochastic term of RD-52 e.m. resolution could be improved from  $\sim 14\%/\sqrt{E}$  up to  $\sim 12.5\%/\sqrt{E}$ .  
(sampling fluctuations:  $\sim 9\%/\sqrt{E}$ ).

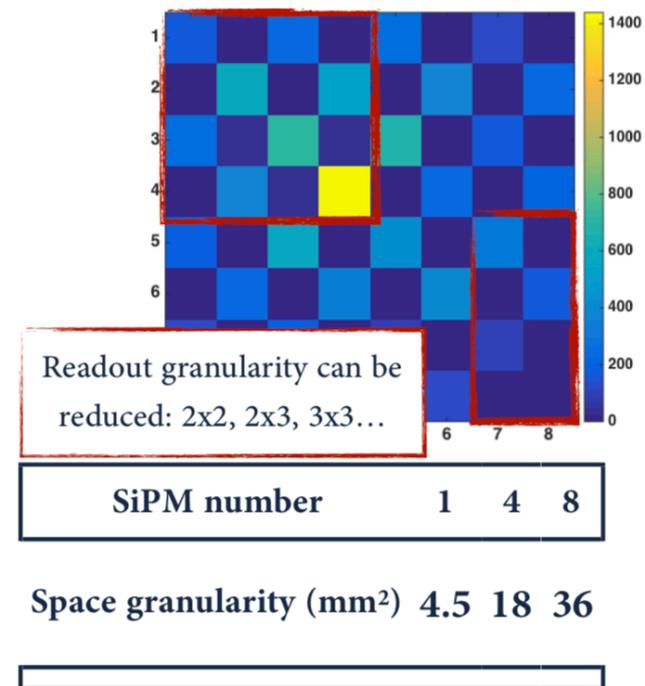
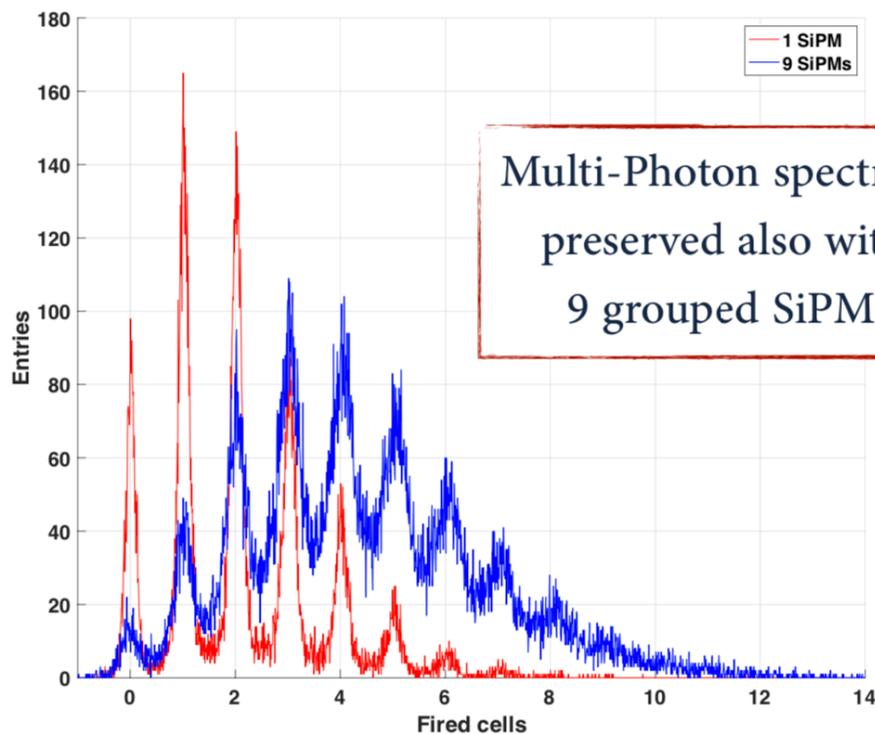
# Grouping channels

In a full scale module, the number of *readout channels* will be of the order of  $10^8$ .

The possibility to **sum up the analog output** is under study:

Number of SiPM that can be grouped guarantying the *Multi-Photon spectrum*.

SiPM *dynamic range*: sensors have to operate in a *linear regime*.



# Outline

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- Why dual readout?
- Towards a combined detector: IDEA test beam
  - Combined data taking with old prototypes
  - SiPM readout
  - Longitudinal segmentation
- Summary

# EM showers: relevant numbers

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- Radiation length  $X_0$ : typical scale of longitudinal shower development

$$X_0 \sim 1433 \frac{A}{Z(Z+1)(11.32 - \ln Z)} \left( \frac{\text{g}}{\text{cm}^2} \right) \implies \sim \frac{1}{Z}$$

- Critical energy  $E_C$  (below which ionisation takes over bremsstrahlung)

$$E_C \sim \frac{160 \text{ MeV}}{(Z + 1.24)} \implies \sim \frac{1}{Z}$$

- Molière radius  $R_M$  (typical scale of lateral shower development)

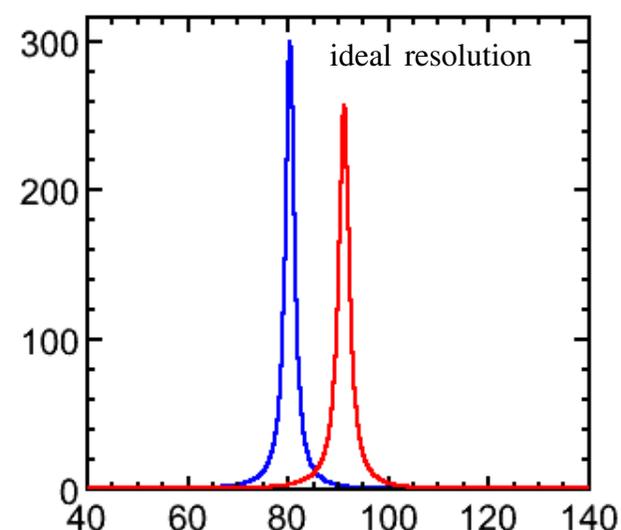
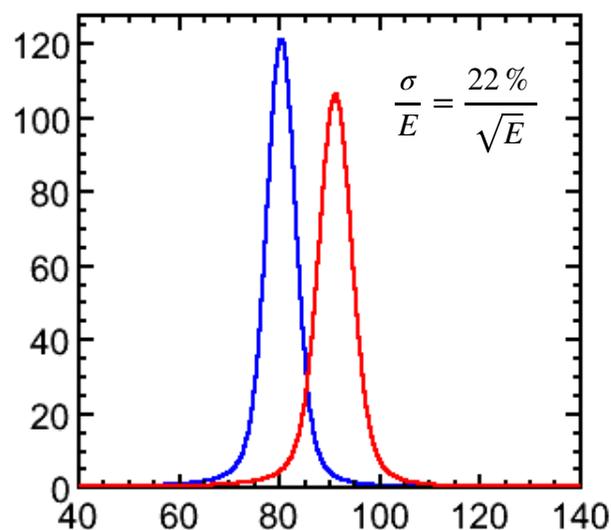
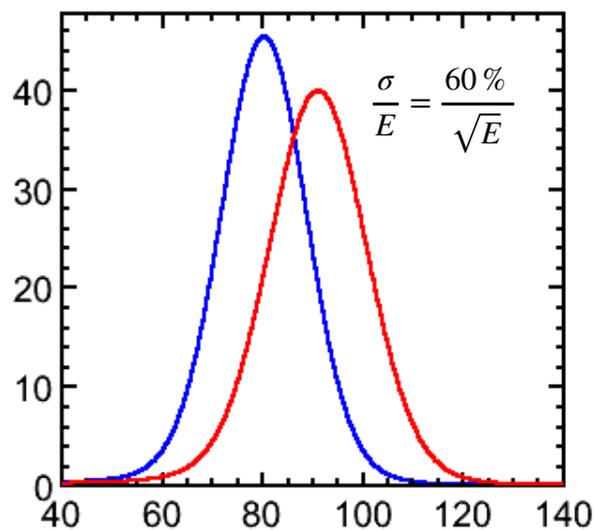
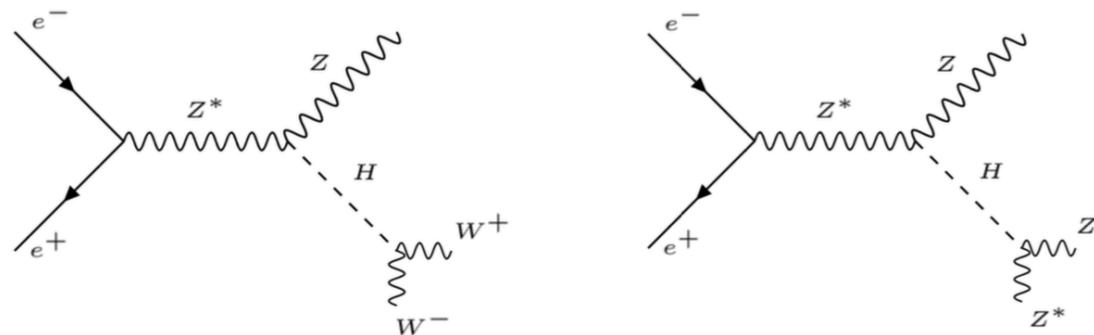
$$R_M \sim \frac{X_0 \times 21.2 \text{ MeV}}{E_C} \implies \sim \text{const}$$

- Shower depth (shower maximum), where the multiplication process stops

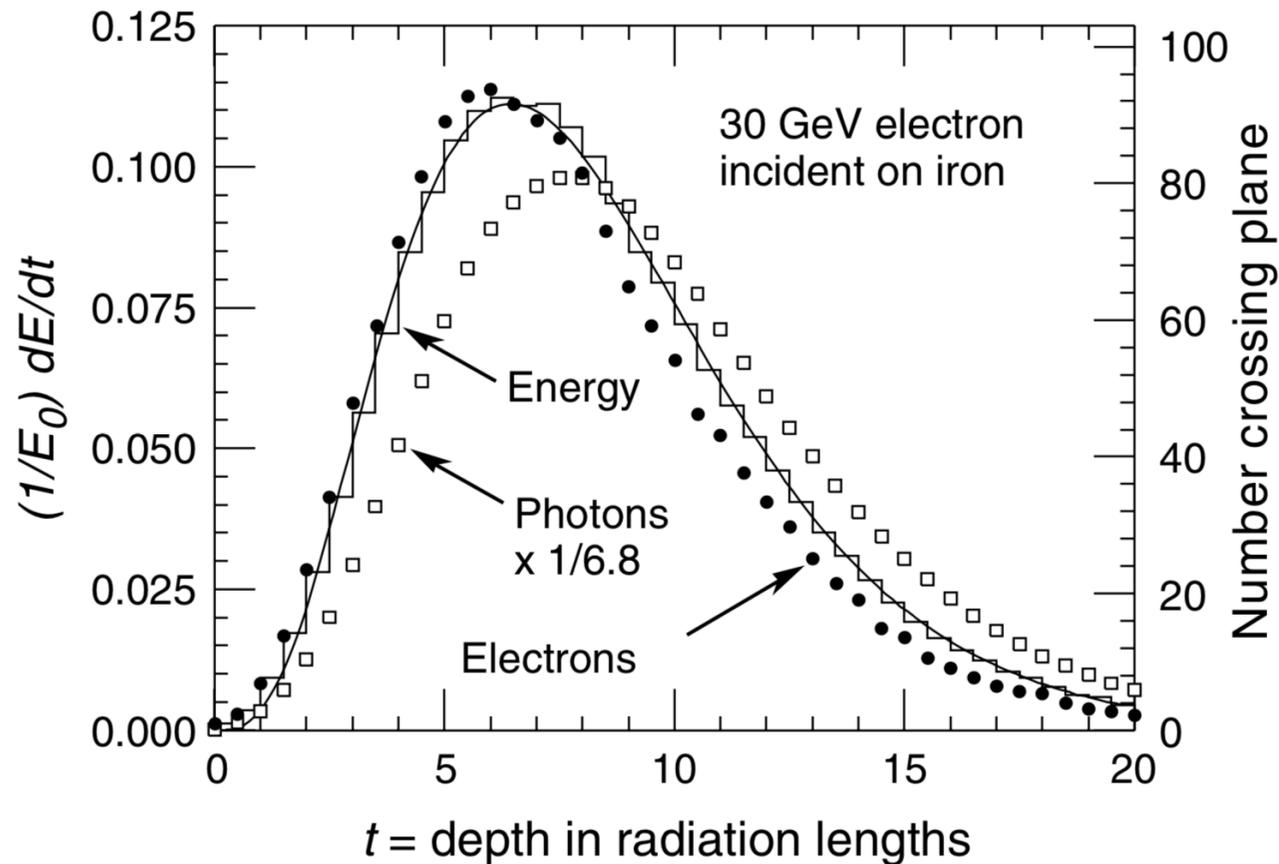
$$X_{\text{max}} = X_0 \frac{\ln \left( \frac{E}{E_C} \right)}{\ln 2} \implies \sim \frac{1}{Z}, \sim \ln E$$

# Reminder: why dual readout?

- Precision physics at  $e^+e^-$  collider calls for high-resolution hadronic calorimetry



# EM shower



- Large number of particles involved
- Regular shape
- Small event-by-event fluctuations

# $f_{em}$ energy dependence

- Simple model: **only pions** are produced at each interaction, **respecting isospin symmetry**

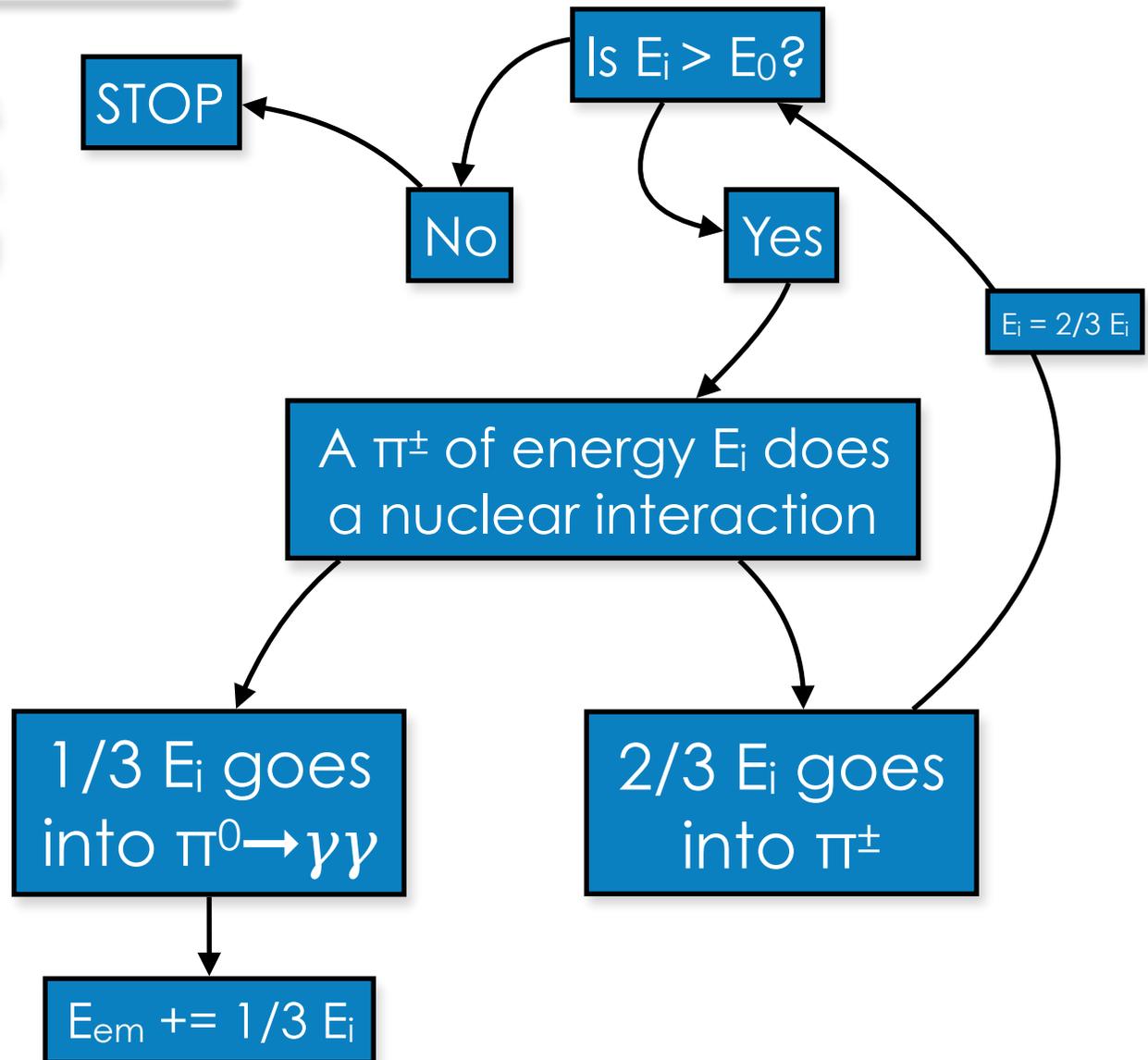
- Then the math is:

$$f_{em} = \frac{E_{em}}{E} = 1 - \left(\frac{2}{3}\right)^n$$

$$\left(\frac{2}{3}\right)^n E = E_0$$

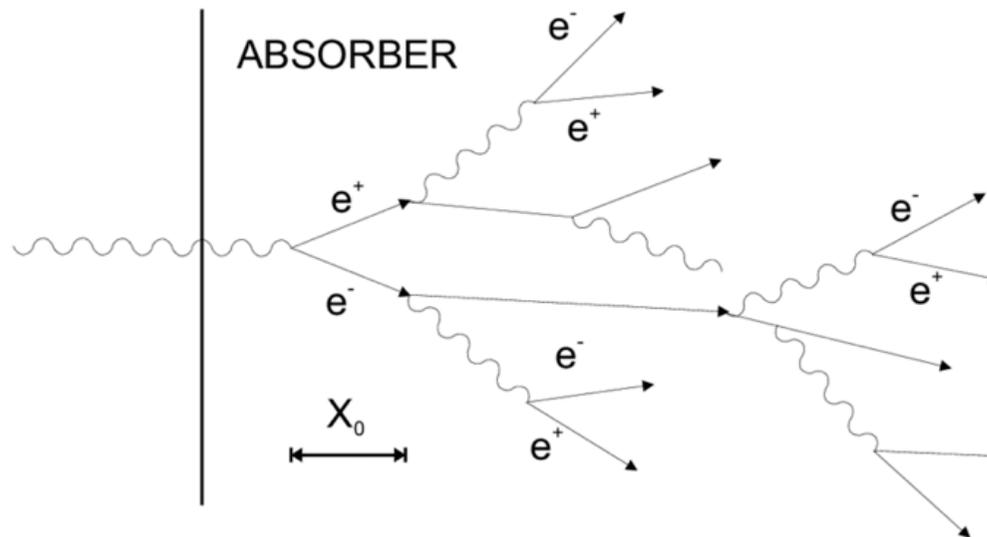
$$n = p \ln \frac{E}{E_0} \quad (\text{for some number } p)$$

$$f_{em} = 1 - \left(\left(\frac{2}{3}\right)^{\ln \frac{E}{E_0}}\right)^p = 1 - \left(\frac{E}{E_0}\right)^{k-1} \quad (\text{for some number } k)$$



# Calorimetry - a primer

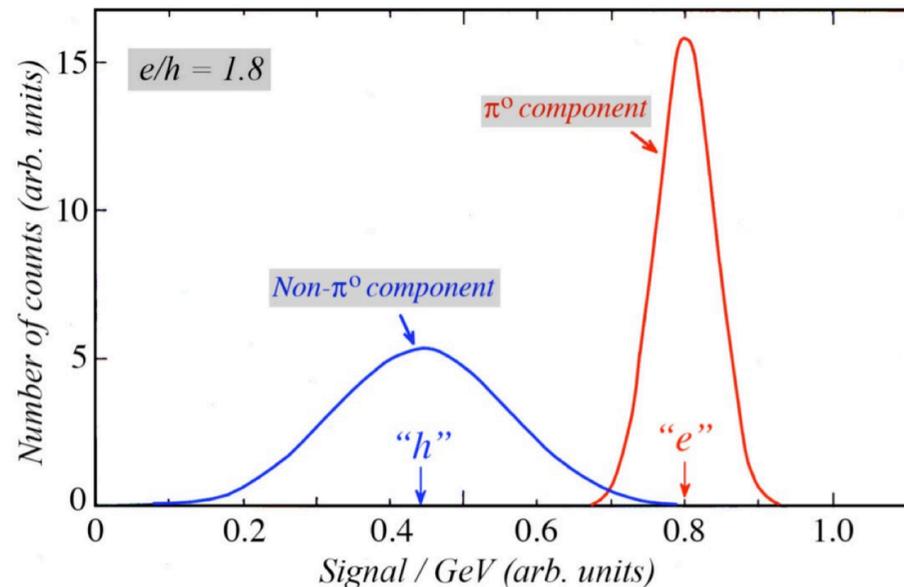
- Electromagnetic shower:
  - Driven by **EM interactions with atom EM field**
  - **Moderate** intrinsic interaction intensity, but **big targets** (scale to bear in mind  $\sim 10^{-10}$  m)
  - Main mechanisms at work: **bremsstrahlung, pair production, ionisation.**
  - **Lots of particles** involved



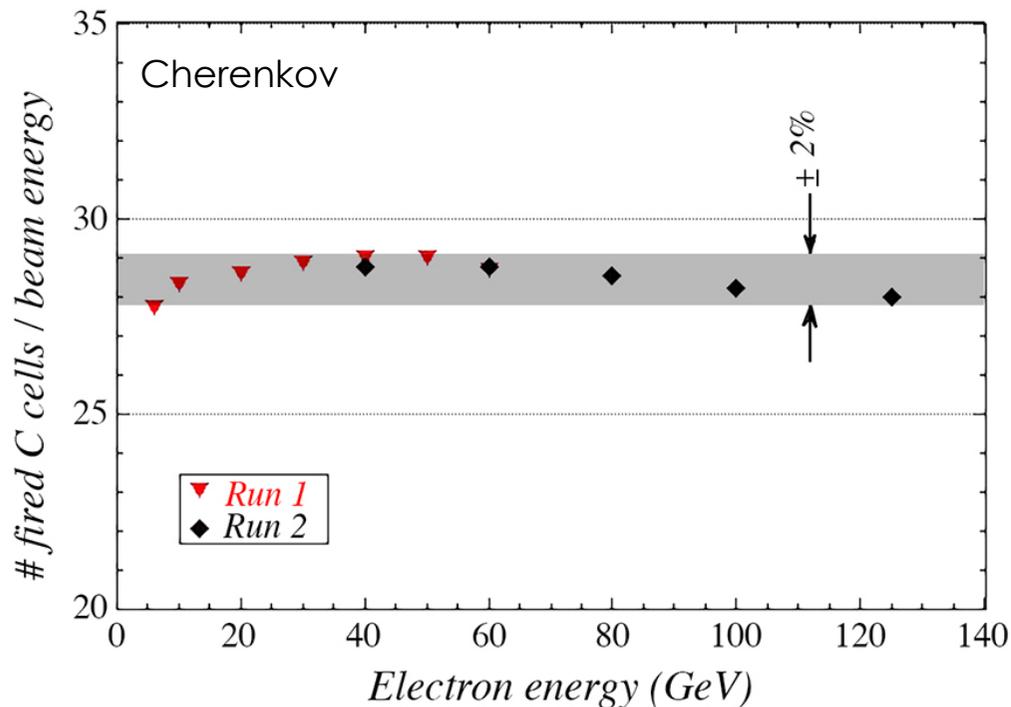
# Measuring hadronic showers

- Definitions:

- **e** is the efficiency of the measurement of the **EM component**
- **h** is the efficiency of the measurement of the **HAD component**
- **e/h** is a characteristic number **of the calorimeter**
- Normally **e/h > 1** (mostly because of invisible component), and the calorimeter is said to be **non-compensating**.

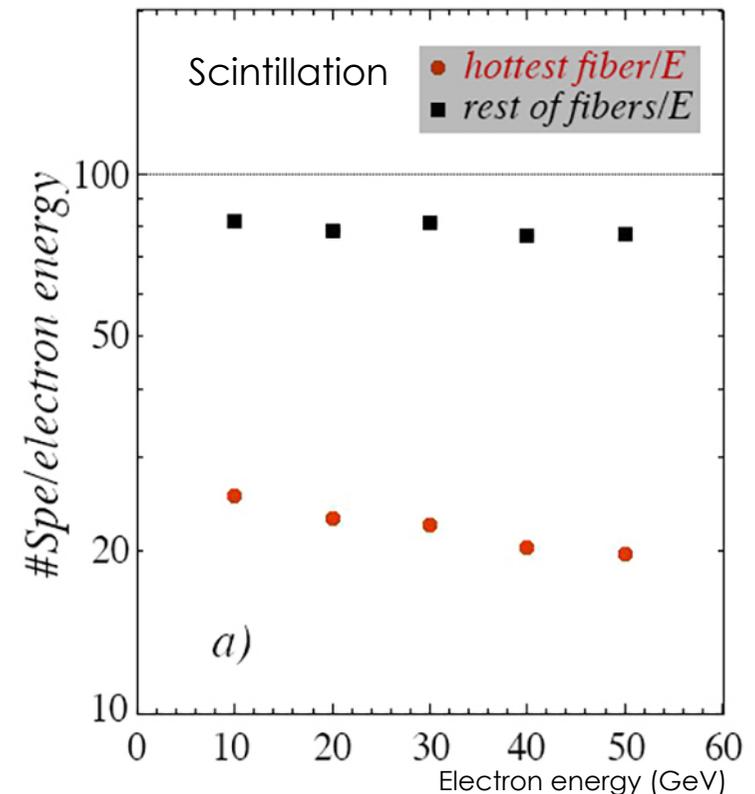


# SiPM readout - results (test beam 2017)



- Scintillation response showing evidence of saturation (addressed in 2018 test)

- Cherenkov light yield linear with beam energy over a wide range
- Light yield 28 Spe/GeV:
  - After correcting for containment ~ 55 Spe/GeV.



# Calorimeter options used during TB

- **RD52 module** (combined data taking with other sub detectors)
- **SiPM-based** readout (standalone)
- **"Staggered"** module (standalone)

